

21 cm Cosmology with HERA

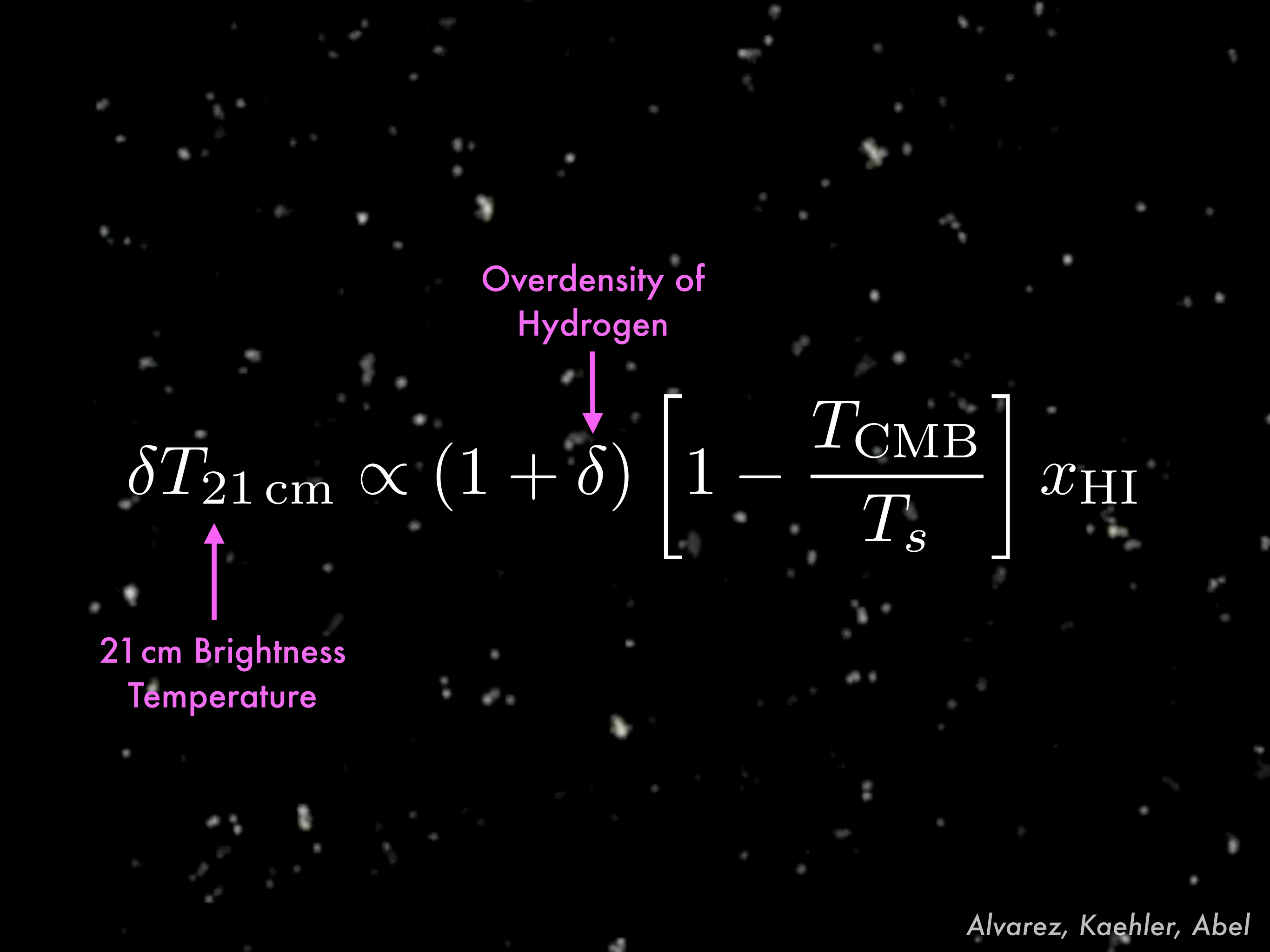
Josh Dillon
UC Berkeley

So we think the cosmic dawn
looked something like this...

So we think the cosmic dawn
looked something like this...

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

↑
21 cm Brightness
Temperature



The background of the slide is a Cosmic Microwave Background (CMB) fluctuation map, showing a complex pattern of white and grey spots against a black background, representing temperature variations in the early universe.

Overdensity of Hydrogen

$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$

21 cm Brightness Temperature

Overdensity of Hydrogen

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

21 cm Brightness Temperature

Spin Temperature

The diagram illustrates the equation for the 21 cm brightness temperature fluctuation. A pink arrow points from the text 'Overdensity of Hydrogen' to the δ term in the equation. Another pink arrow points from the text '21 cm Brightness Temperature' to the $\delta T_{21 \text{ cm}}$ term. A third pink arrow points from the text 'Spin Temperature' to the T_s term in the denominator of the bracketed expression.

The diagram illustrates the relationship between the 21 cm brightness temperature and several physical parameters. The central equation is:

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

Annotations with arrows indicate the following dependencies:

- An upward arrow from "21 cm Brightness Temperature" points to $\delta T_{21 \text{ cm}}$.
- A downward arrow from "Overdensity of Hydrogen" points to δ .
- An upward arrow from "Spin Temperature" points to T_s .
- A downward arrow from "Neutral Fraction" points to x_{HI} .

The brightness temperature probes different physics at different times.

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

$z = 1100$

$z \approx 50$

$z \approx 8$

$z < 6$

The brightness temperature probes different physics at different times.

Dark Ages

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

$z = 1100$

$z \approx 50$

$z \approx 8$

$z < 6$

The brightness temperature probes different physics at different times.

Dark Ages

First Black Holes

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

First Stars

$z = 1100$

$z \approx 50$

$z \approx 8$

$z < 6$

The brightness temperature probes different physics at different times.

Dark Ages

First Black Holes

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

First Stars

The Epoch of Reionization

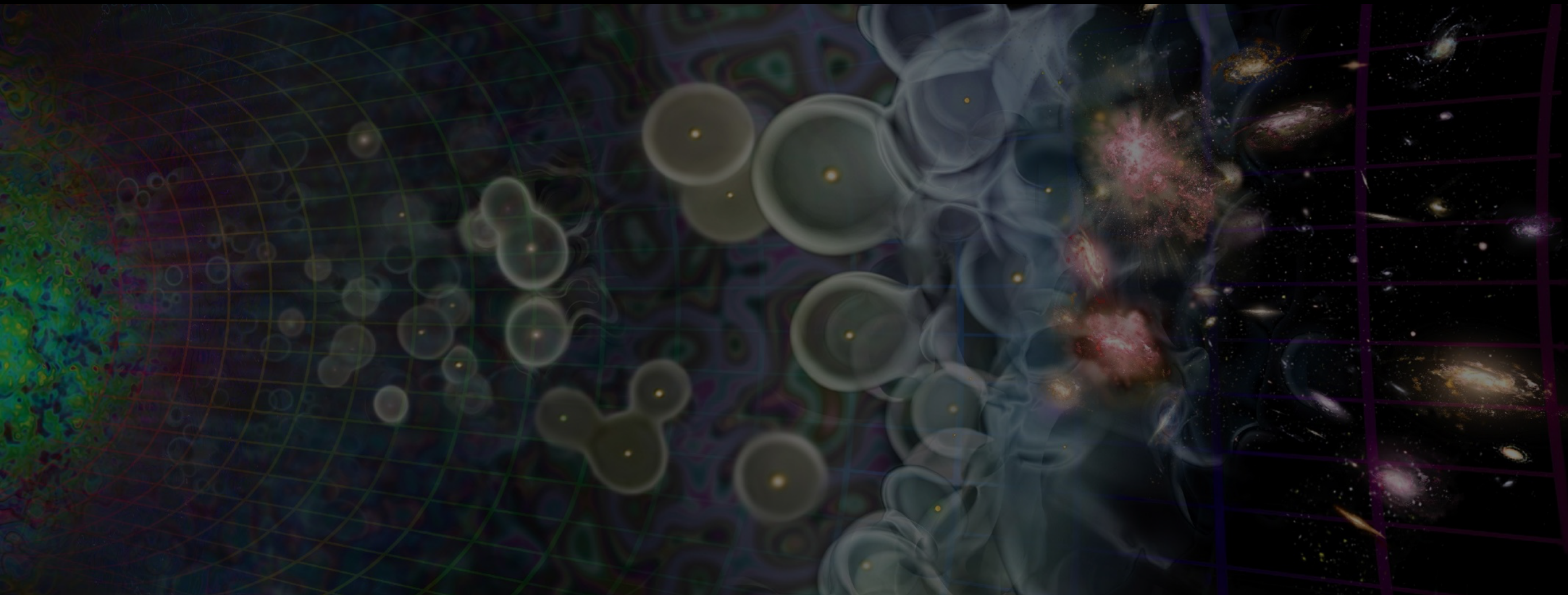
$z = 1100$

$z \approx 50$

$z \approx 8$

$z < 6$

**There's still a lot of open
astrophysical questions.**



There's still a lot of open astrophysical questions.

- What did the first stars look like? How and when how did they form?



There's still a lot of open astrophysical questions.

- What did the first stars look like? How and when how did they form?
- How did they die and were they the LIGO black hole progenitors? Or the seeds of supermassive black holes?

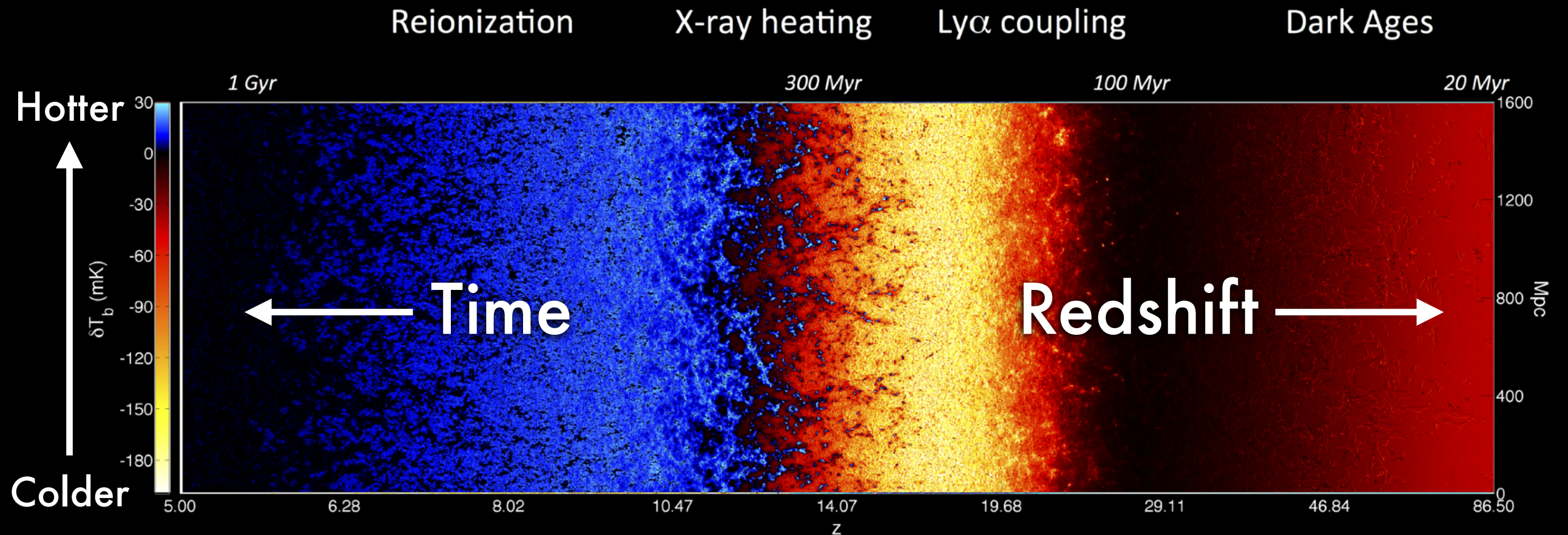
There's still a lot of open astrophysical questions.

- What did the first stars look like? How and when how did they form?
- How did they die and were they the LIGO black hole progenitors? Or the seeds of supermassive black holes?
- What determined the thermal history of the intergalactic medium? Are there new physics at play?

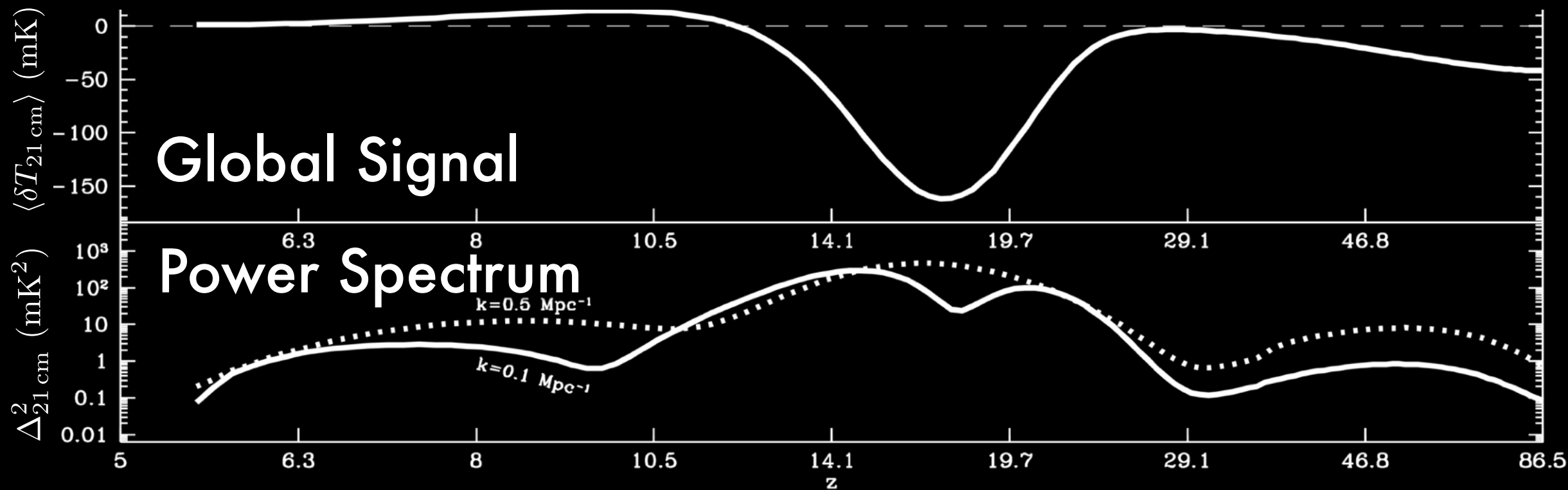
There's still a lot of open astrophysical questions.

- What did the first stars look like? How and when how did they form?
- How did they die and were they the LIGO black hole progenitors? Or the seeds of supermassive black holes?
- What determined the thermal history of the intergalactic medium? Are there new physics at play?
- What reionized the universe and when?

If we want to understand the history of 21 cm signal, we have two primary statistical probes.



Mesinger et al. (2016)

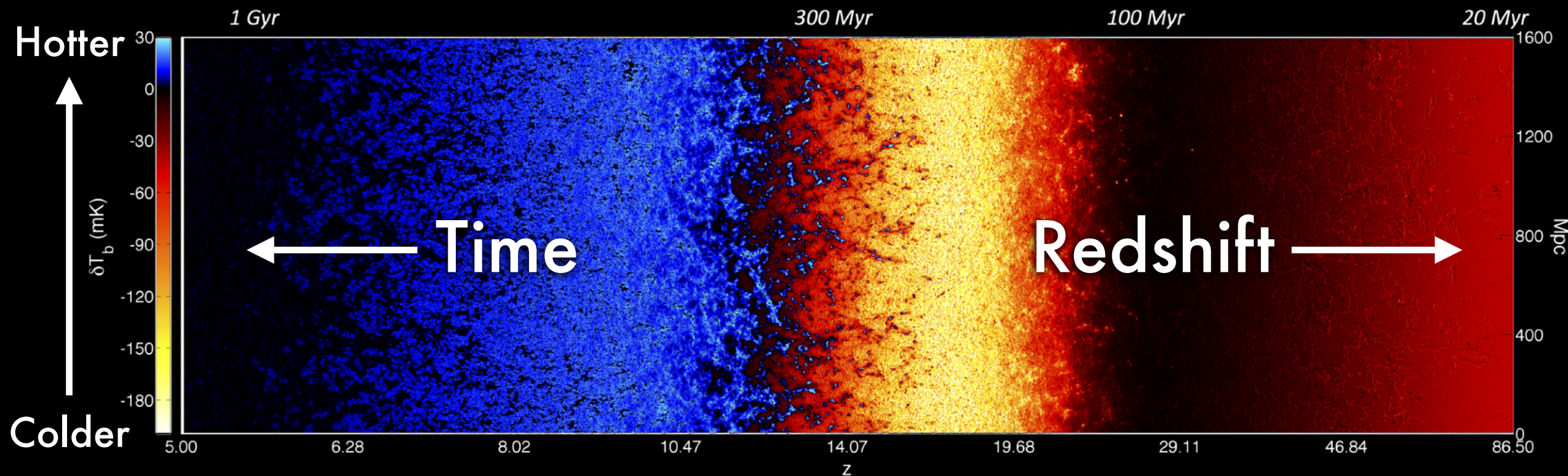


Reionization

X-ray heating

Ly α coupling

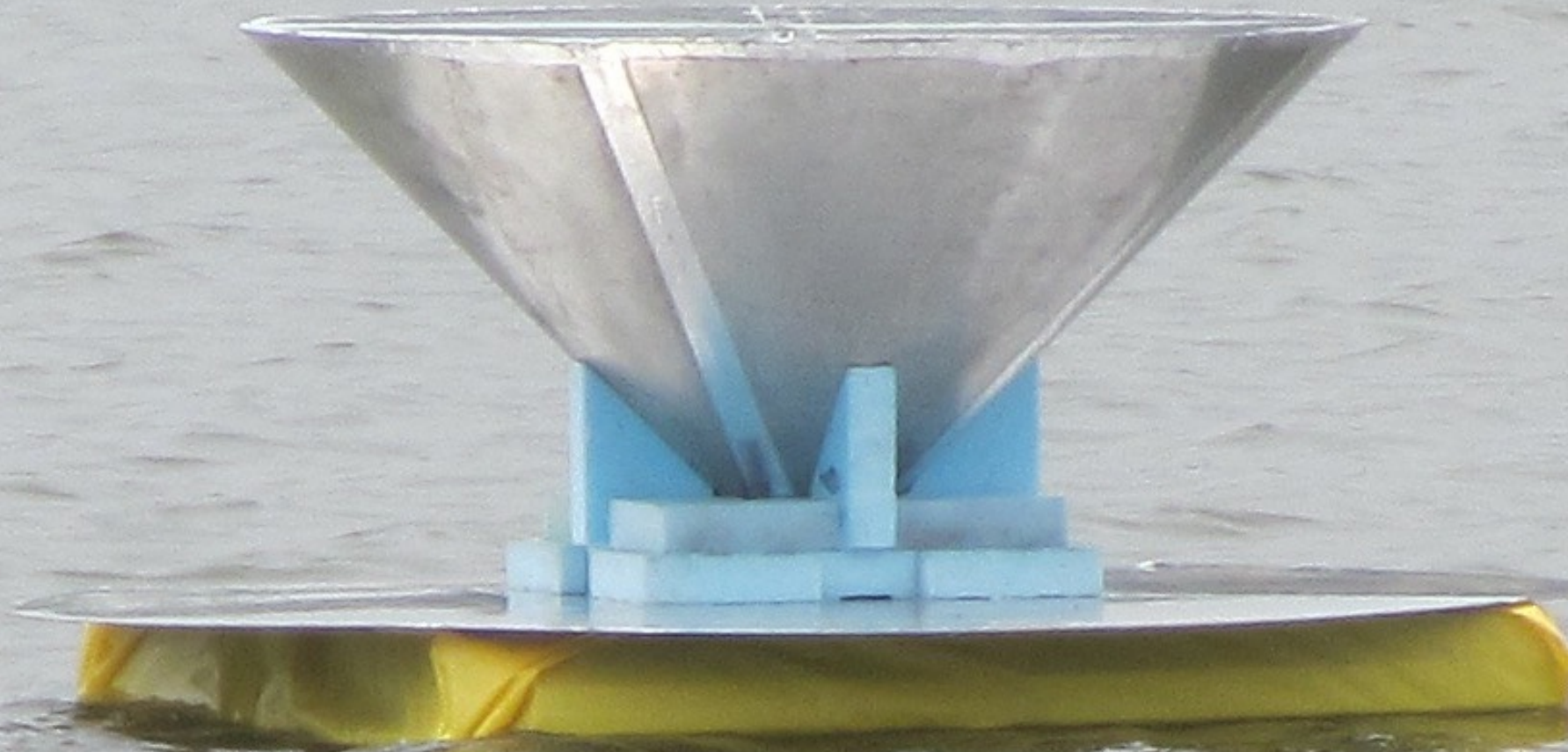
Dark Ages



EDGES

On the global signal side, there's a tension between a reported EDGES detection at $z \approx 17$ and a SARAS non-detection.

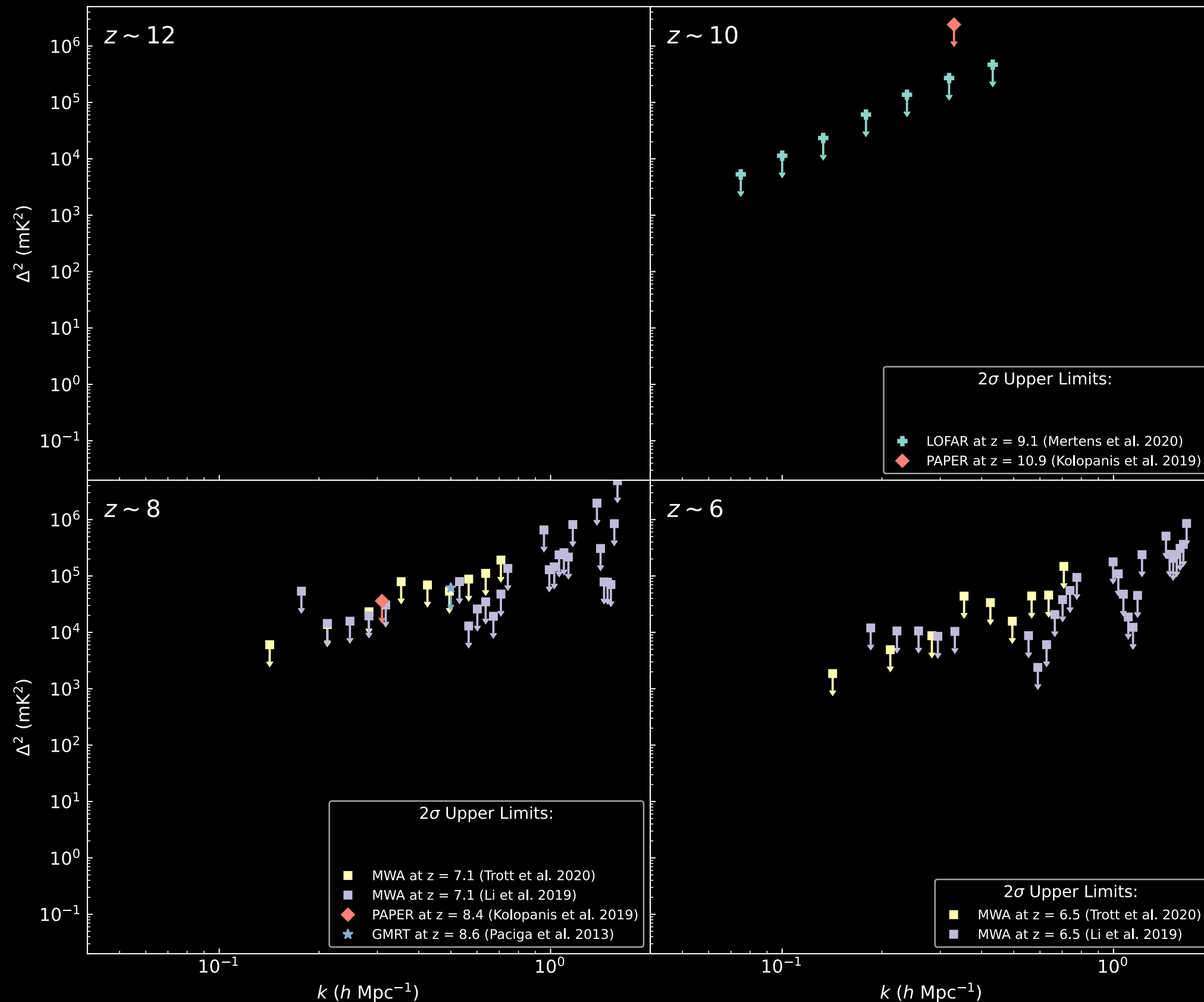
SARAS-3



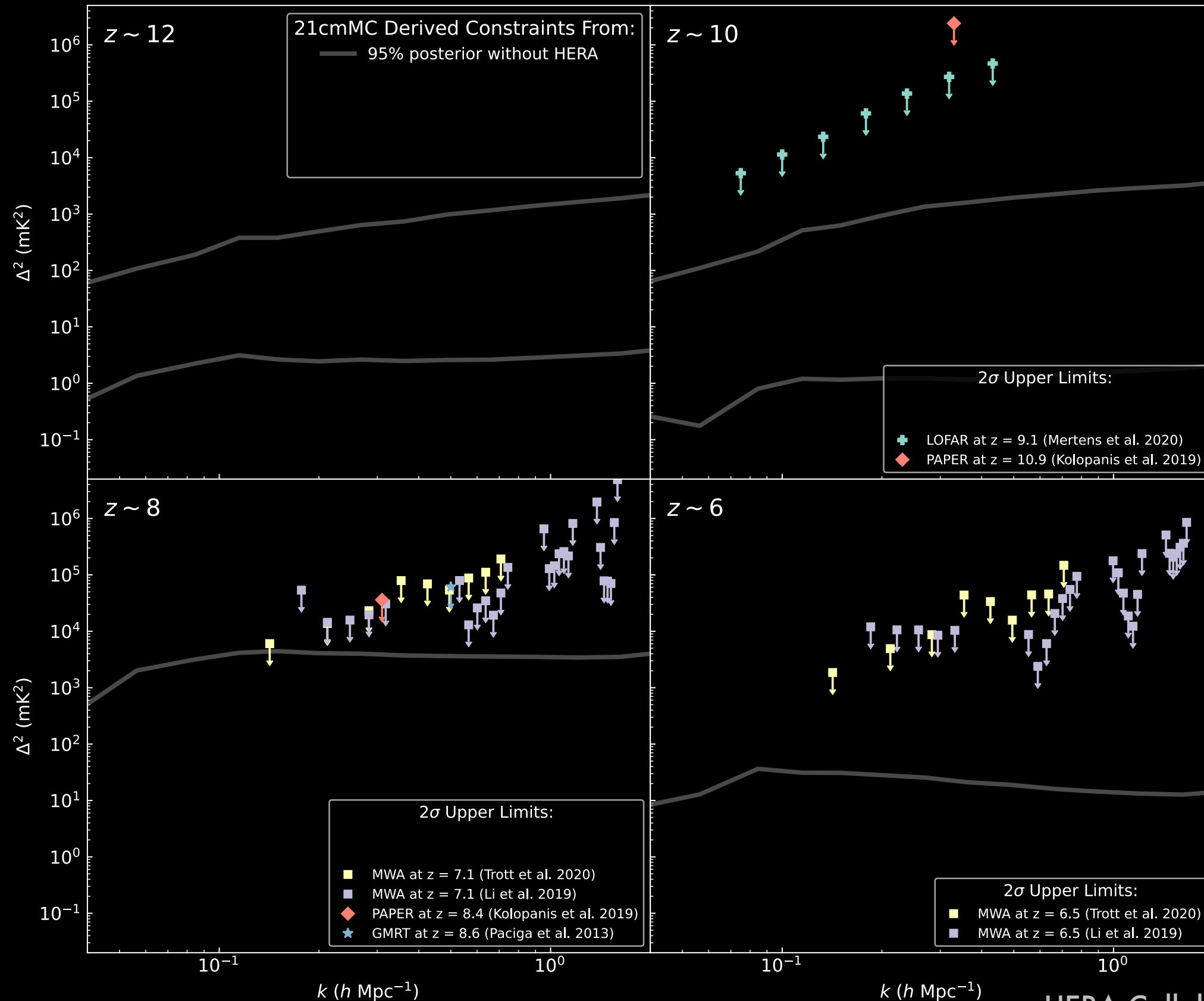
For the 21 cm power spectrum, a first generation of interferometers got us started, deploying different strategies.



Until last year, this was the state of the field:



Using 21cmMC, a wide range of power spectra were still possible, even with CMB and galaxy LF constraints.





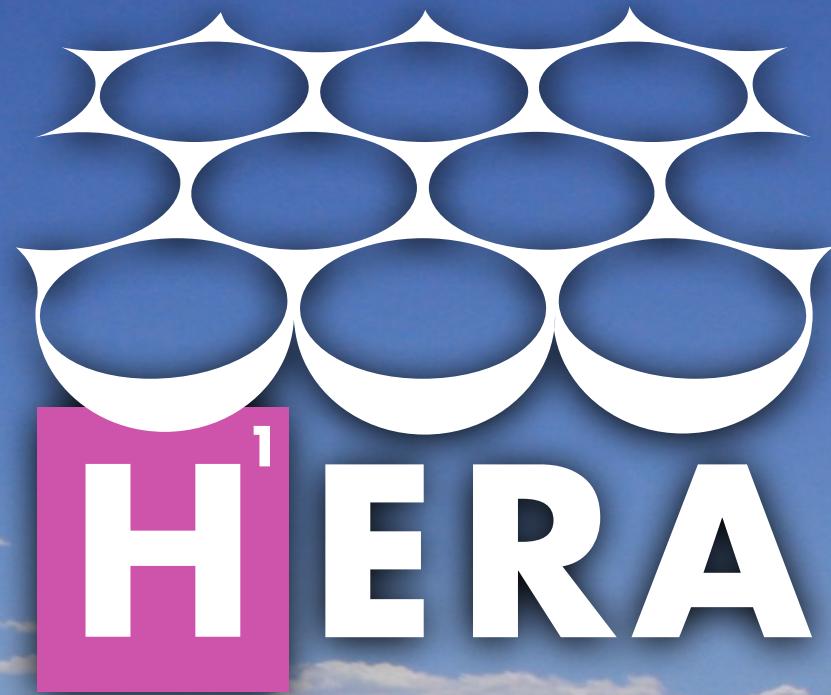
Google Earth
Data SIO, NOAA, U.S. Navy

So we went bigger...

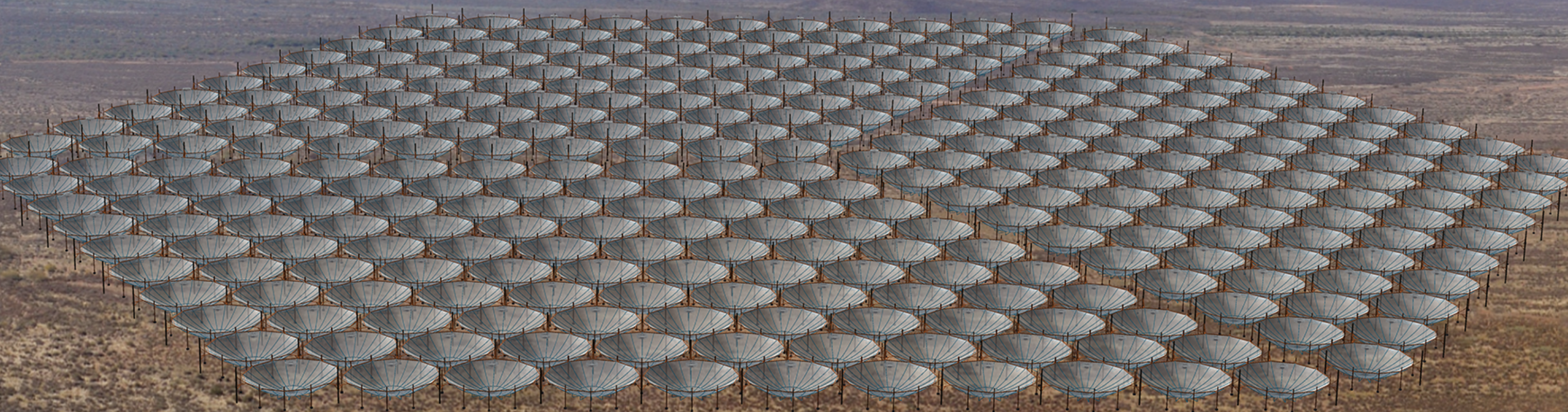


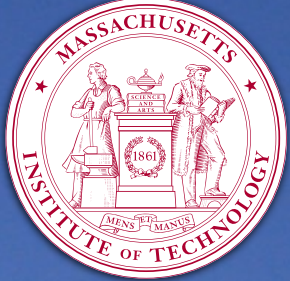
Google Earth
Data SIO, NOAA, U.S. Navy

So we went bigger...



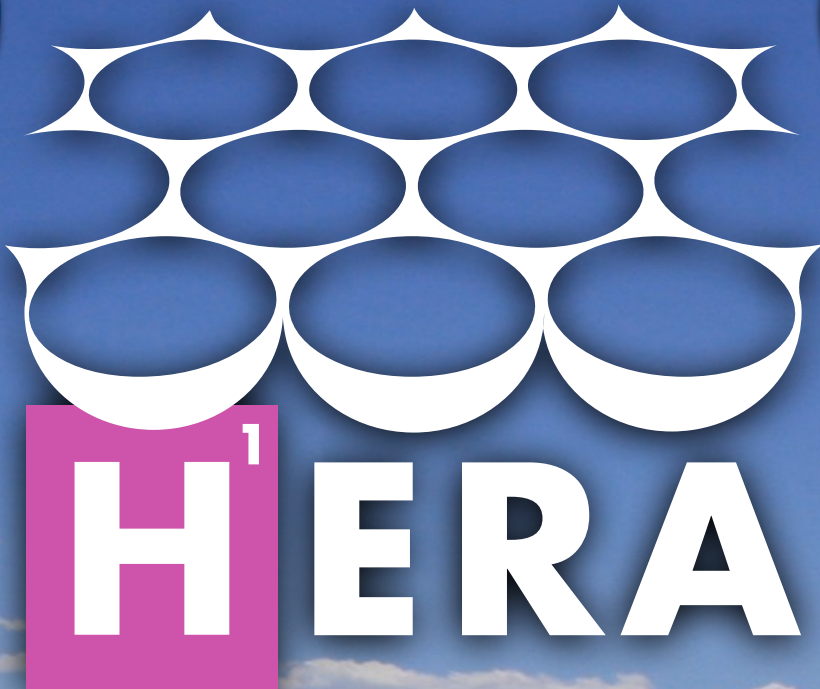
The Hydrogen Epoch of Reionization Array



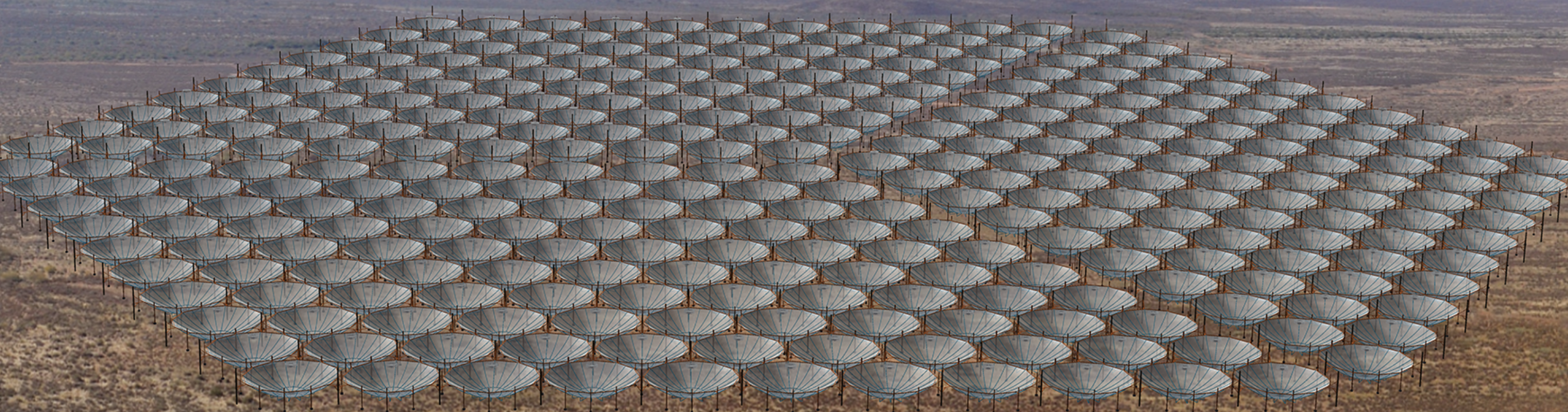


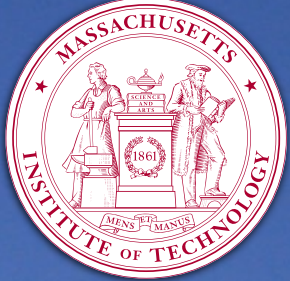
SARAO

South African Radio Astronomy Observatory



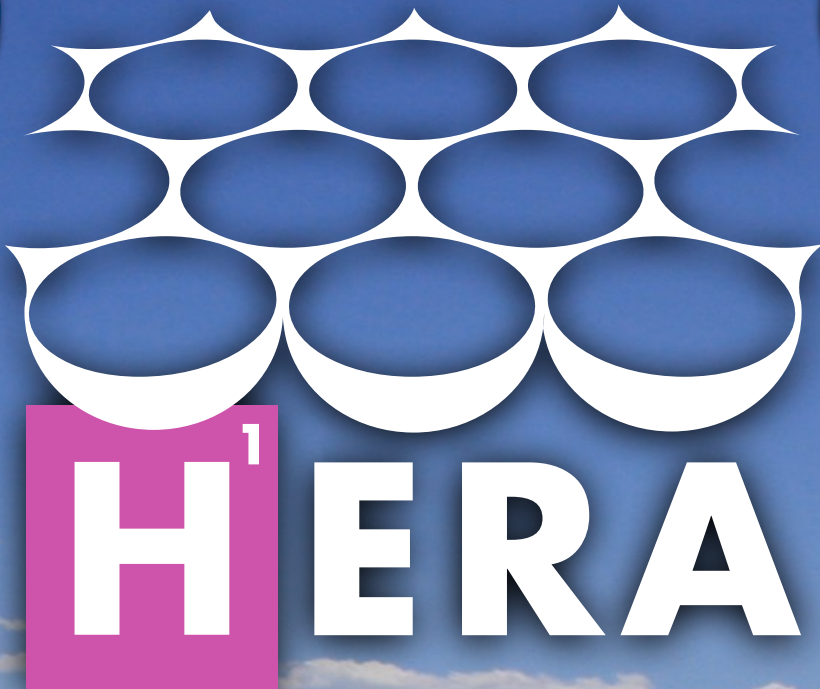
The Hydrogen Epoch of Reionization Array





SARAO

South African Radio Astronomy Observatory



BROWN



UNIVERSITY OF KWAZULU-NATAL



The Hydrogen Epoch of Reionization Array



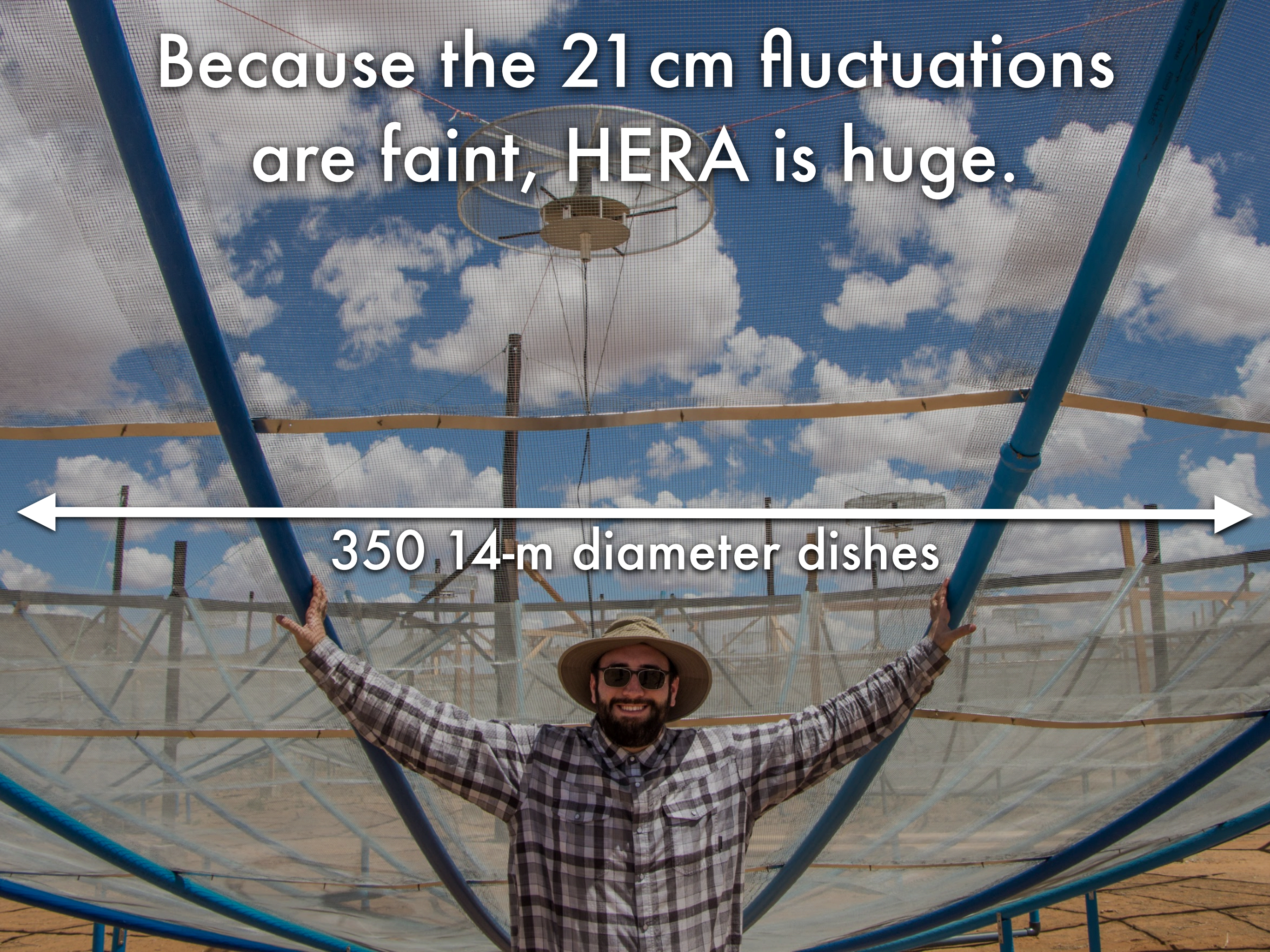
UNIVERSITY of the WESTERN CAPE

GORDON AND BETTY MOORE FOUNDATION



Because the 21 cm fluctuations are faint, HERA is huge.

← 350 14-m diameter dishes →





The HERA Stripe

HERA is a drift scan instrument that maps out a stripe of constant declination.

Our biggest problem
is foregrounds.

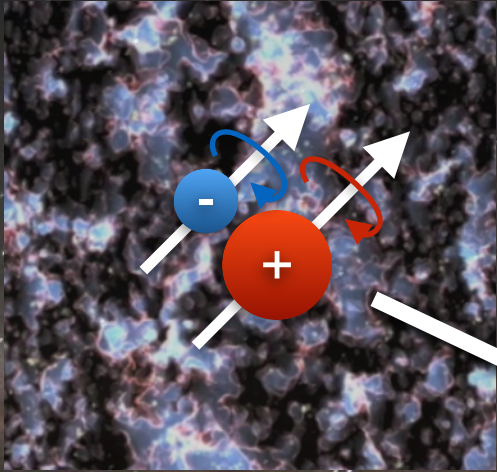
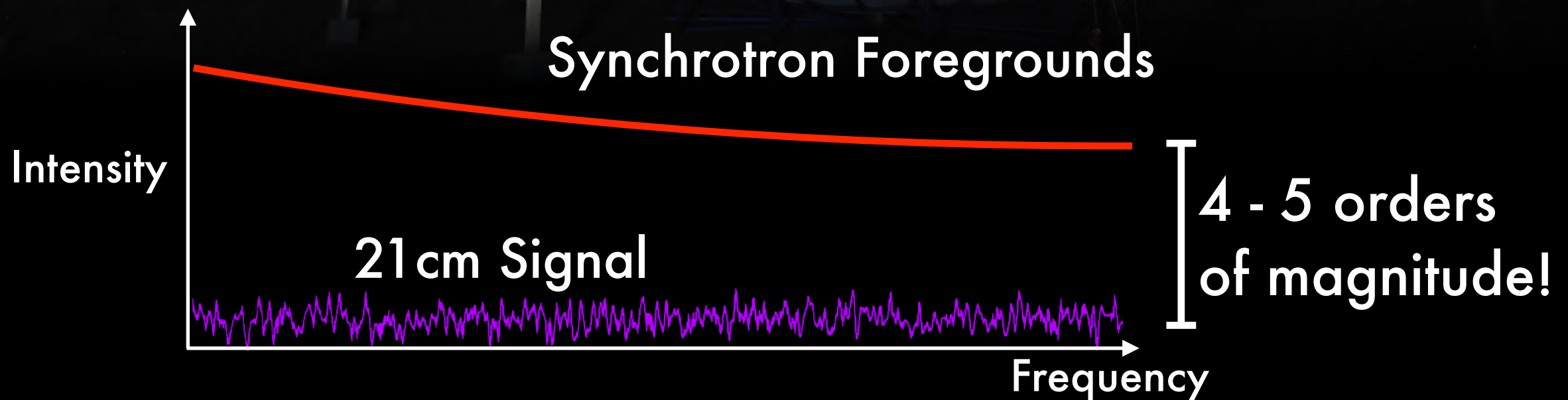
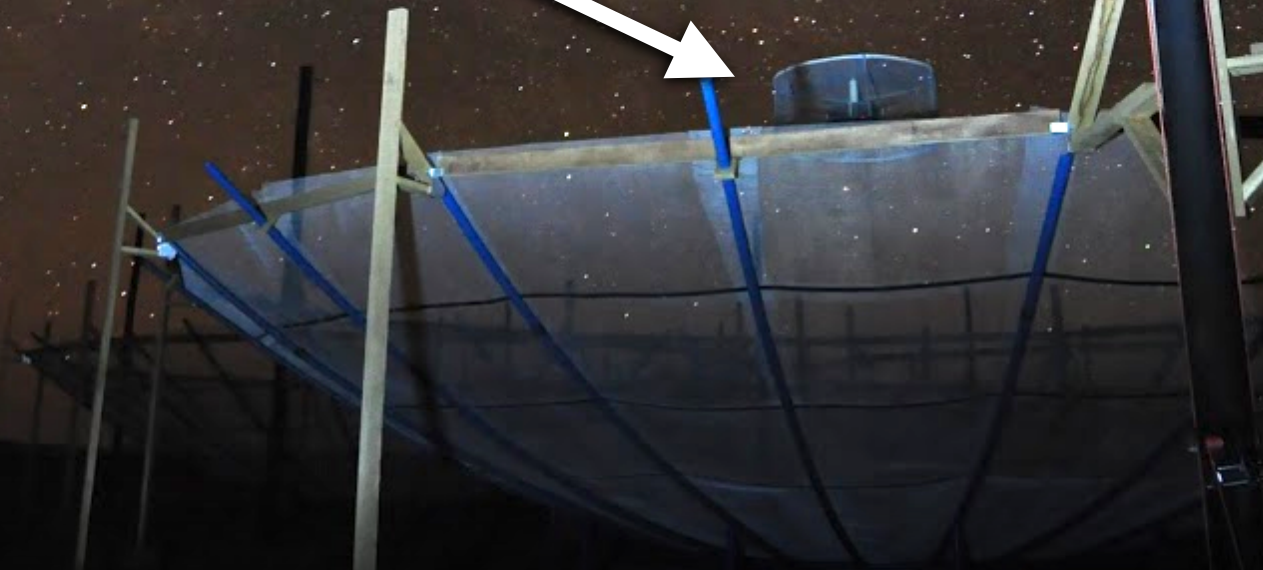
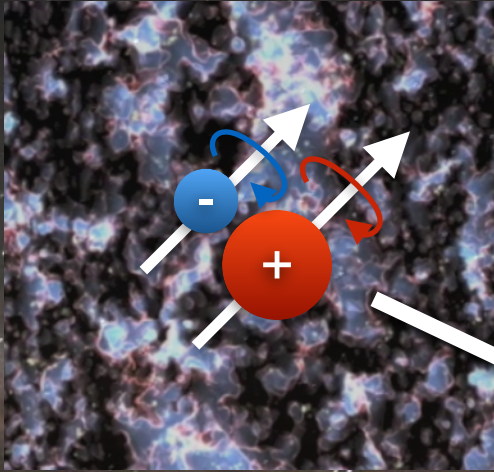
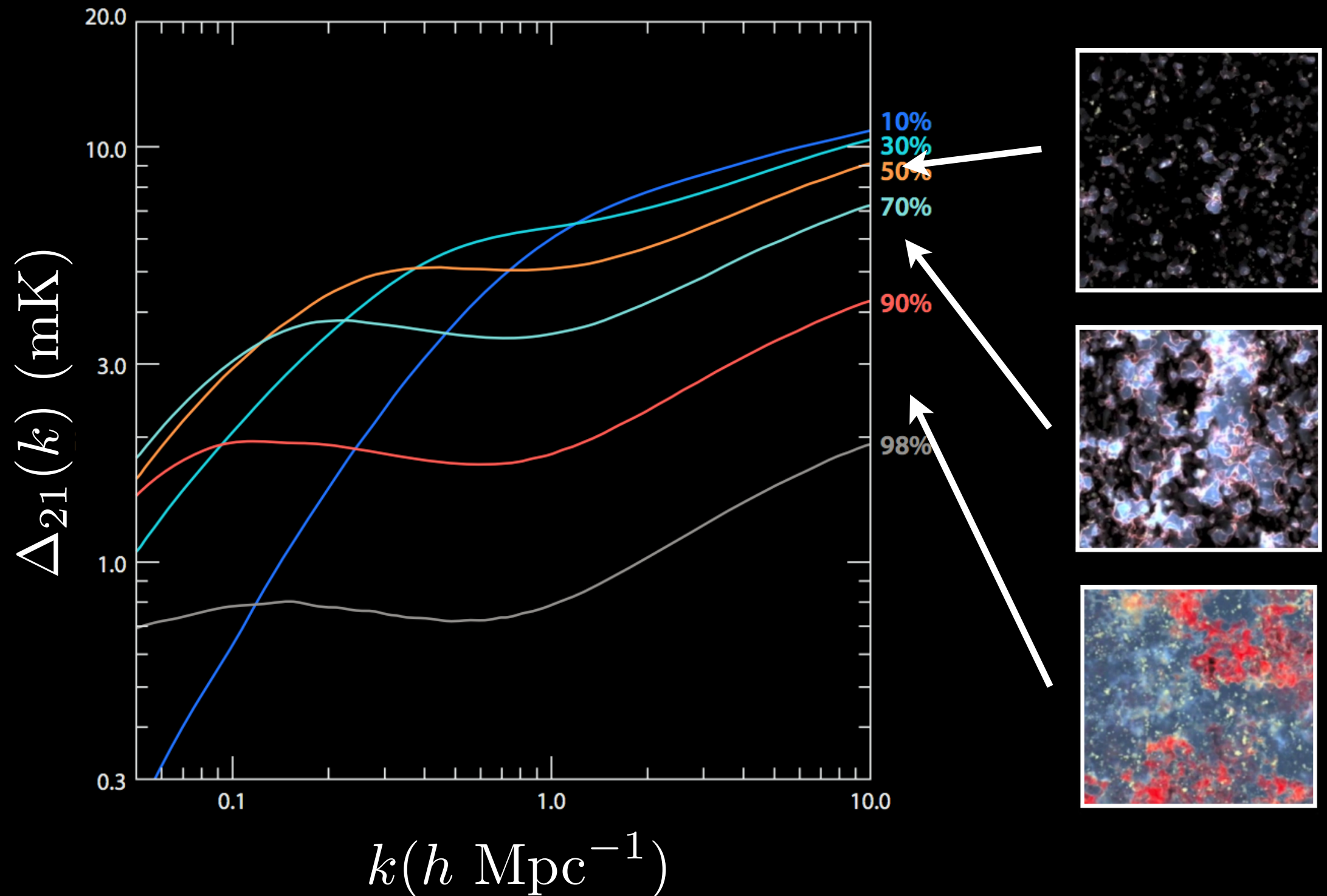


Photo: Carina Cheng

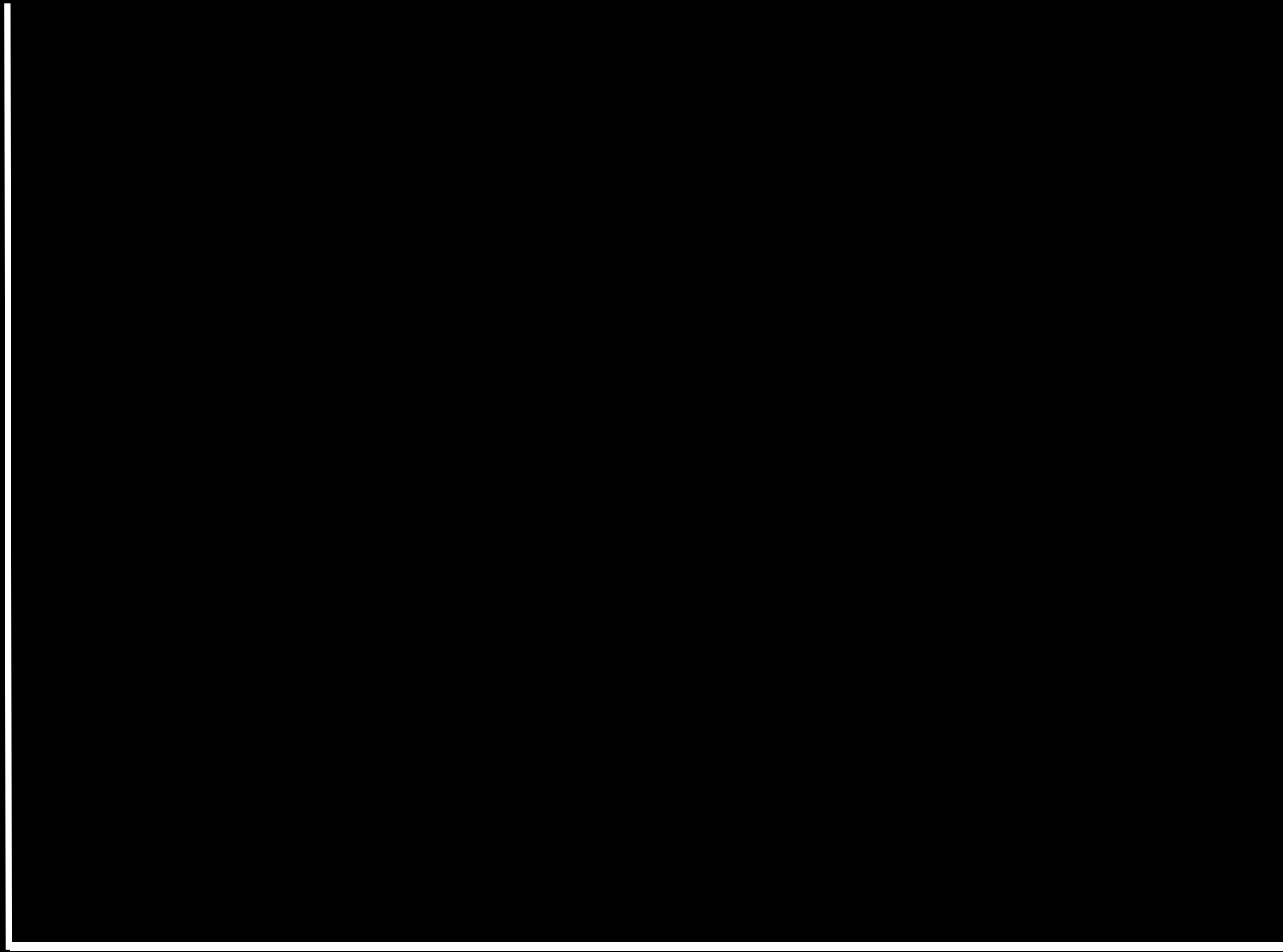
The key to separating out foregrounds is their spectral smoothness.



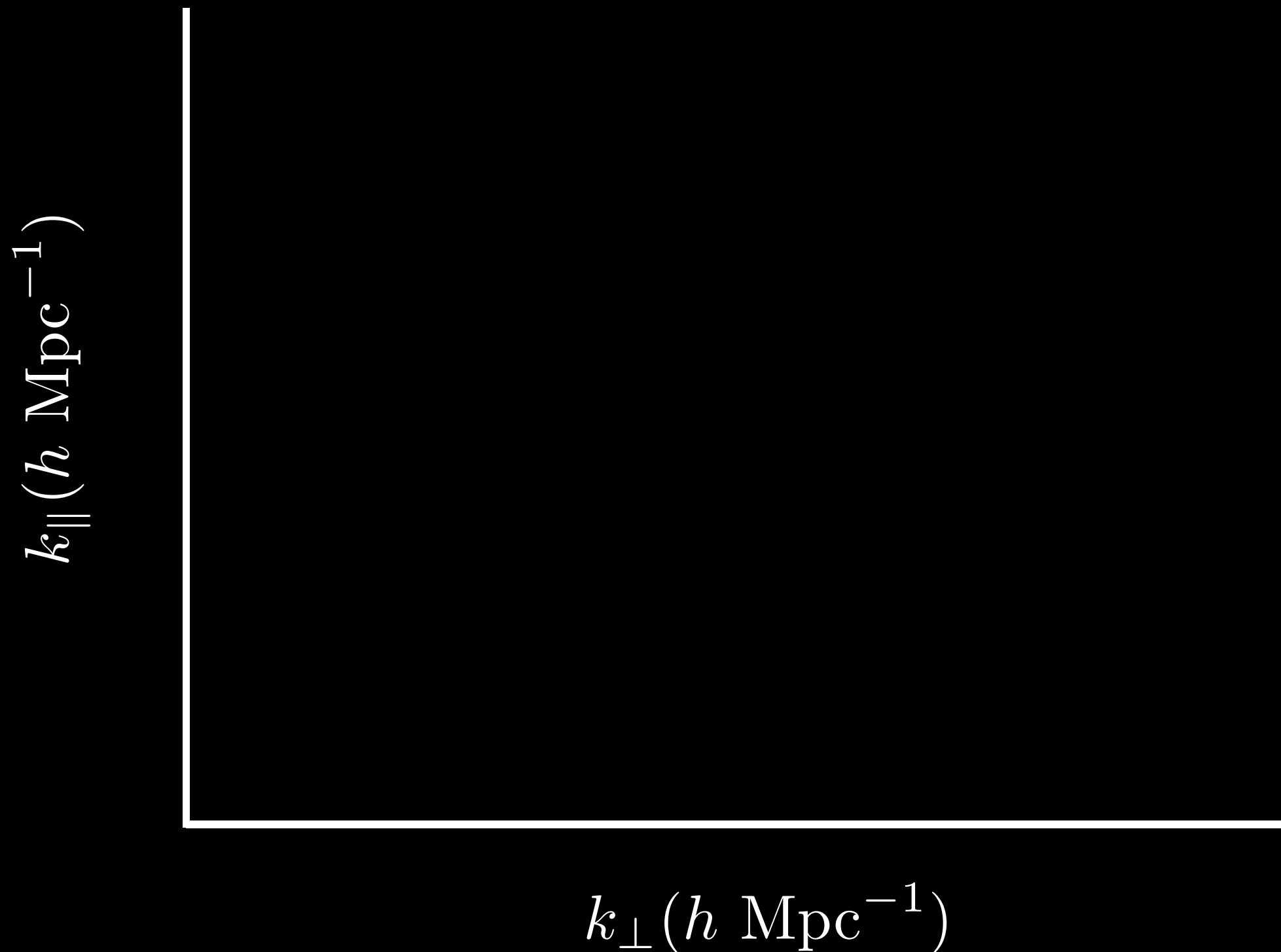
So instead of spherically averaged Fourier space...



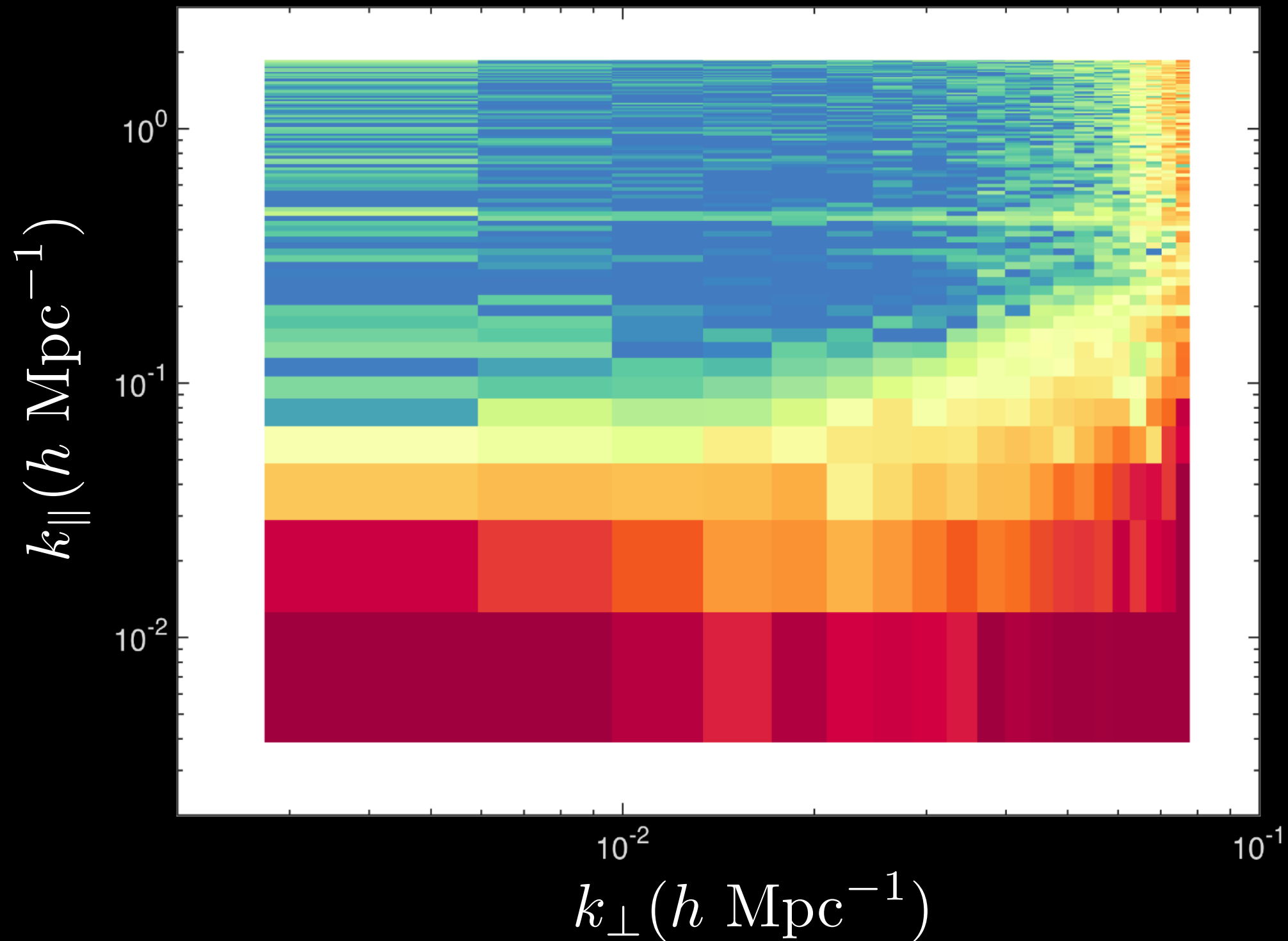
So instead of spherically averaged Fourier space...



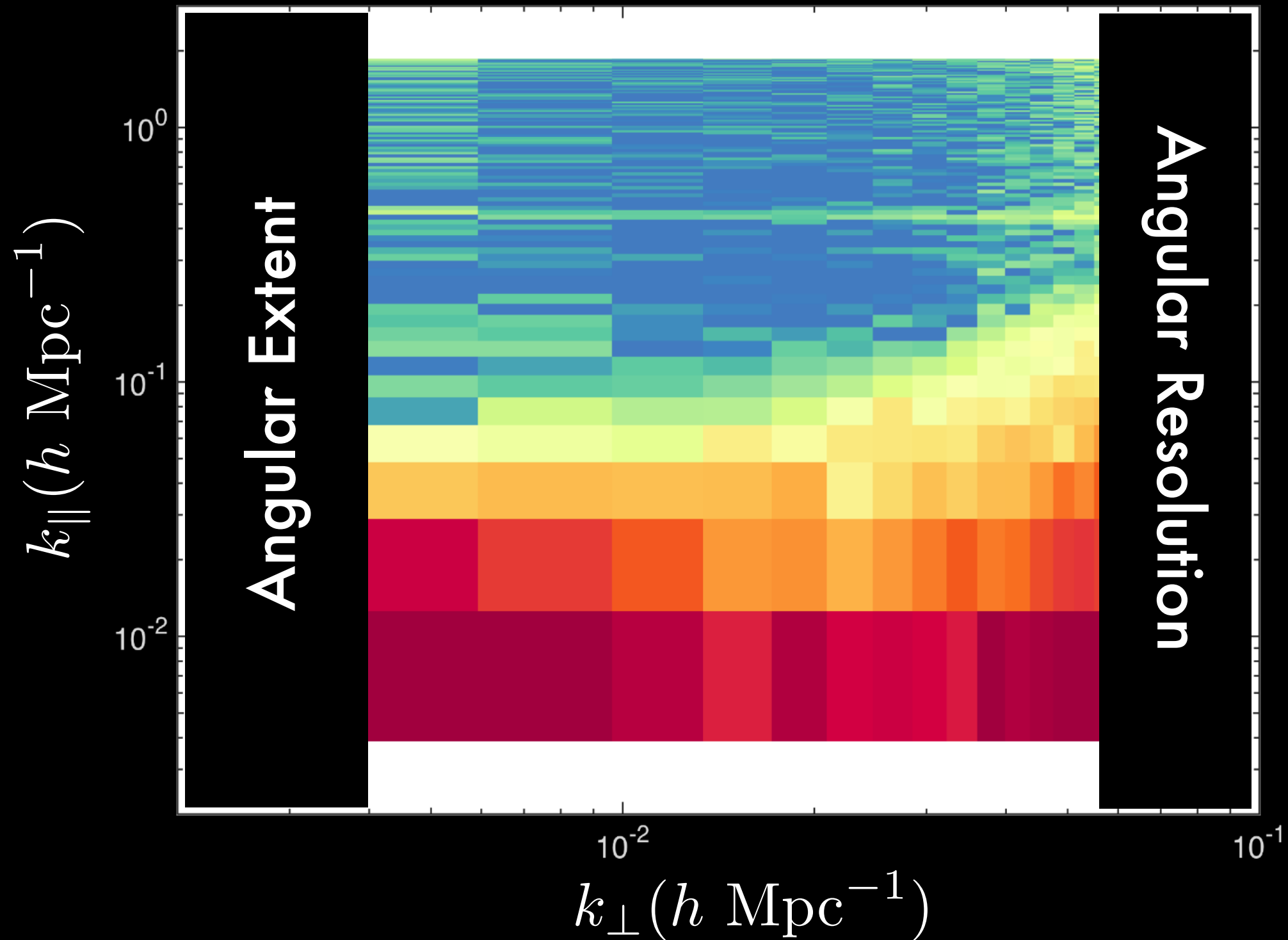
We separate out Fourier modes parallel and perpendicular to the line of sight.



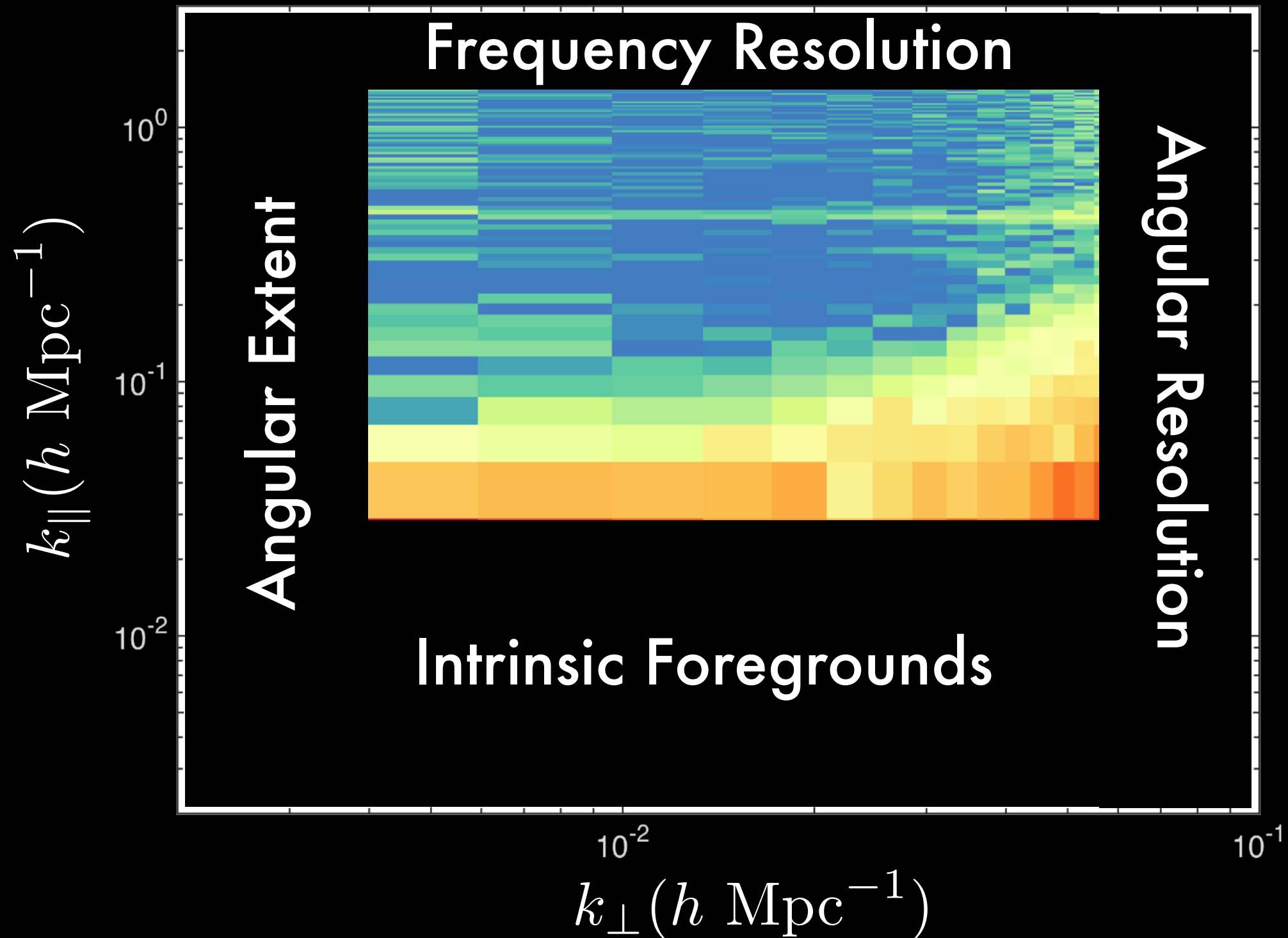
And we find a "window."



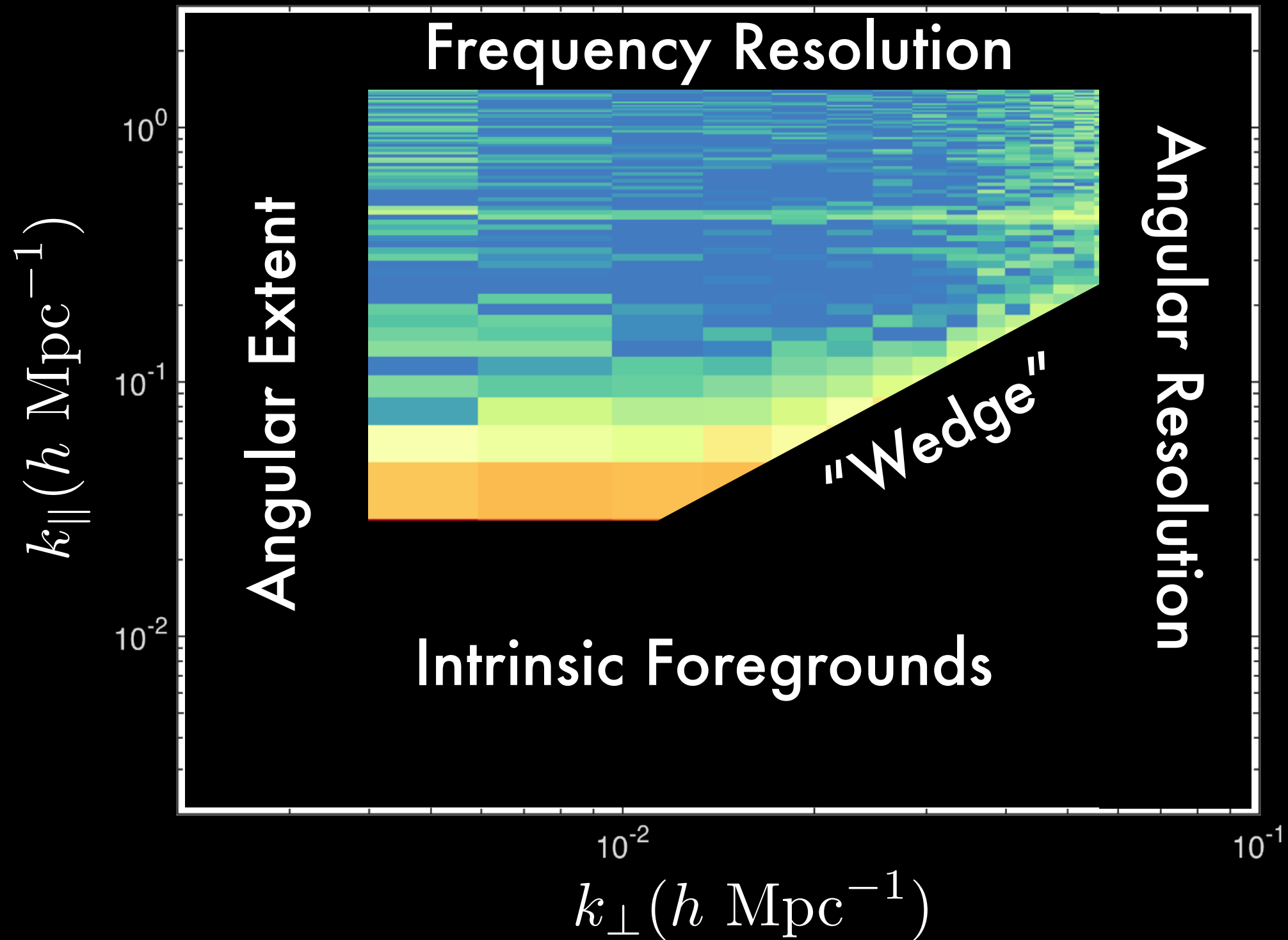
And we find a "window."



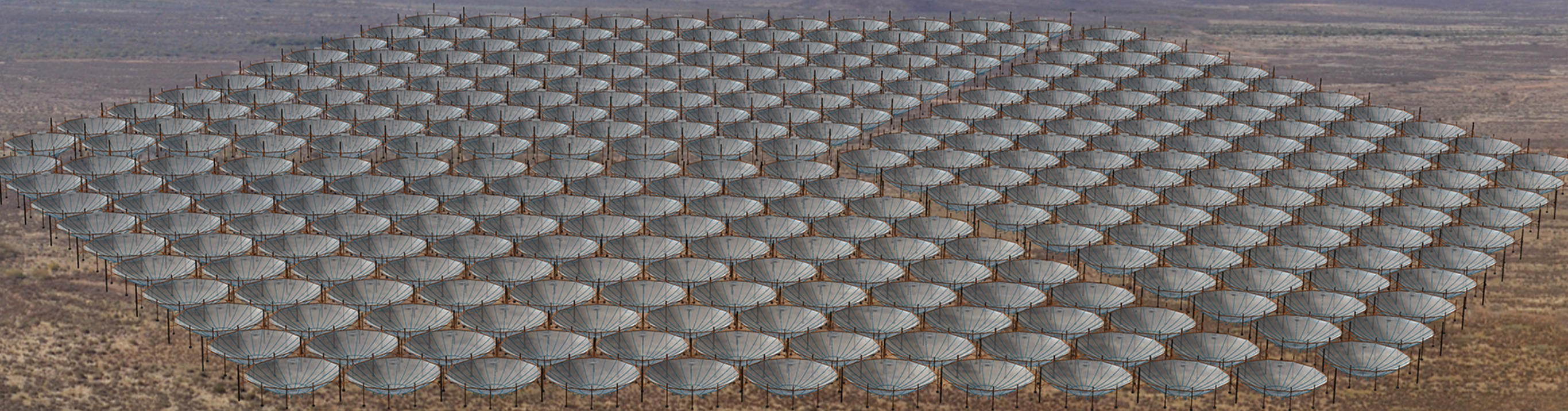
And we find a "window."

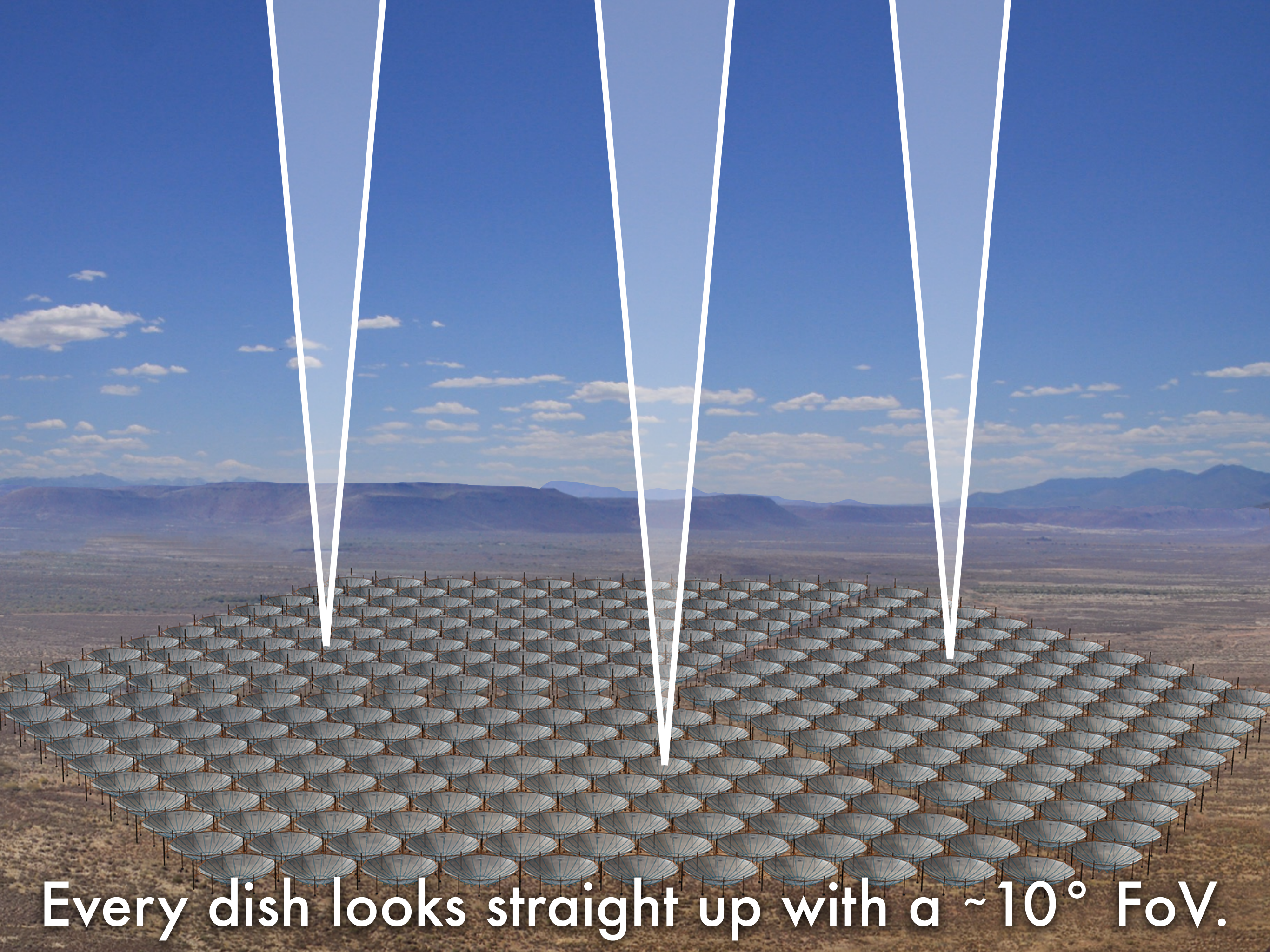


And we find a "window."



What does HERA actually measure?





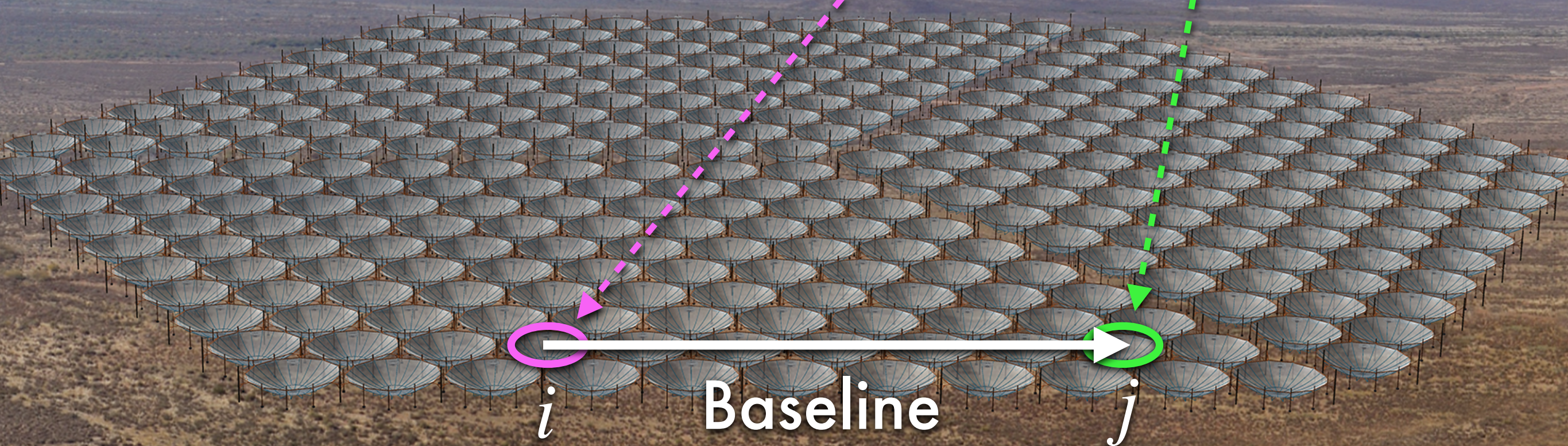
Every dish looks straight up with a $\sim 10^\circ$ FoV.

Interferometers measure Fourier modes on the sky, which we call “visibilities.”



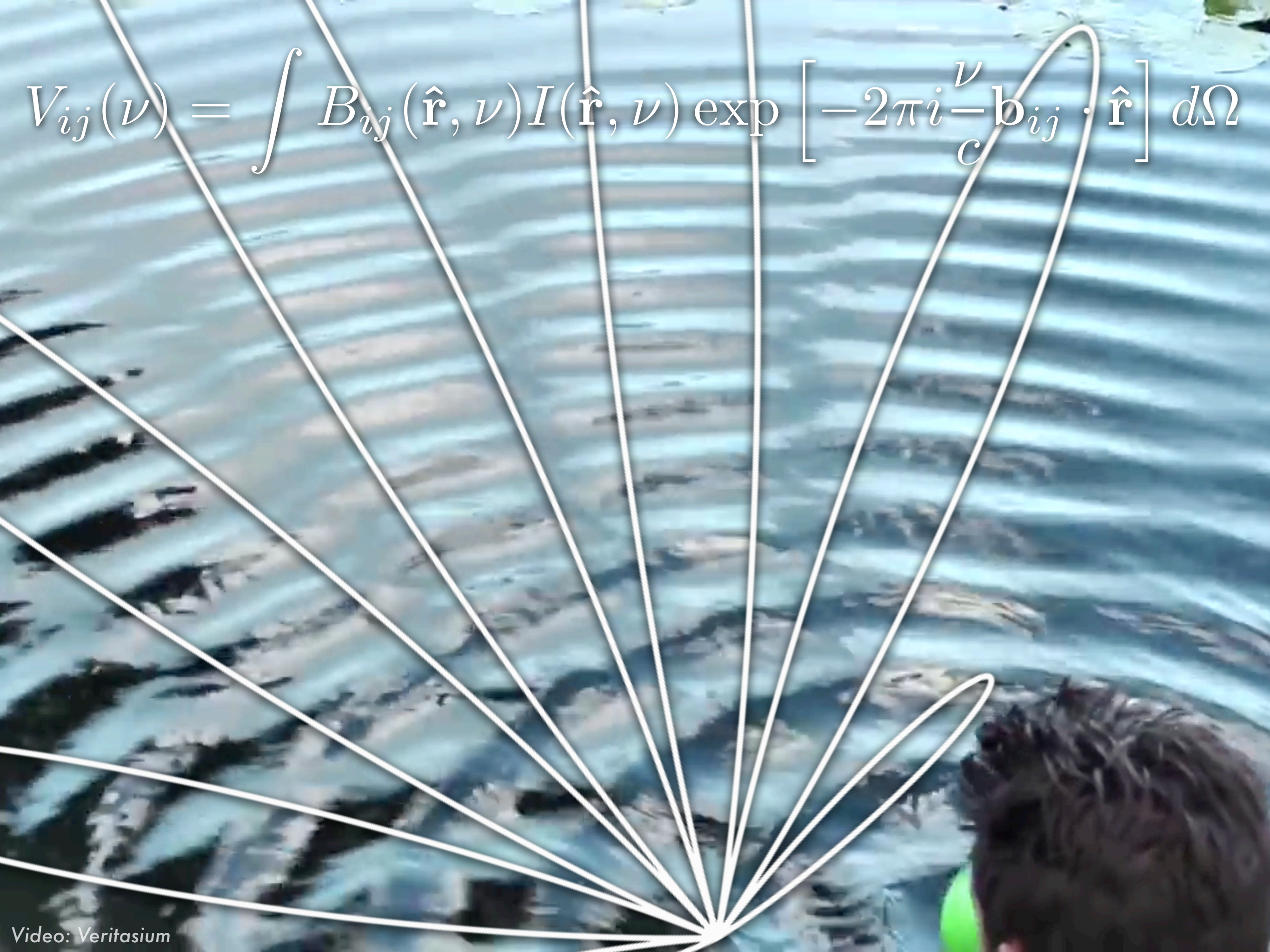
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

“Visibility” Beam Sky



$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

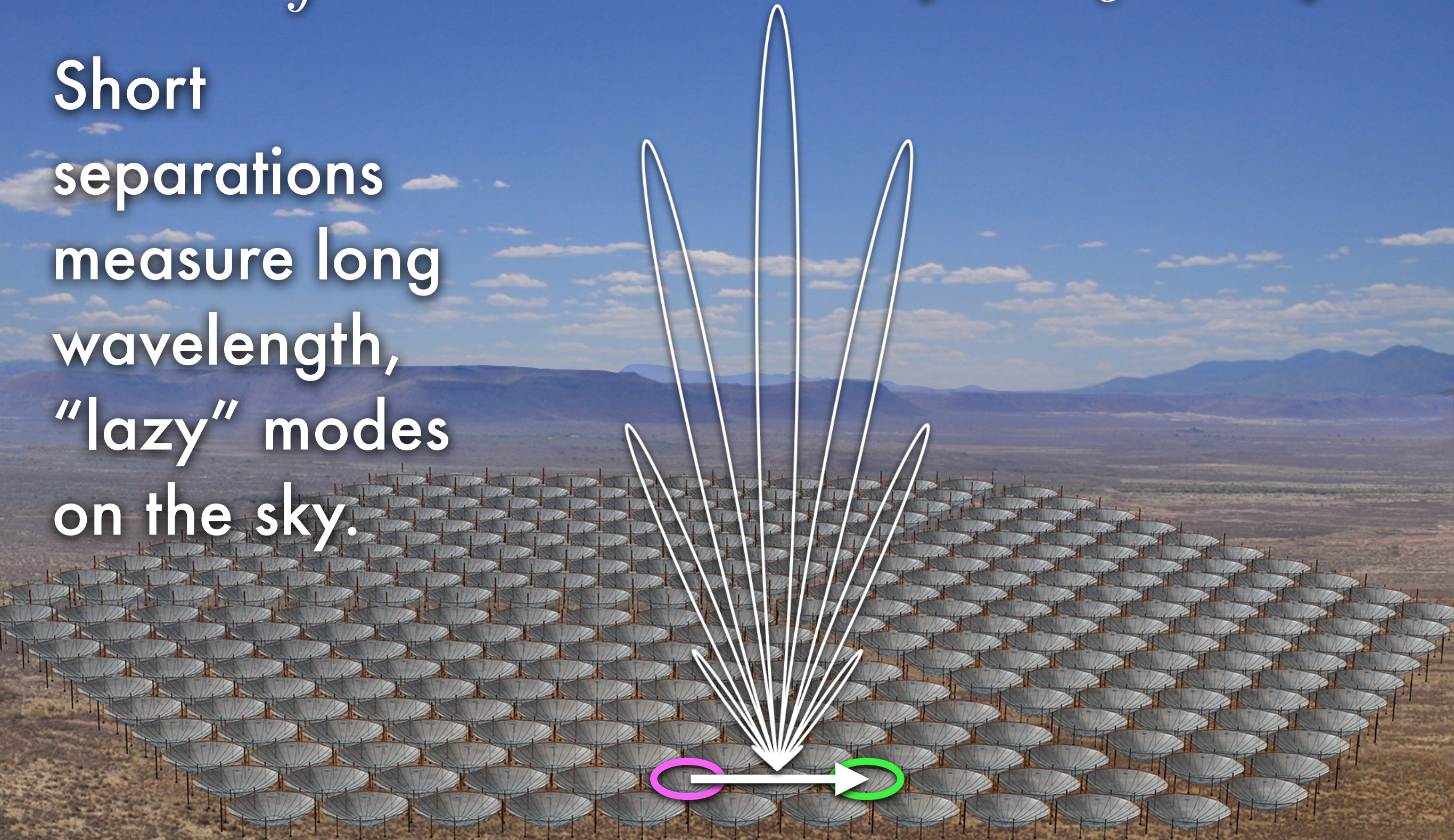
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

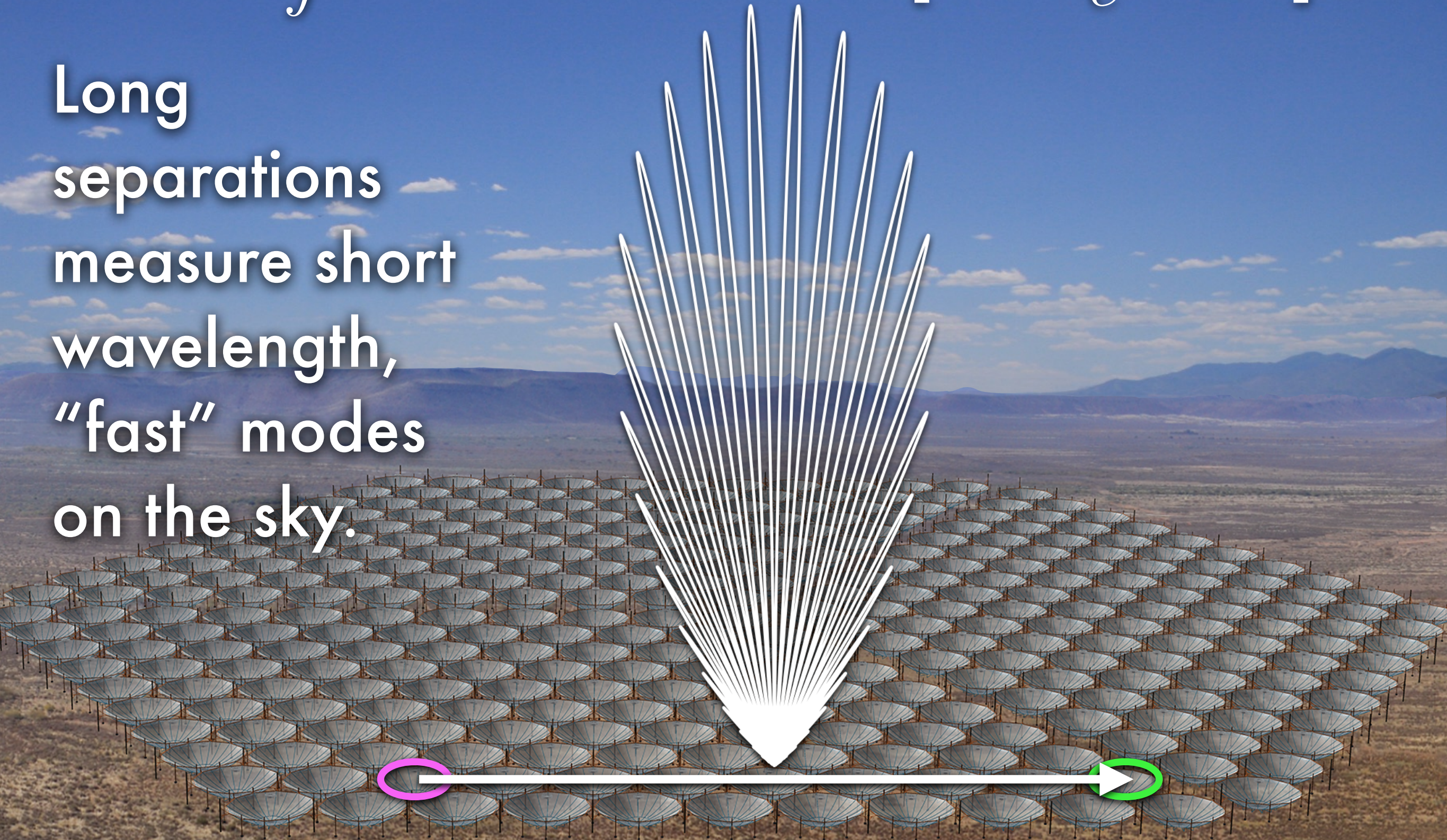
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Short
separations
measure long
wavelength,
“lazy” modes
on the sky.

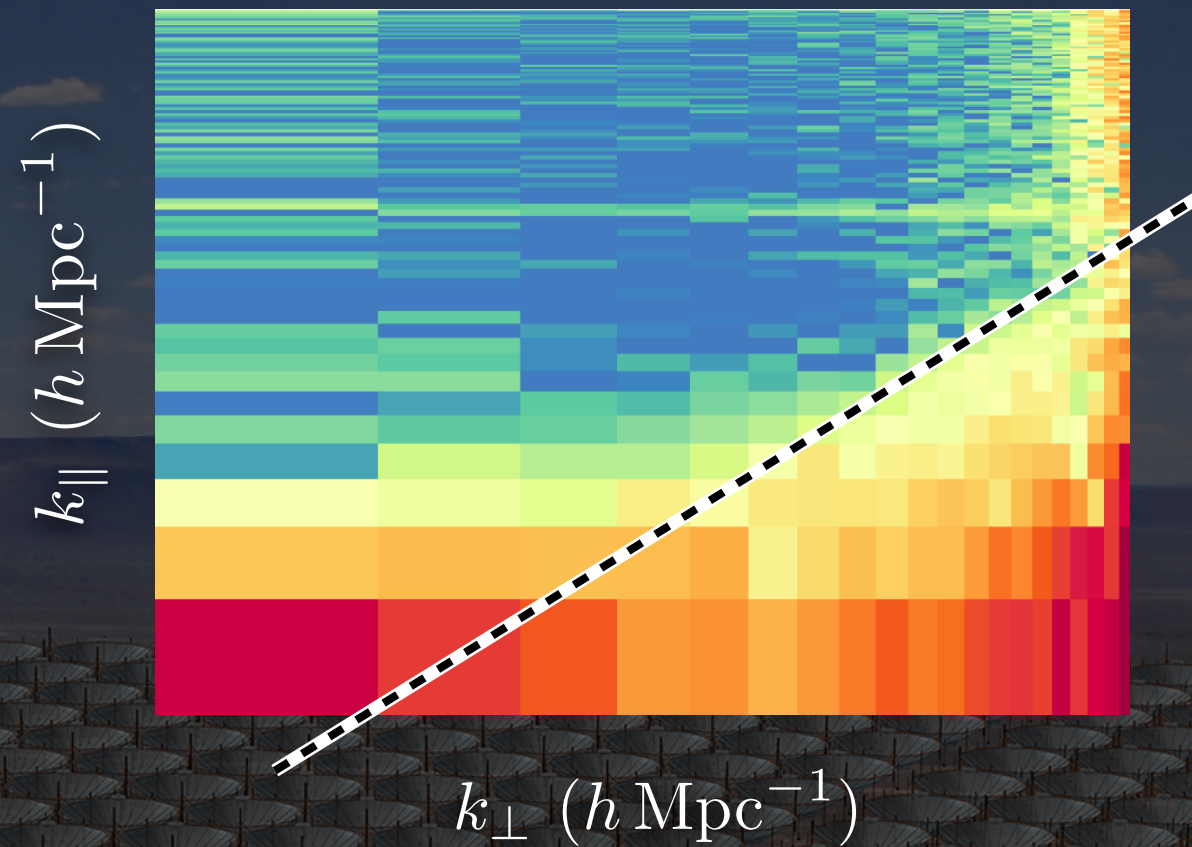


$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Long
separations
measure short
wavelength,
“fast” modes
on the sky.

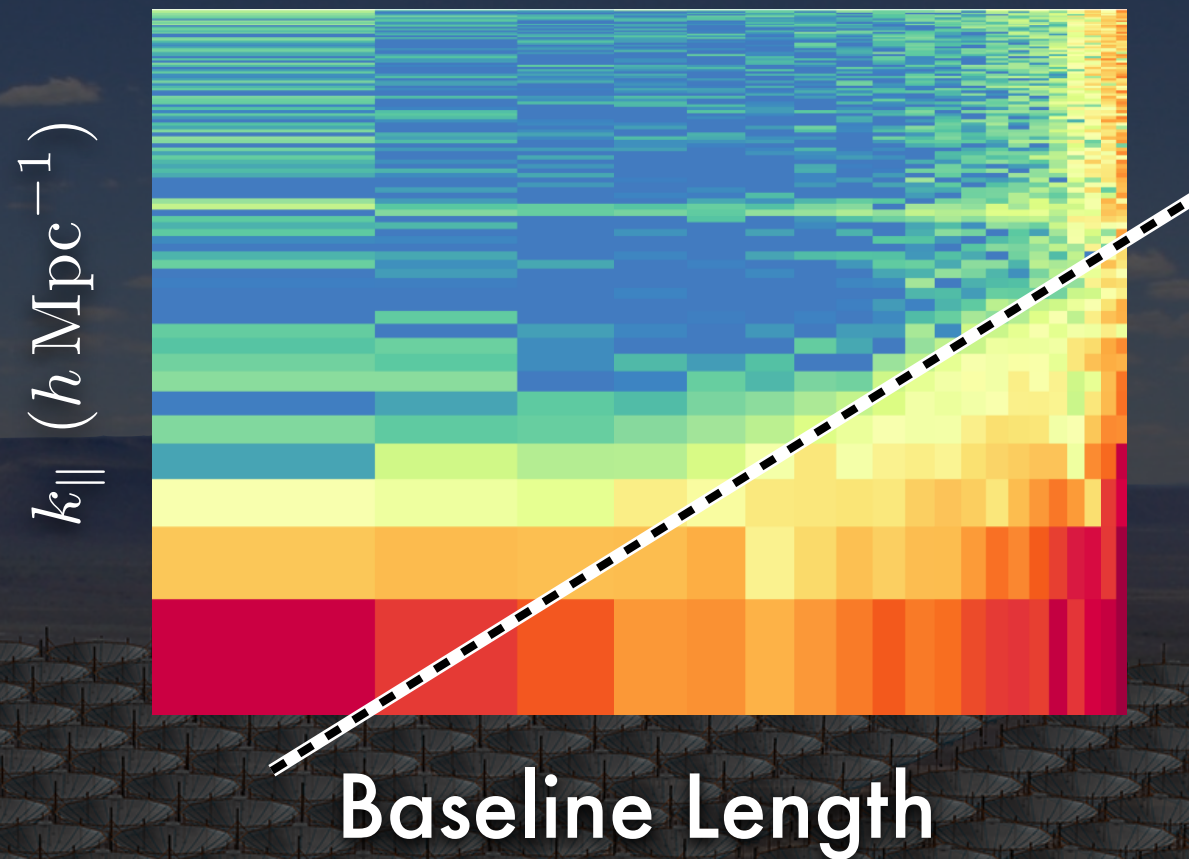


$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



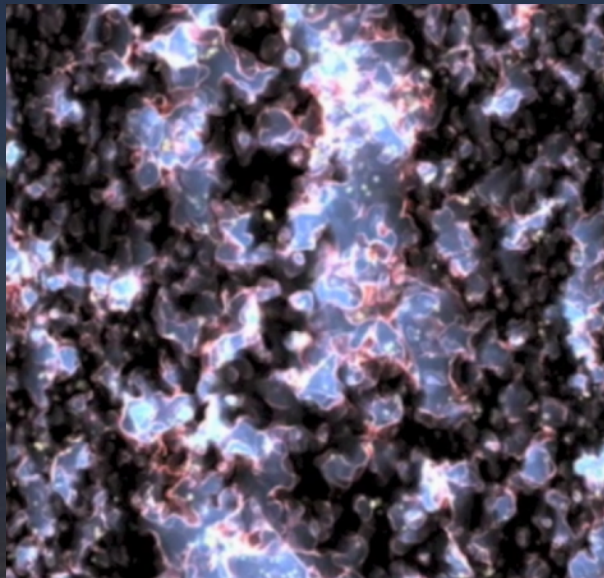
\mathbf{k}_{\perp} is effectively baseline length.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

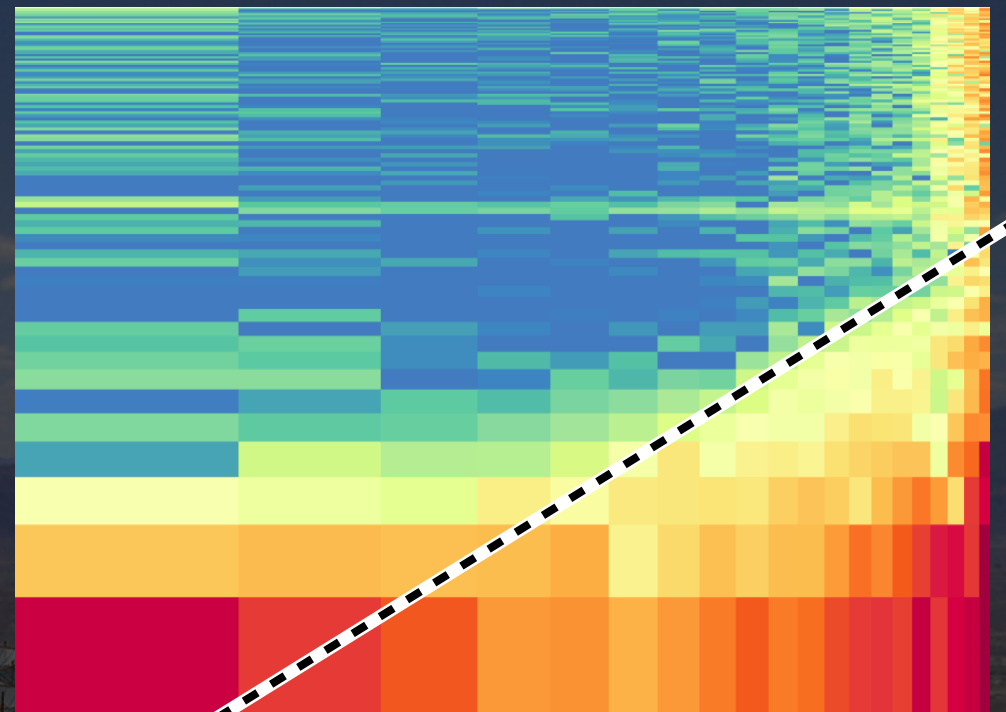


\mathbf{k}_{\perp} is effectively baseline length.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

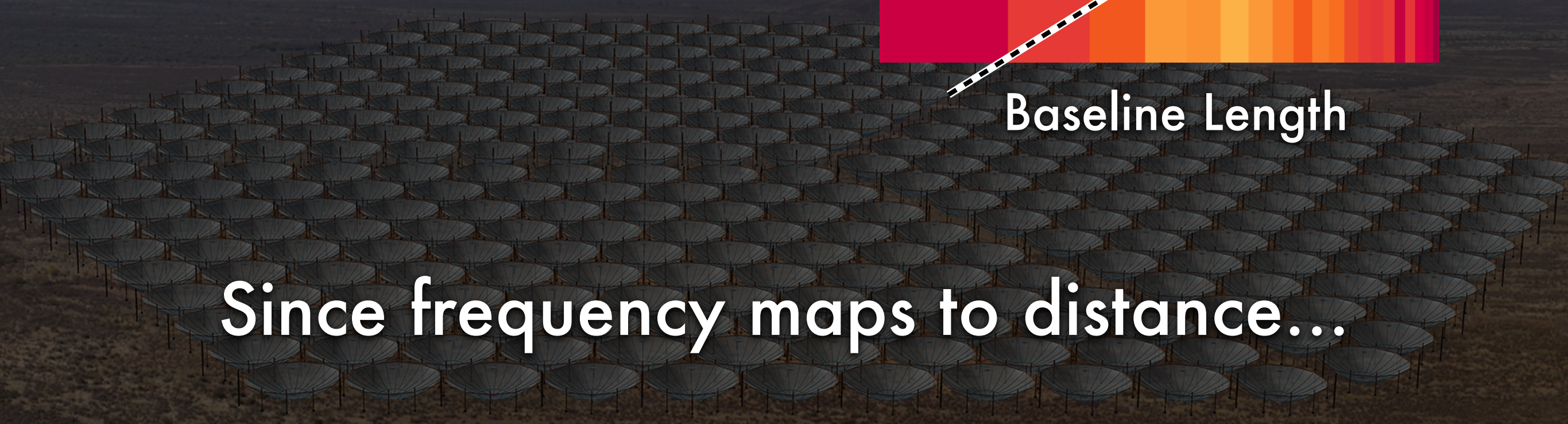


$k_{\parallel} (h \text{ Mpc}^{-1})$

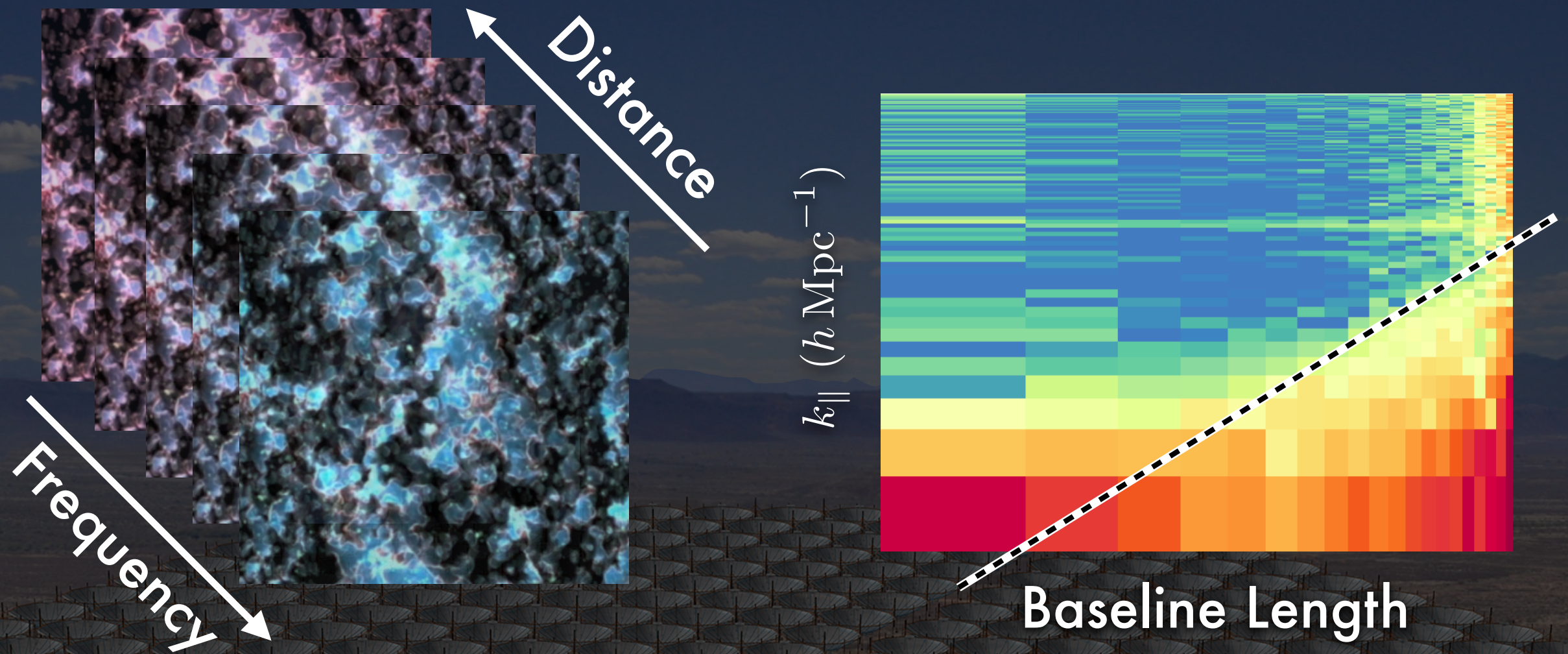


Baseline Length

Since frequency maps to distance...

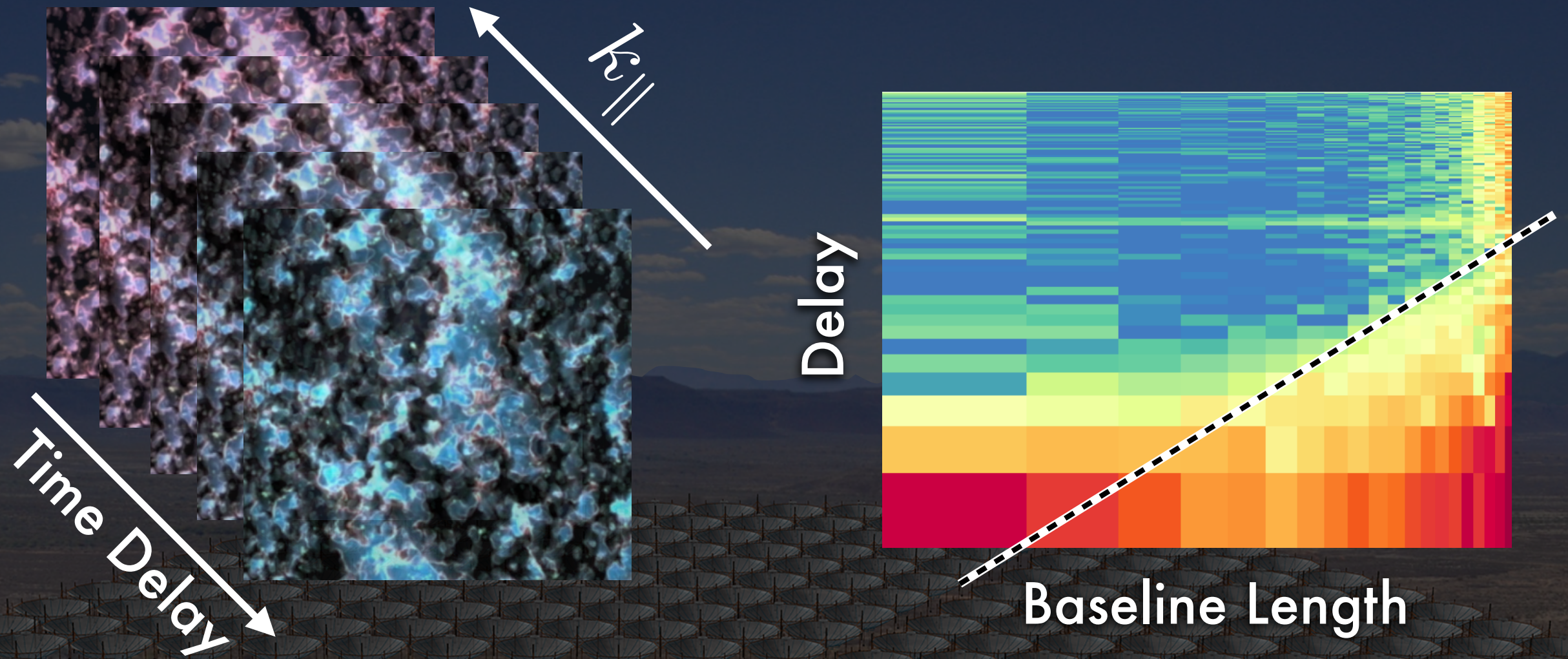


$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



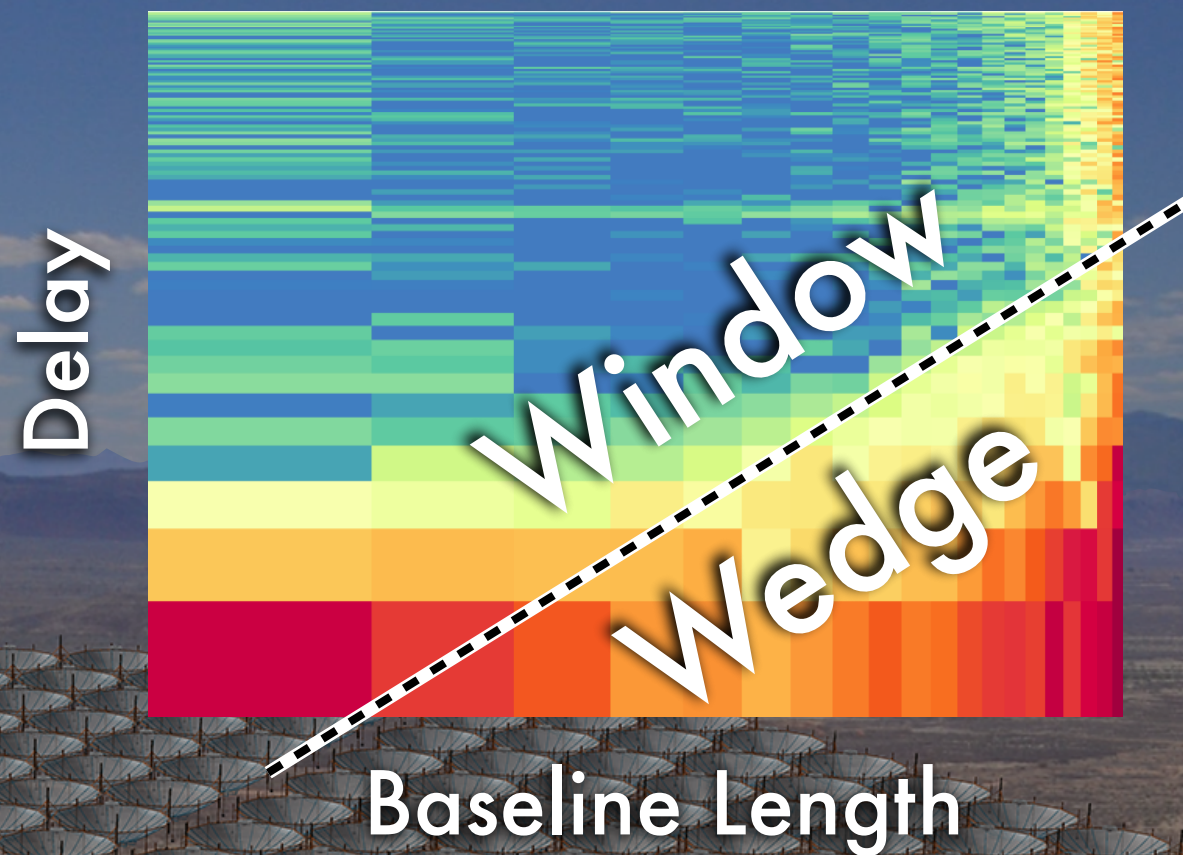
Since frequency maps to distance...

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

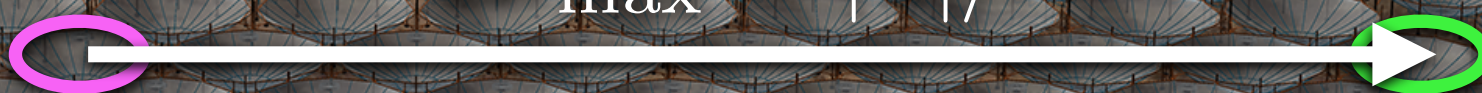


k_{\parallel} is effectively time delay.

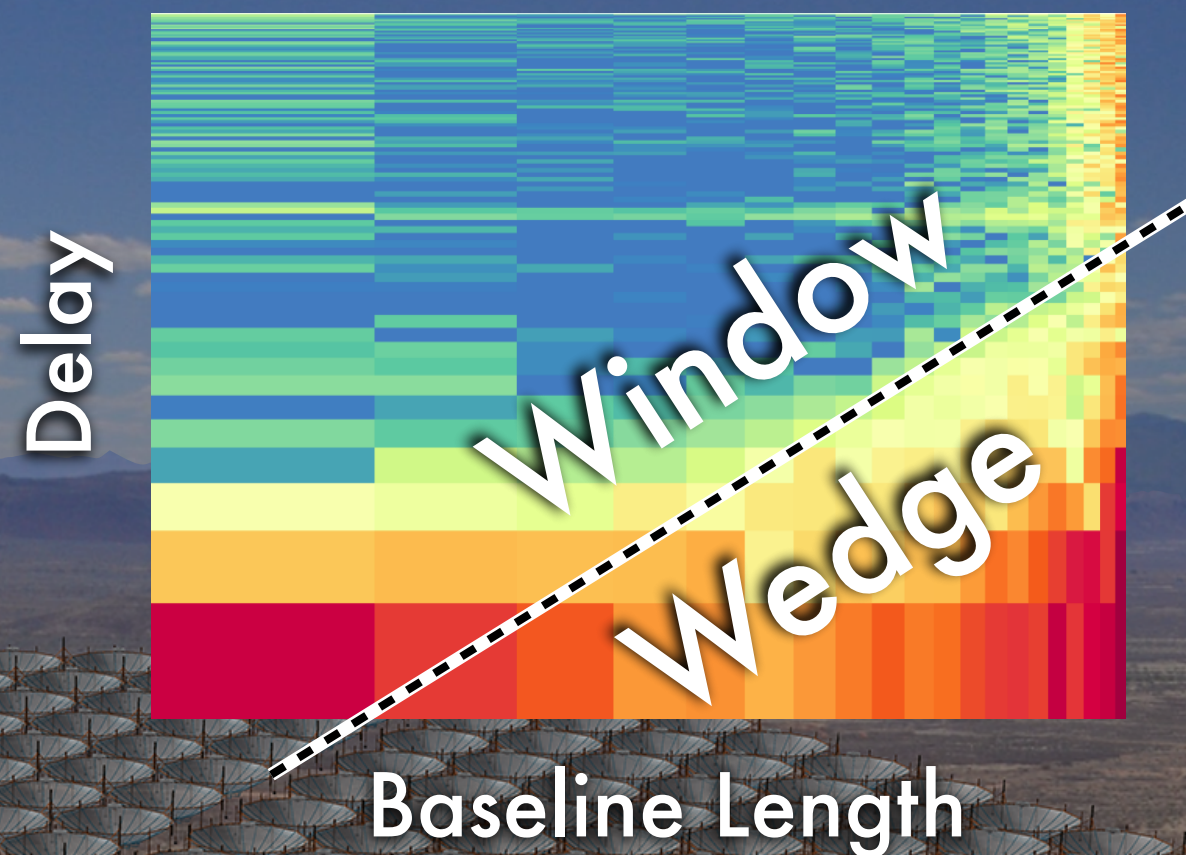
The maximum delay of foregrounds for a baseline is simply the light travel time.



$$\Delta t_{\max} = |\mathbf{b}|/c$$



Our design for
HERA's configuration
maximizes sensitivity
on short baselines.

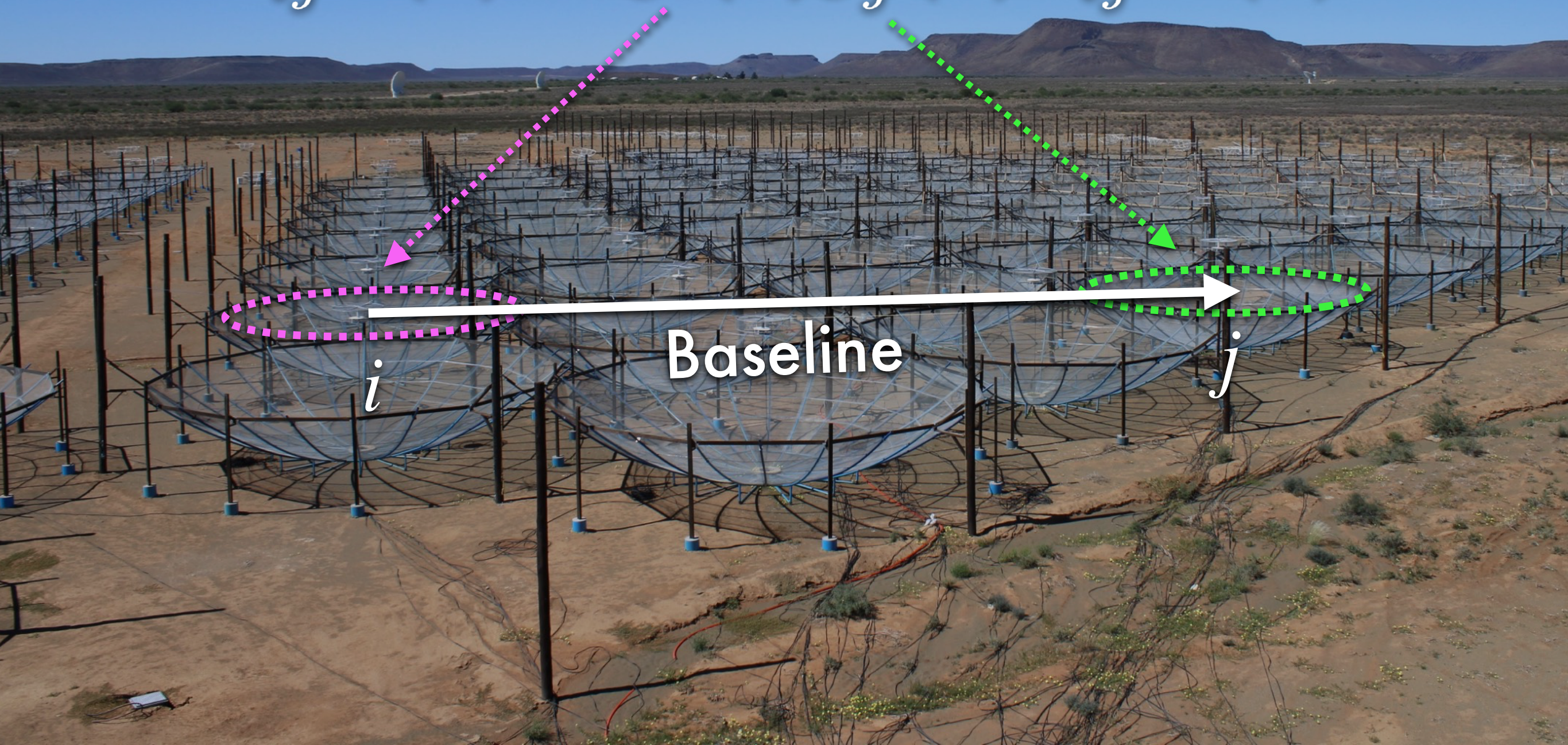


$$\Delta t_{\max} = |\mathbf{b}|/c$$

Working outside the wedge
manages our ignorance — we
trade sensitivity for robustness.

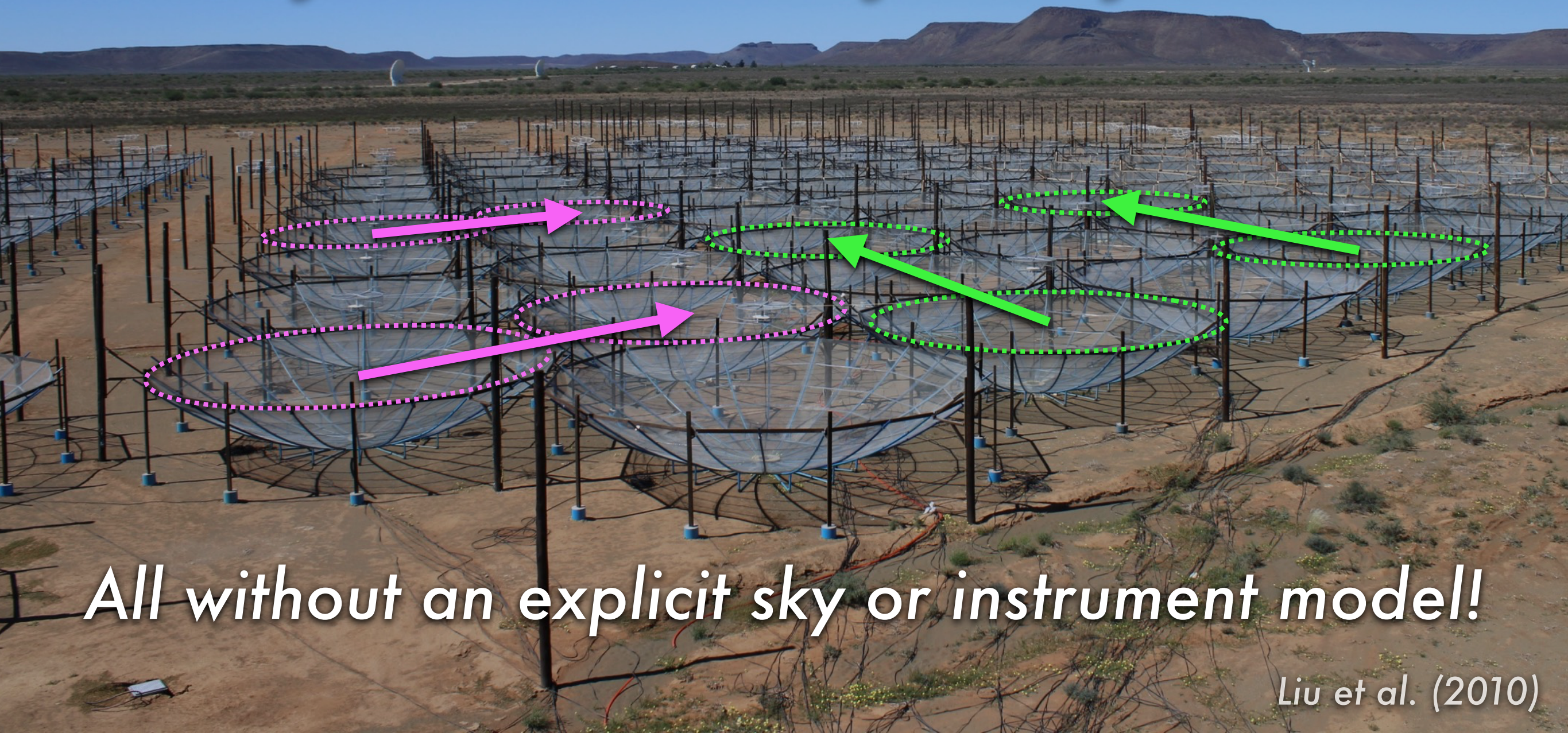
Foreground avoidance won't work without precision calibration.

$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$



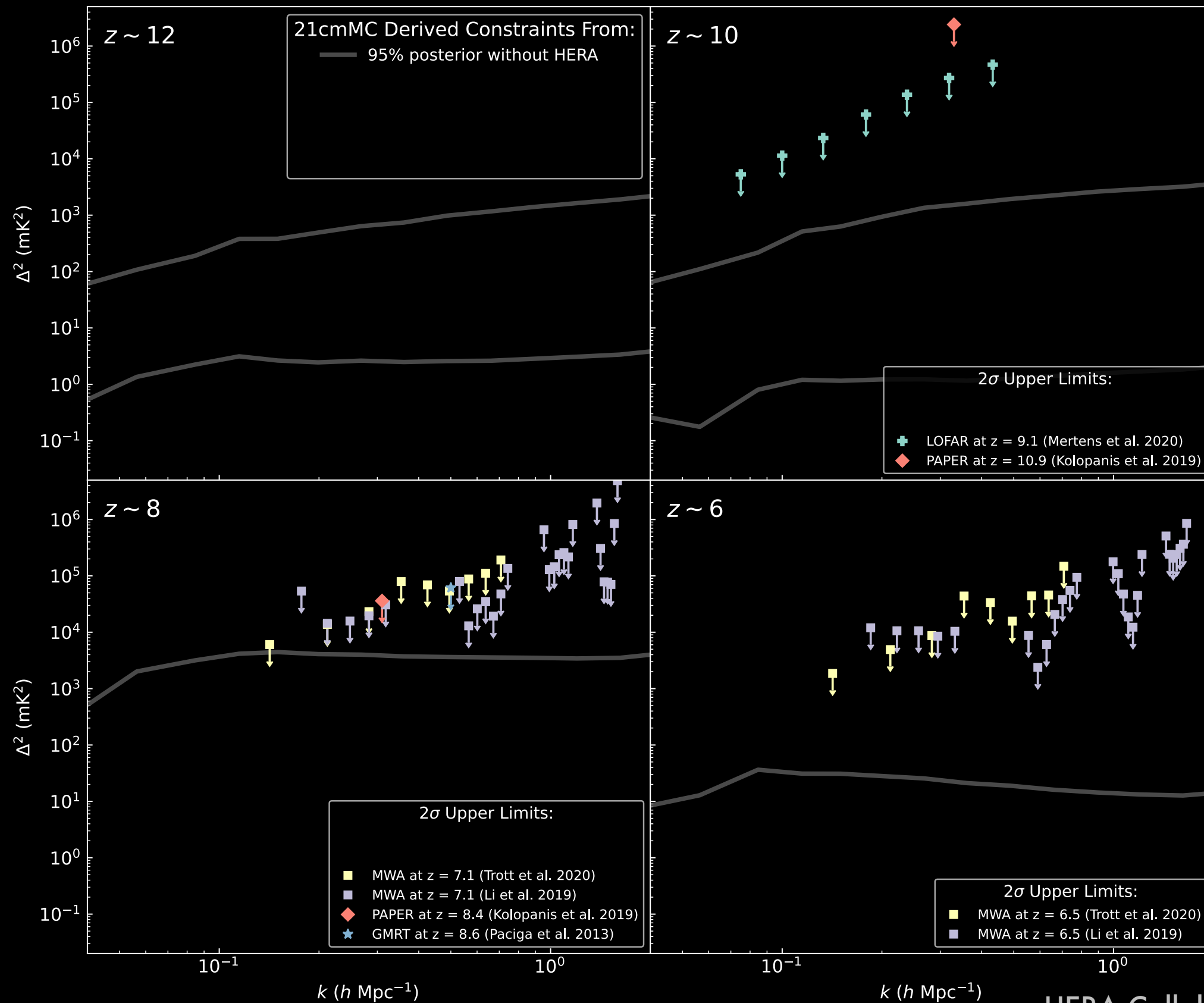
HERA was designed to be calibrated using the internal consistency of redundant baselines.

$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$

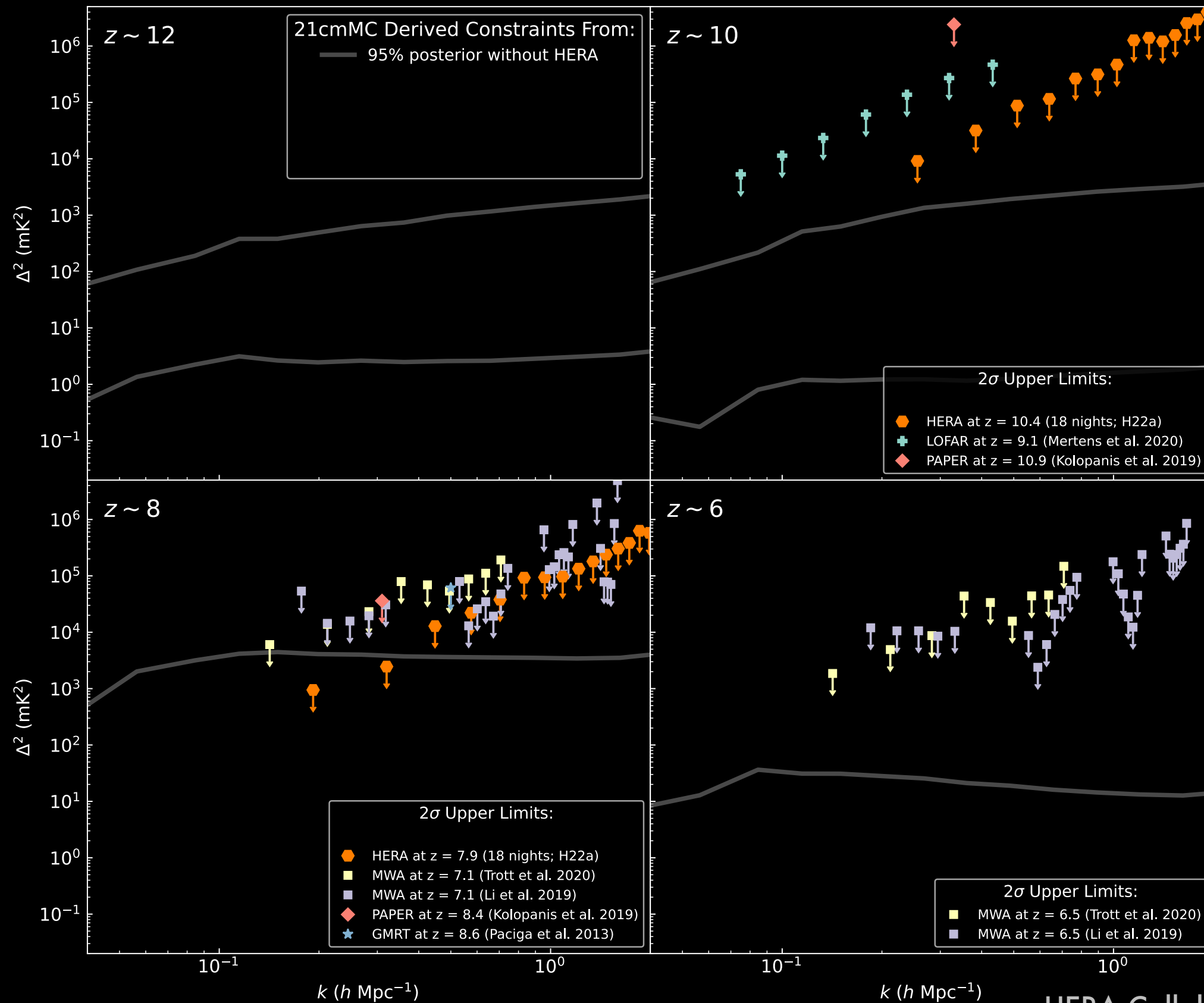


All without an explicit sky or instrument model!

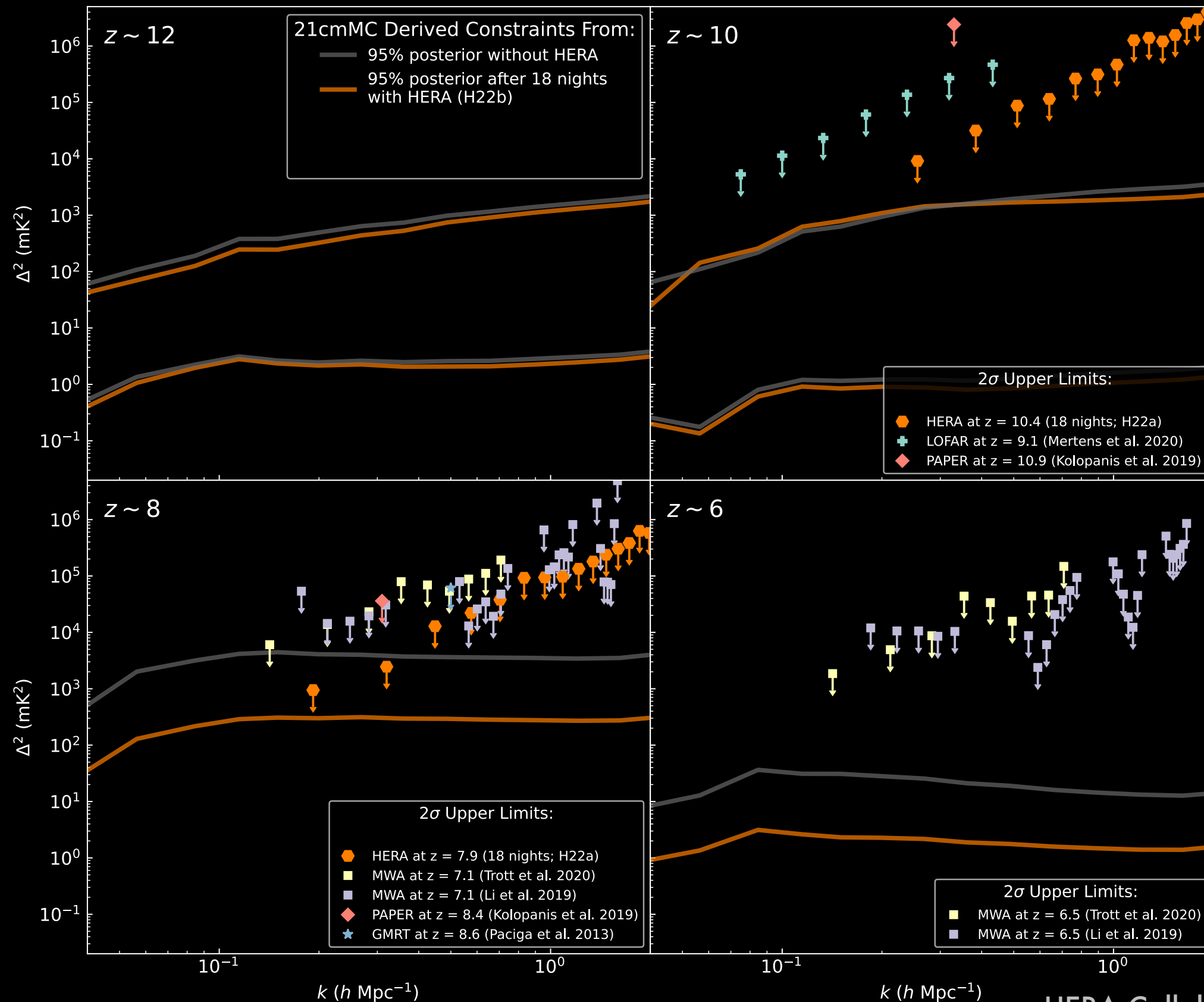
So again, here's where we were before HERA.

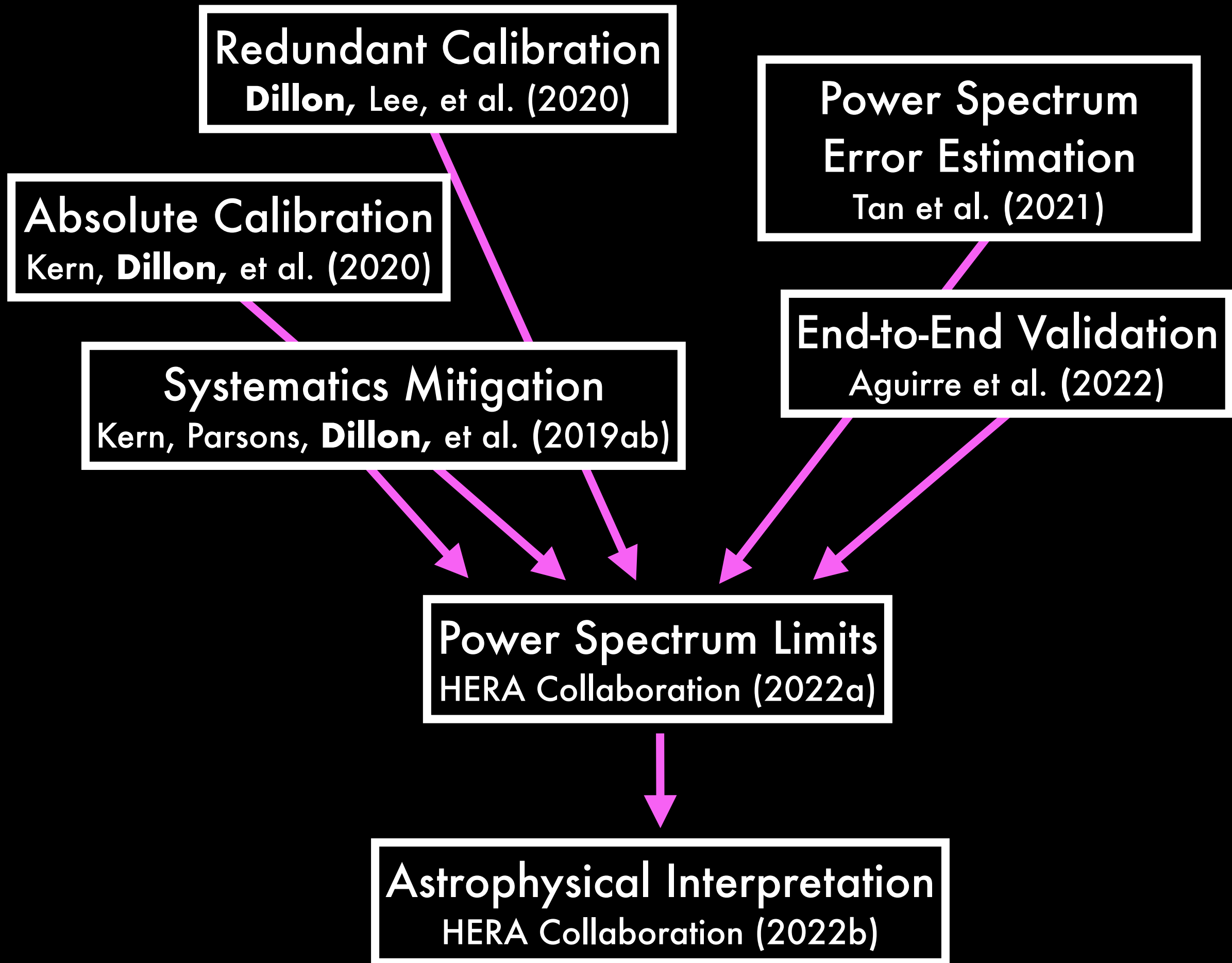


With just 18 nights and 40 antennas, HERA set world-leading upper limits on the 21 cm power spectrum.

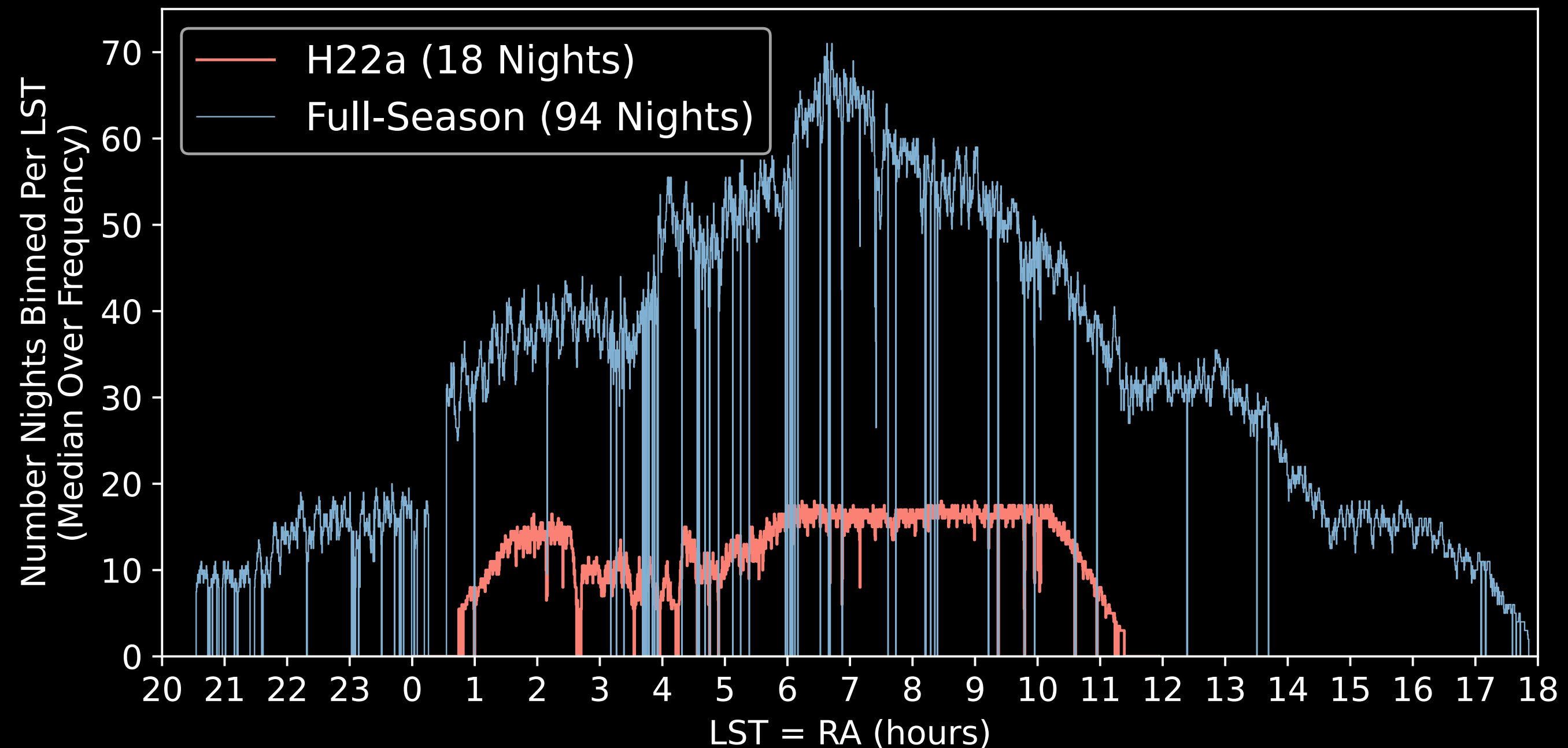


Which constrained the space of models, largely by ruling out an IGM unheated by X-rays at $z=8$.

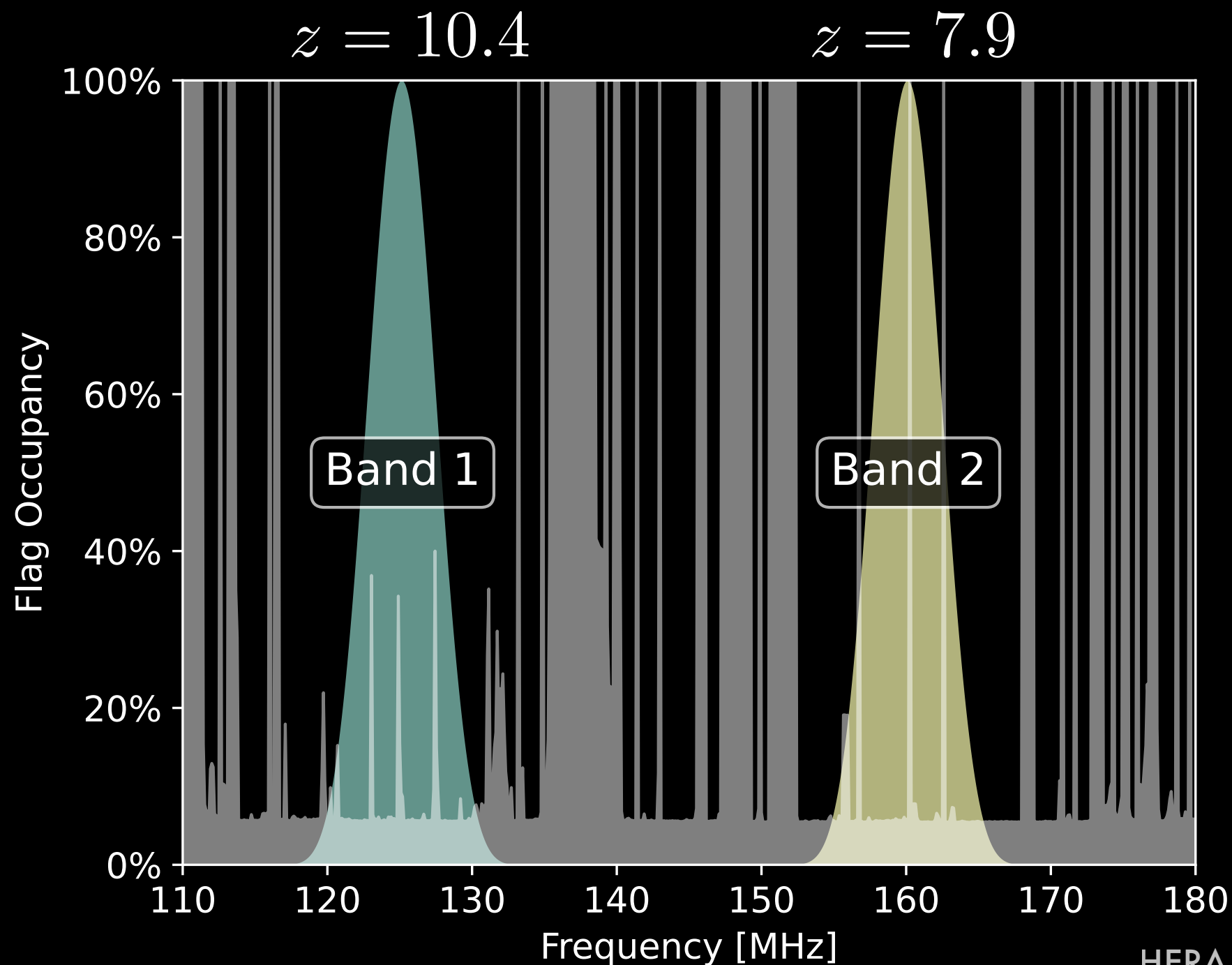




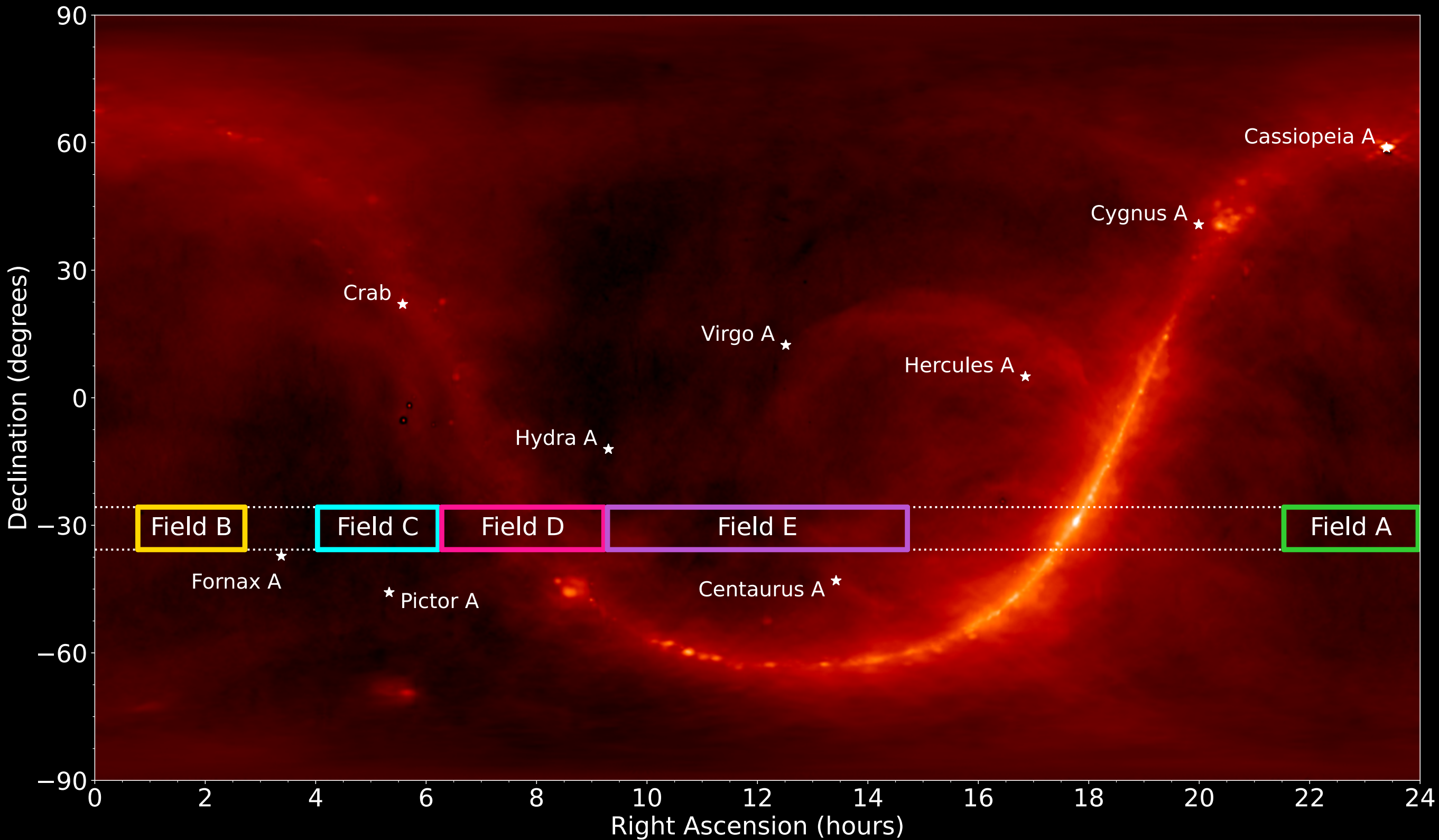
But all that was with only 18 nights
of data... and we had 94 good
nights from that season.



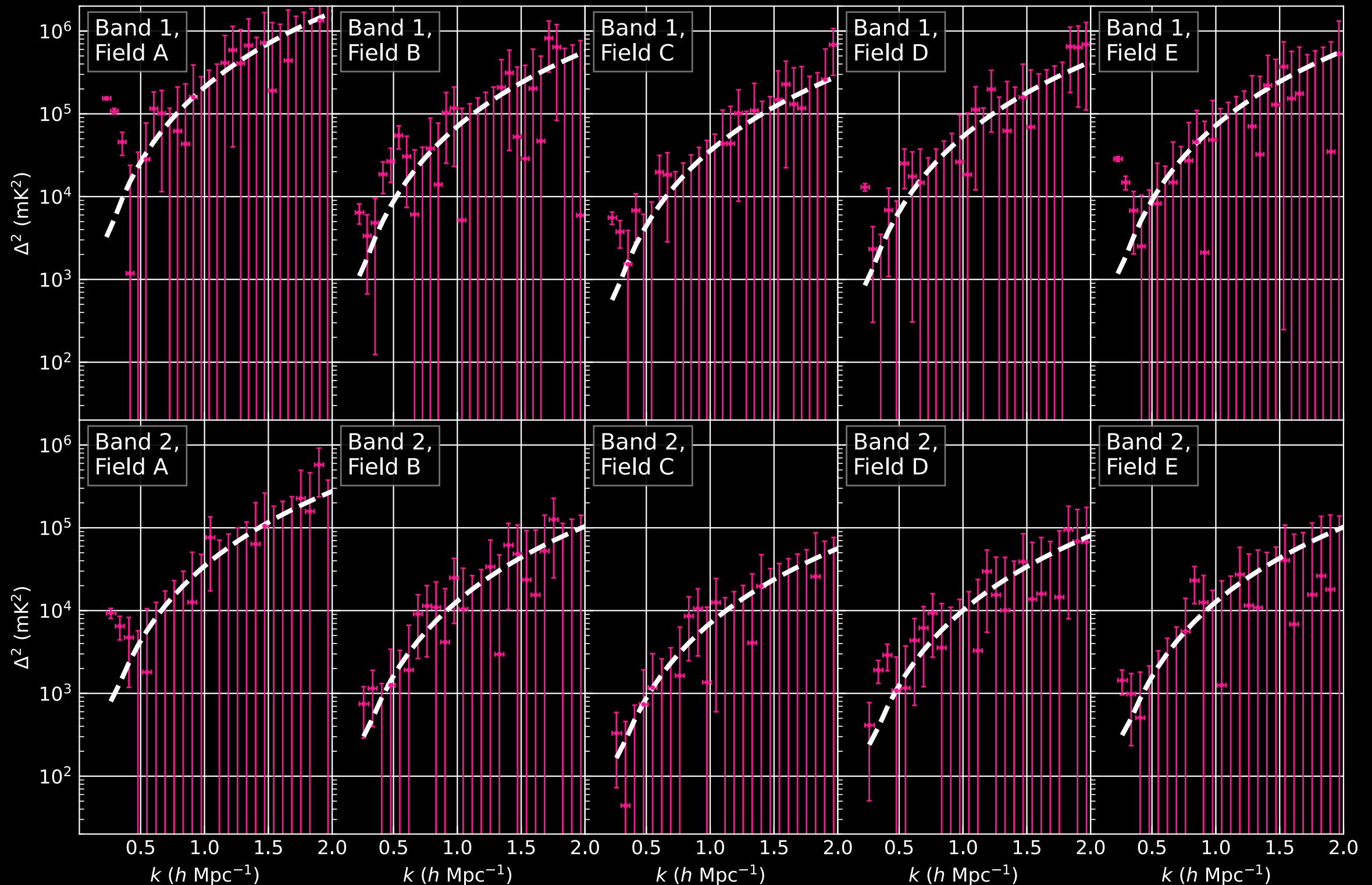
Adapting the analysis techniques to the larger data set, we picked two frequency bands with minimal RFI contamination.



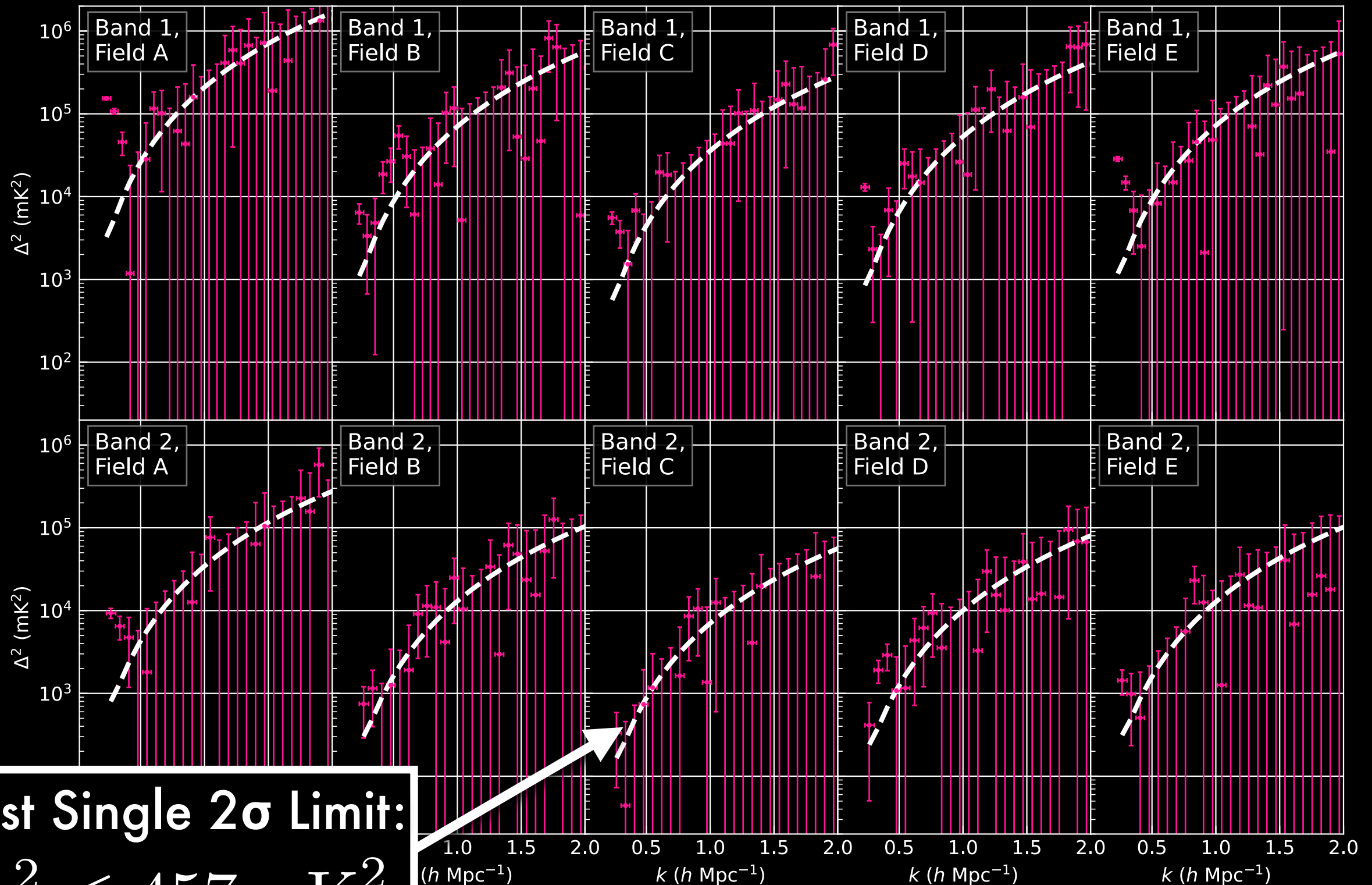
We divided our observed LSTs into five fields.



And thus set power spectrum upper limits across bands and fields.



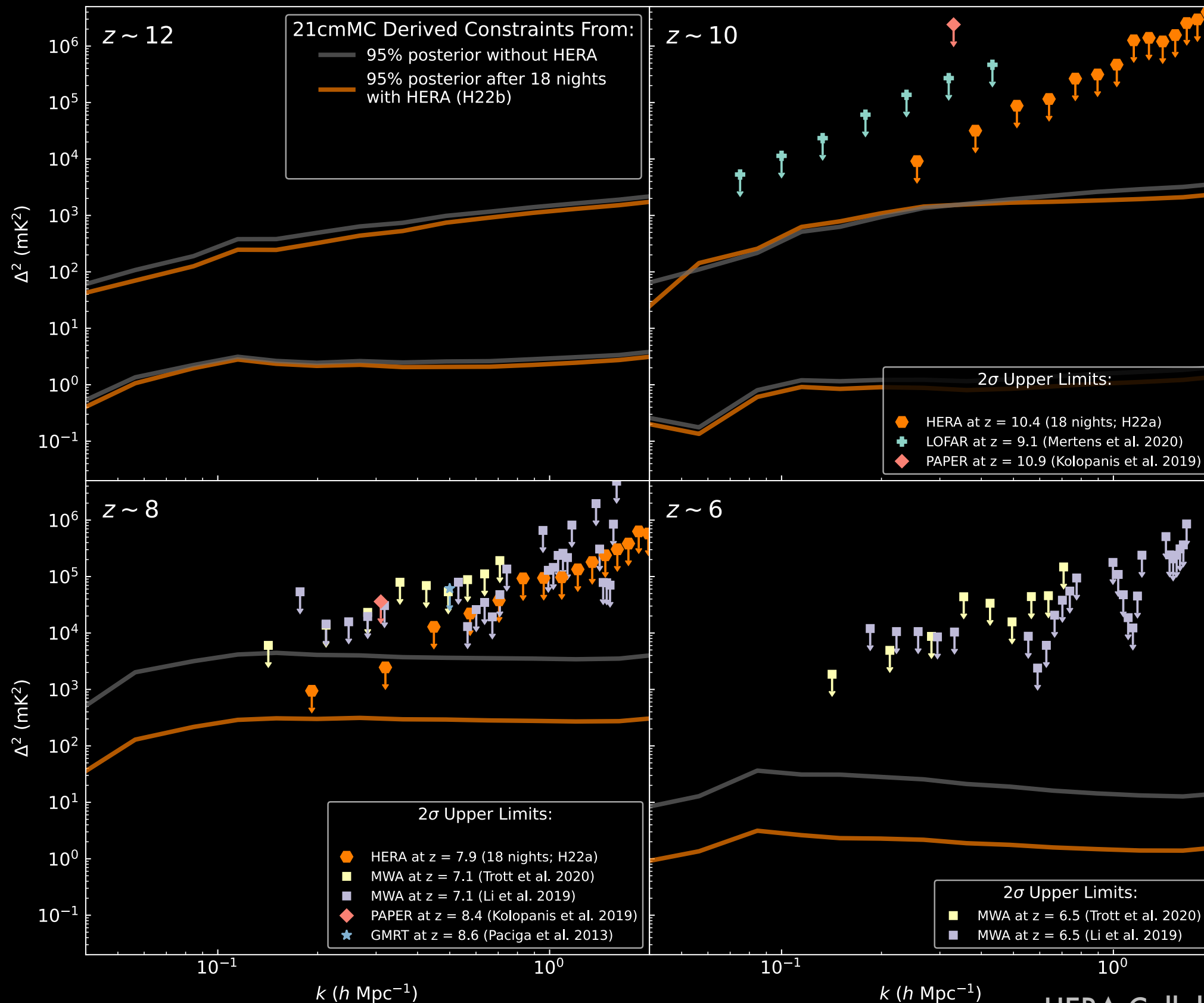
And thus set power spectrum upper limits across bands and fields.



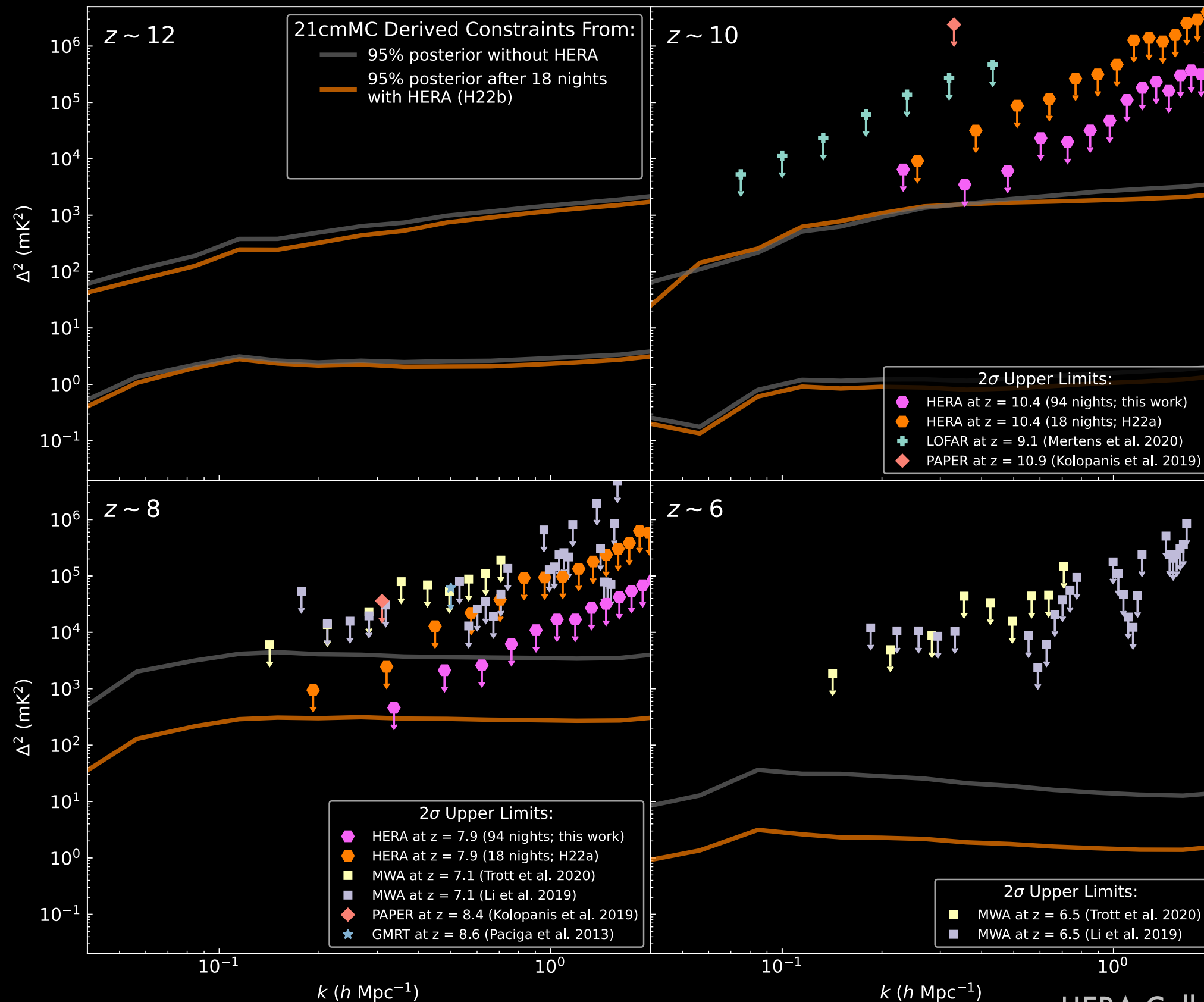
Best Single 2 σ Limit:

$$\Delta^2 \leq 457 \text{ mK}^2$$

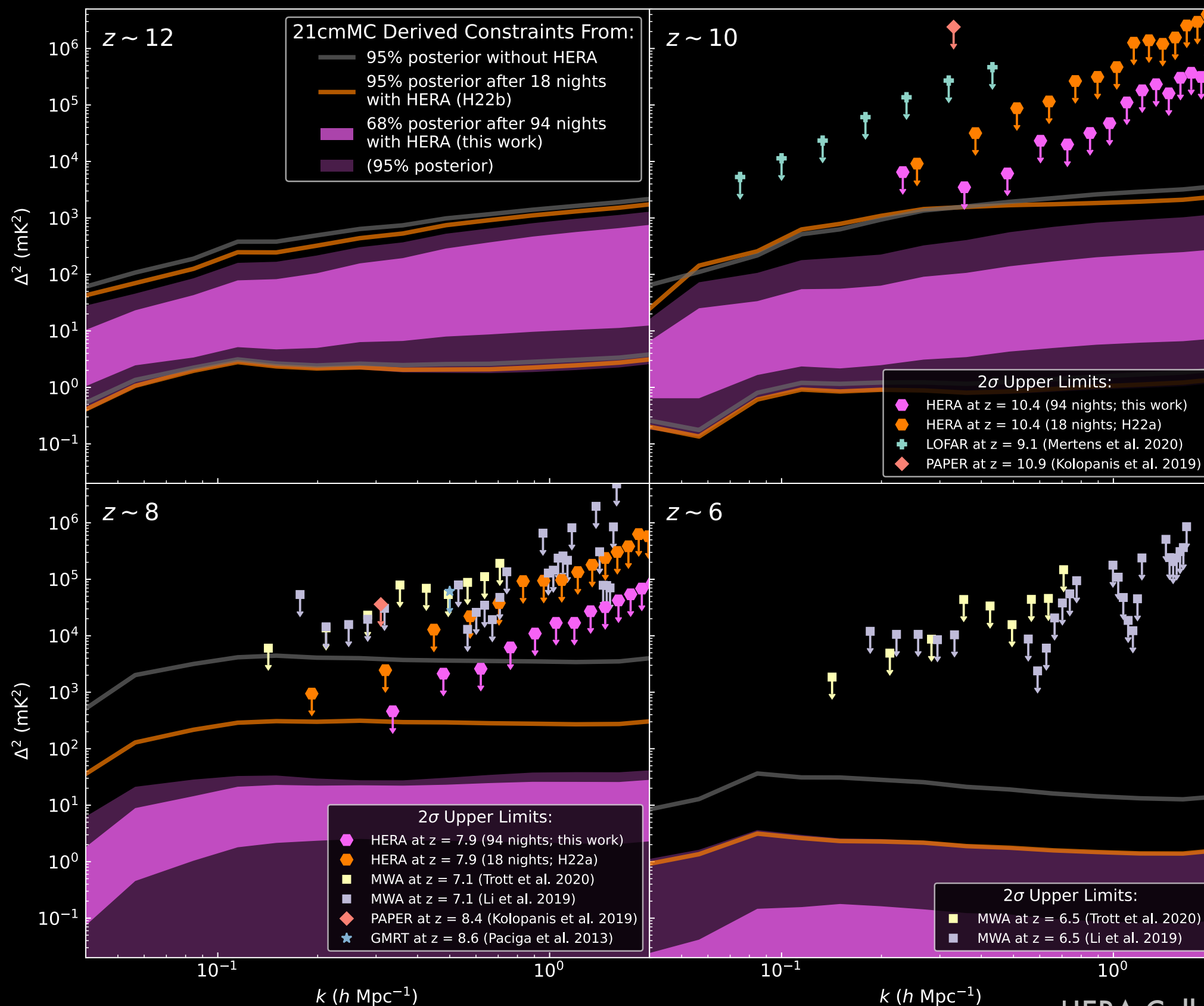
So here's where we were again...



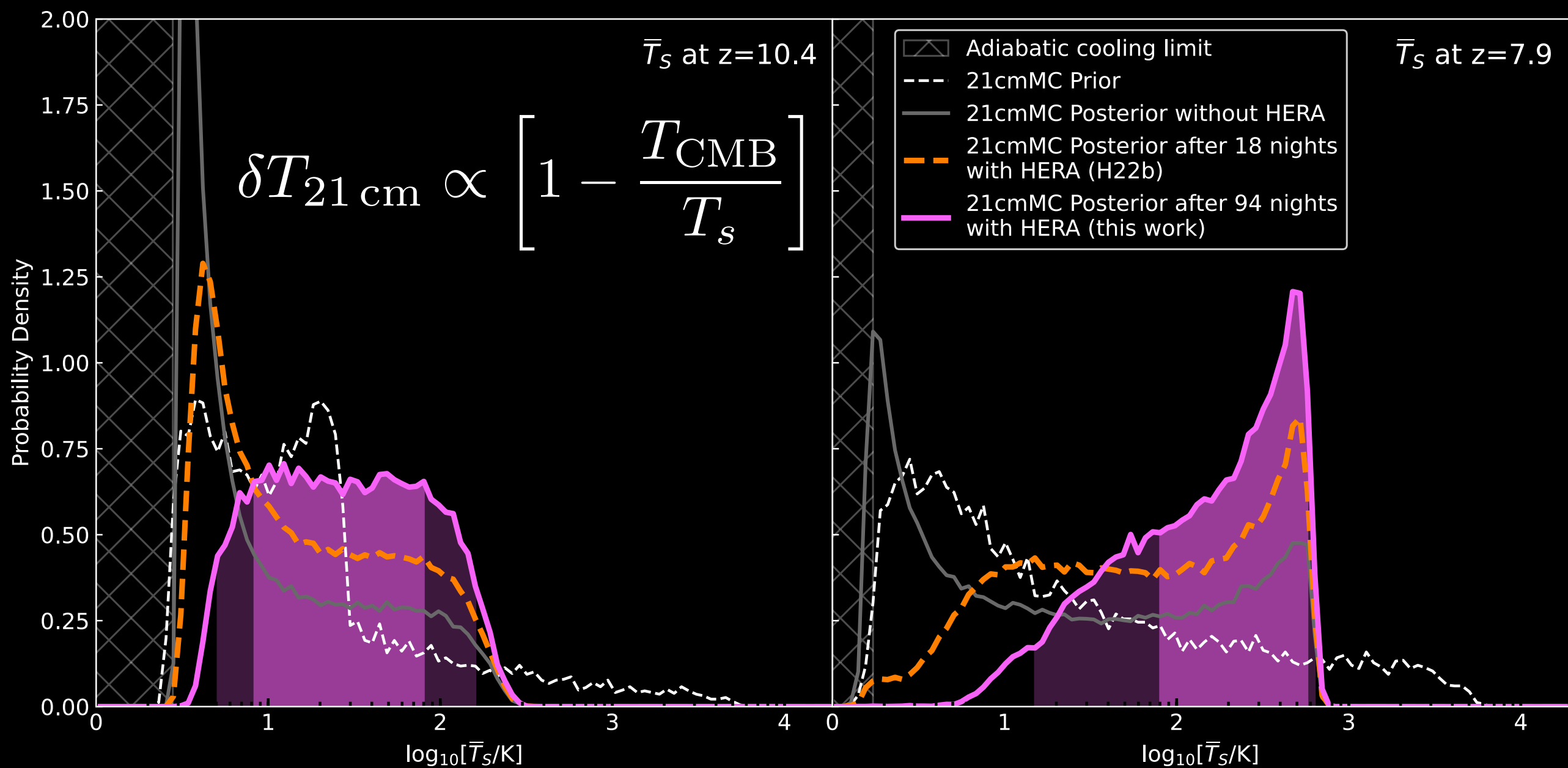
With a full season, our limits come down by more than a factor of 2 at both redshifts.



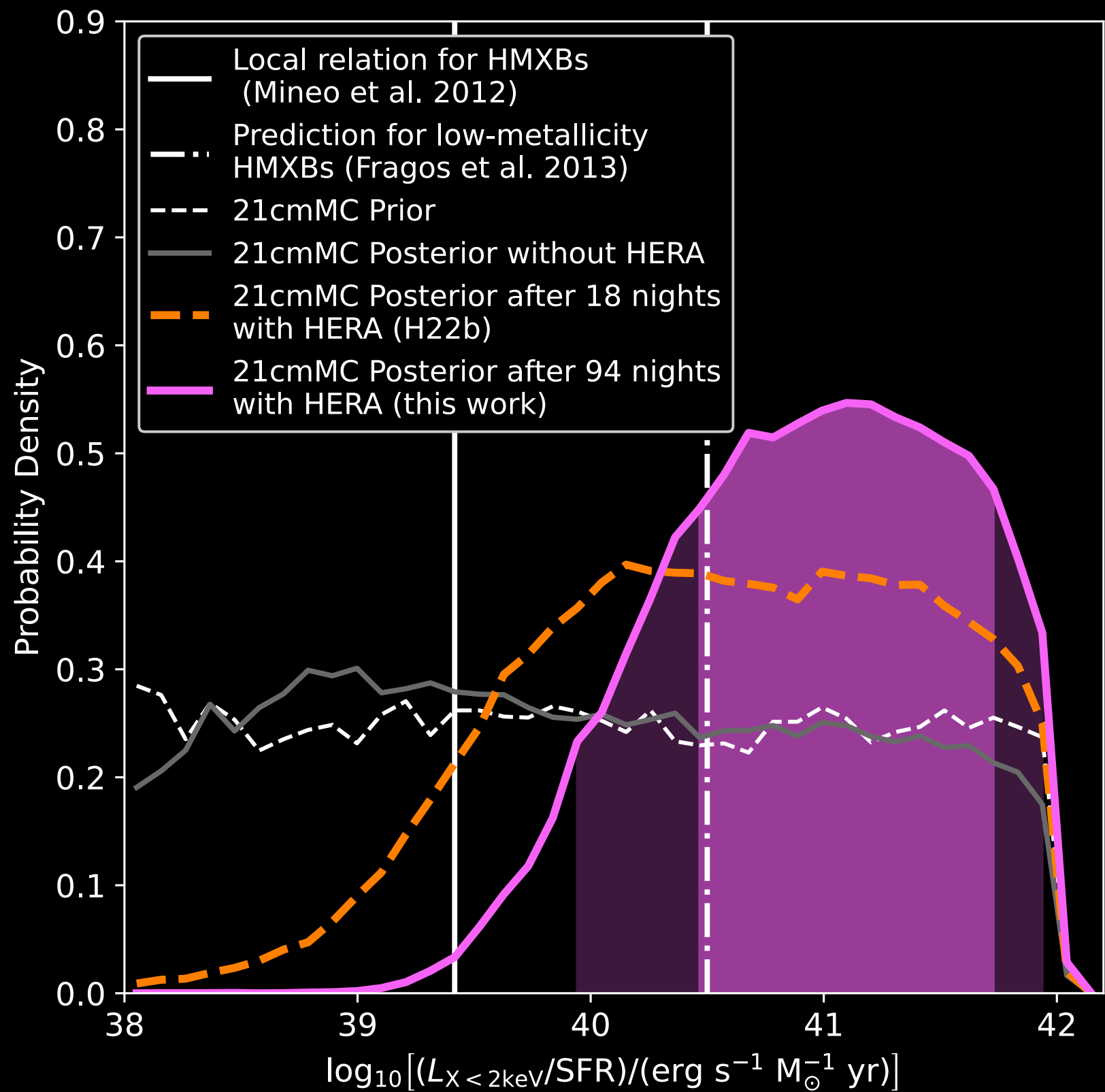
Our posterior for the power spectrum with 21cmMC tightens substantially.



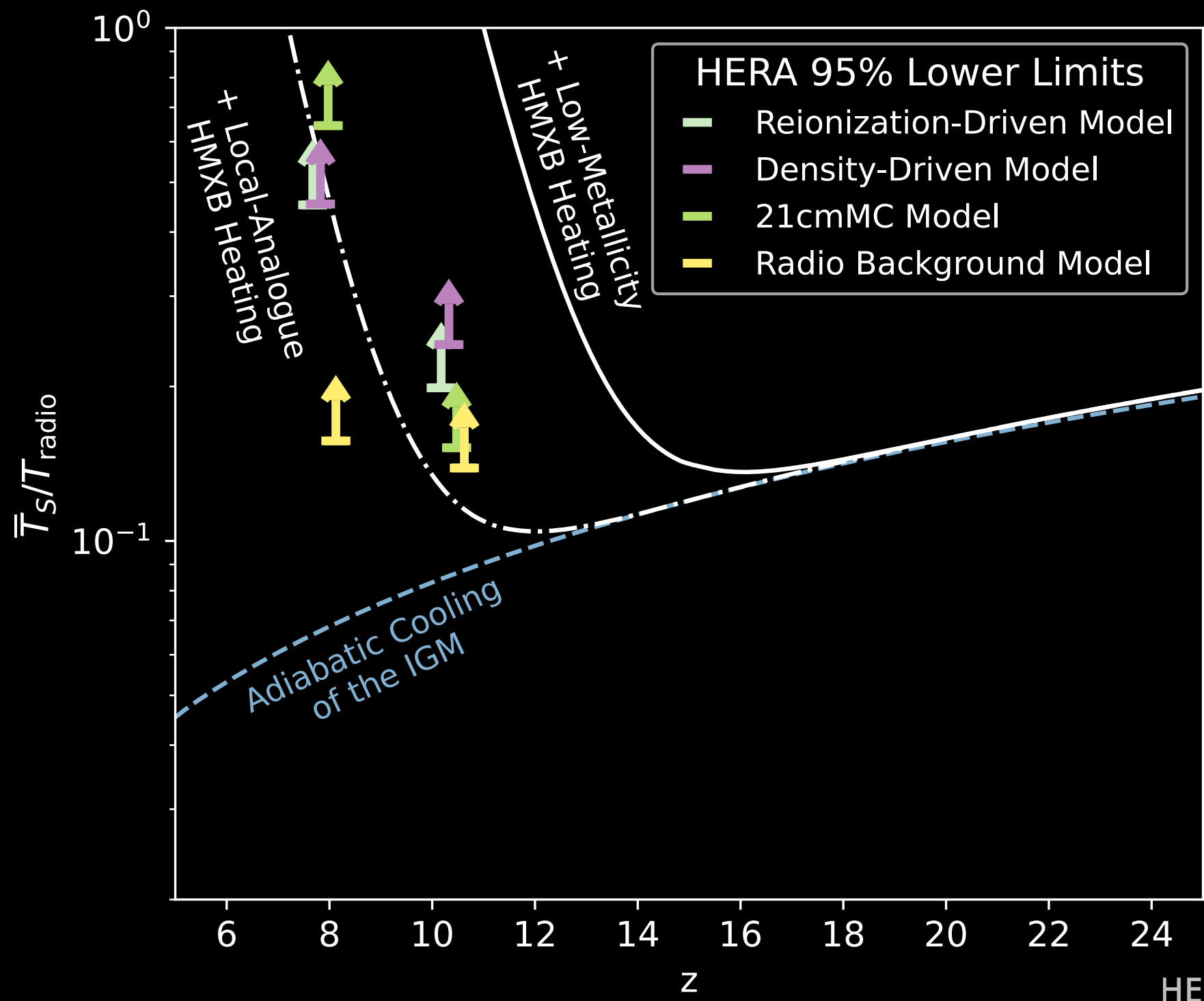
The big shift comes from showing the IGM was heated by $z = 10.4$, since a cold IGM produces a bright 21 cm signal.



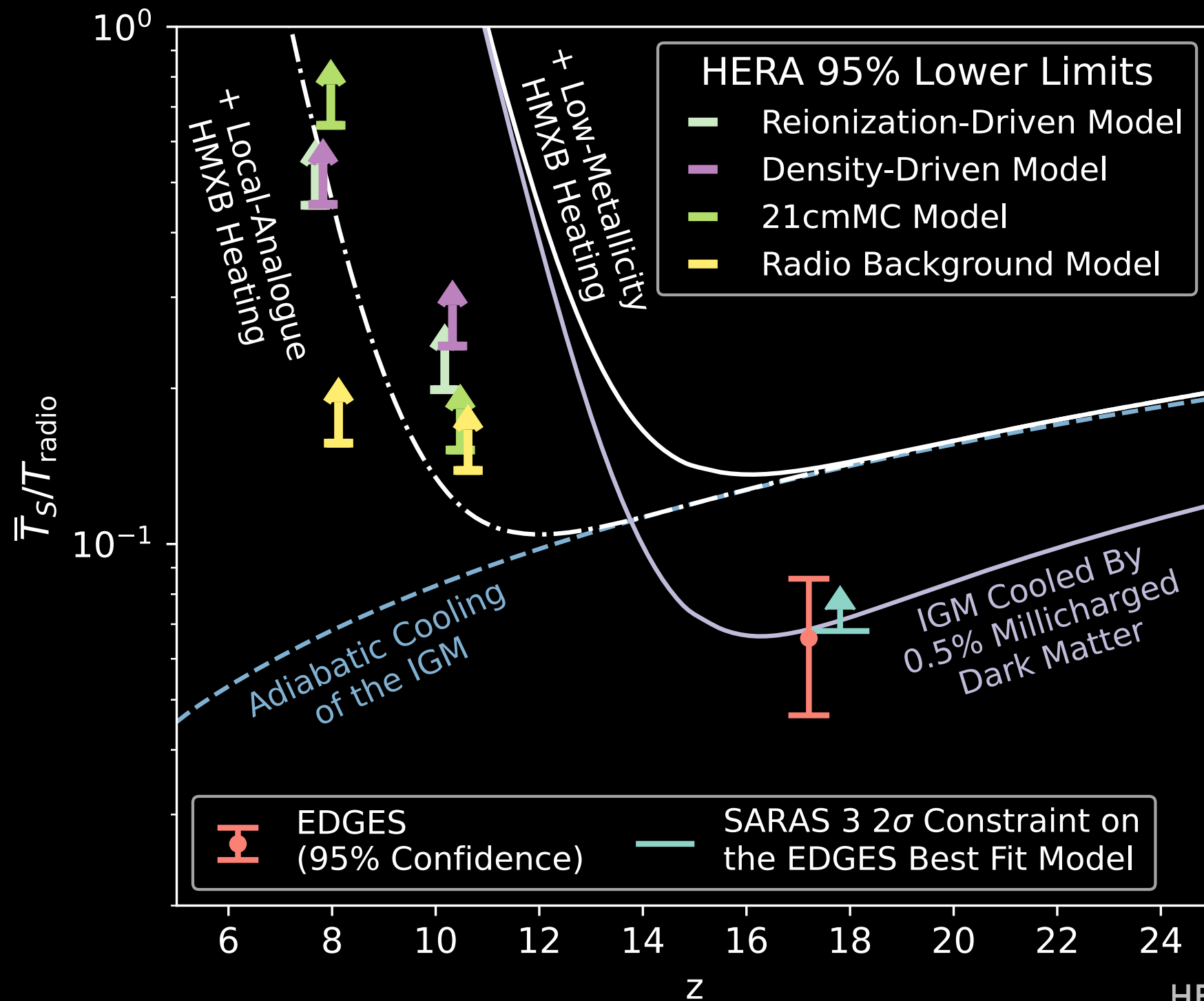
If the IGM was heated by high-mass X-ray binaries — as is generally believed — this result rules out high-metallicity HMXBs (which are less X-ray-efficient per unit SFR) and thus requires heating driven by evolved low-metallicity stars.



Four independent theoretical models agree the IGM was heated before $z=10.4$, likely by low-metallicity HMXB.



However, we are not yet able to say much about the tension between EDGES and SARAS or the exotic models invoked to explain EDGES.



What's next for HERA?

**We just finished an observing
season (~150 good nights)
observing with over 200
antennas as we build out to 350.**



Everything but the dishes is new, including our wide-band Vivaldi feeds that go from 50 – 250 MHz ($4.7 > z > 29$).



Photo: Ziyaad Halday

With a full season and the full array, we'll have the sensitivity to detect the 21 cm signal and distinguish between models.

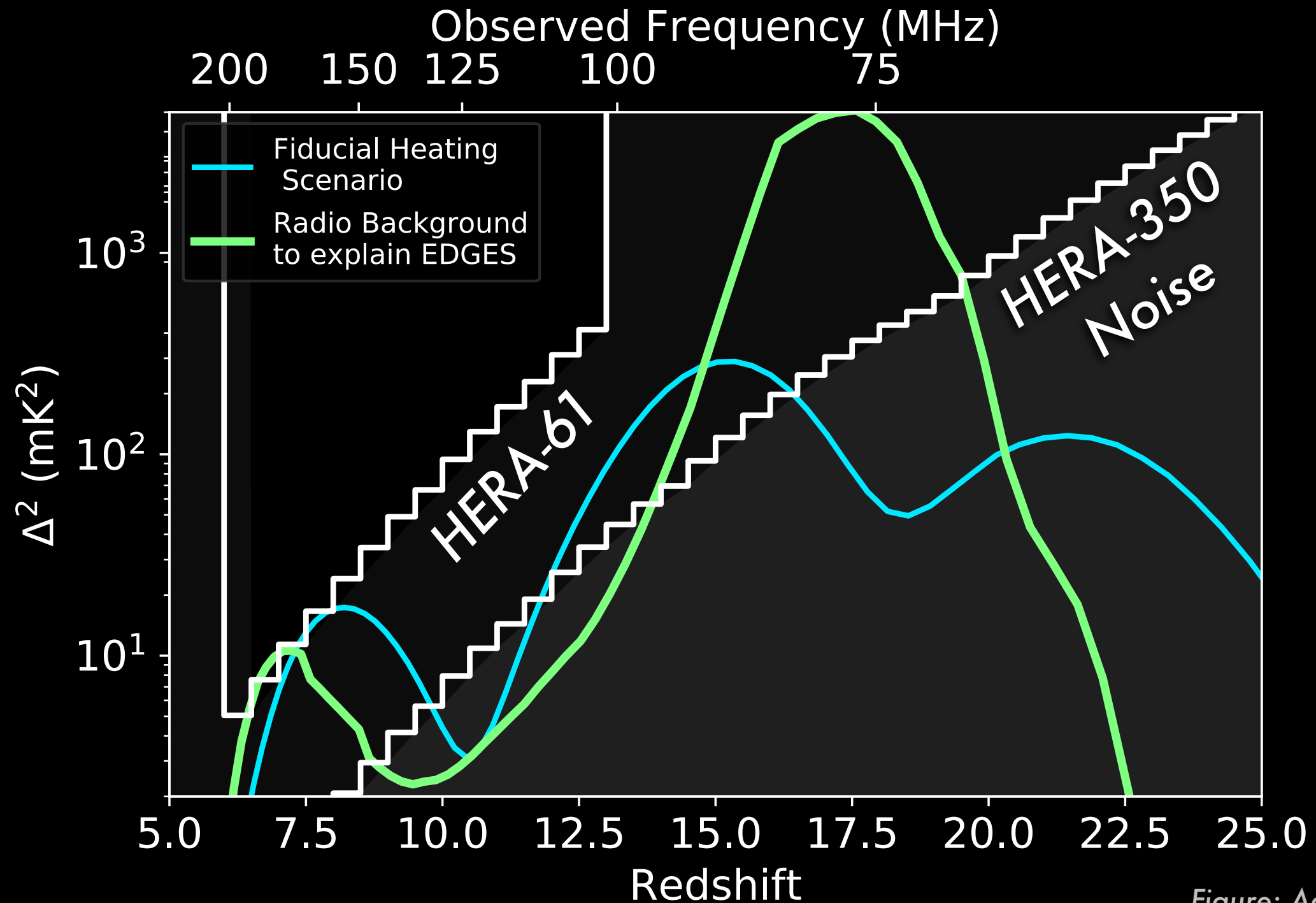
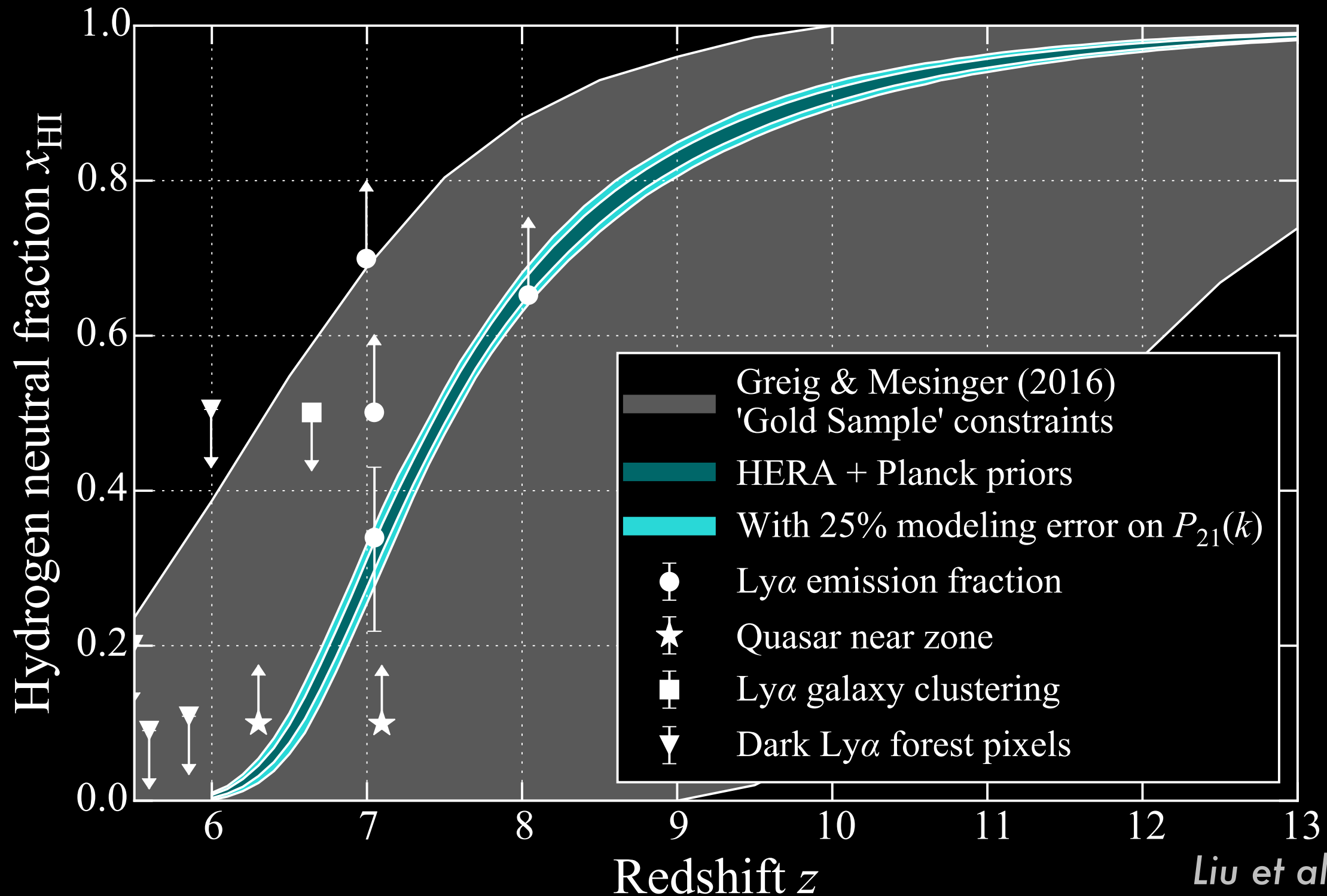
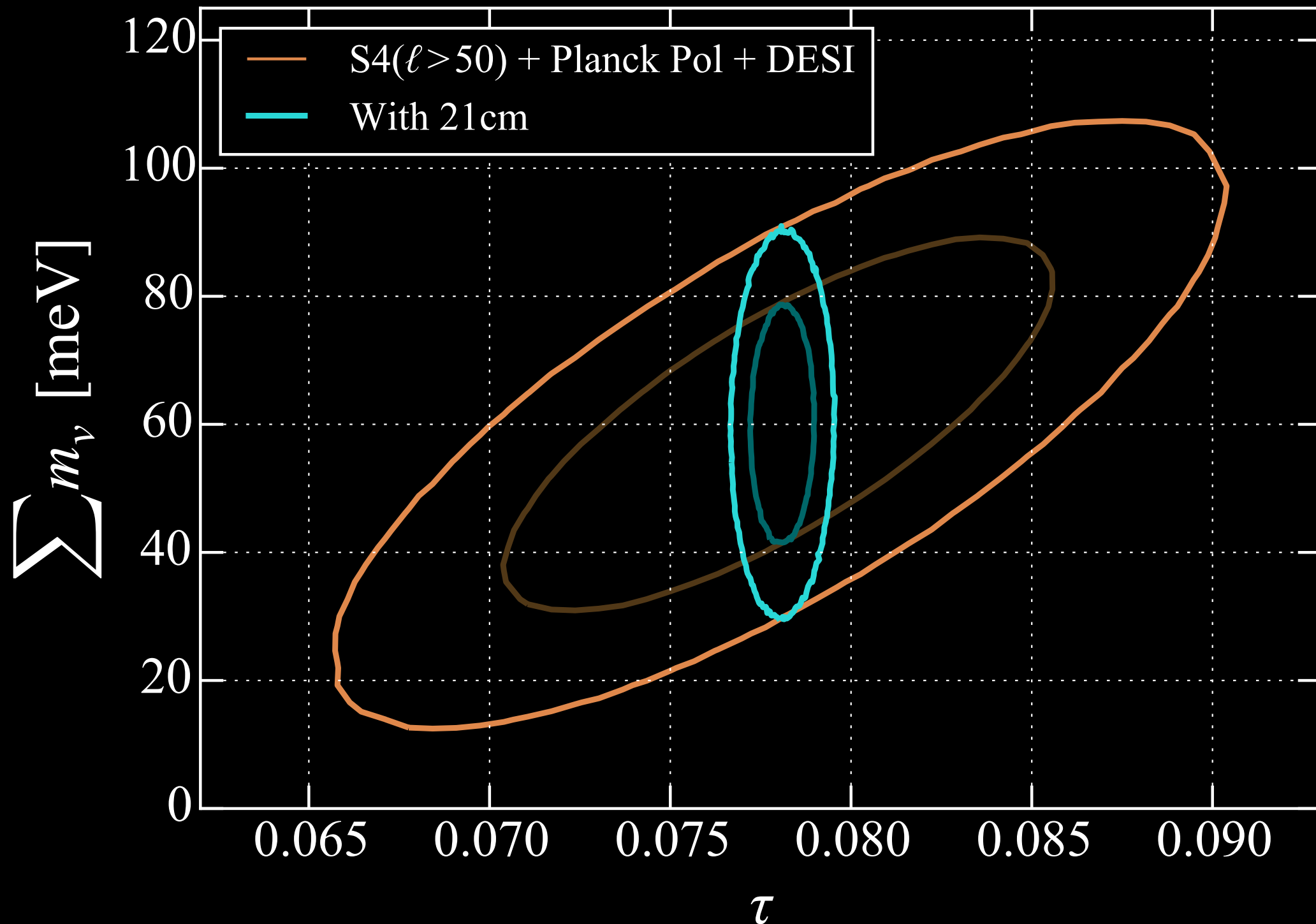


Figure: Aaron Ewall-Wice

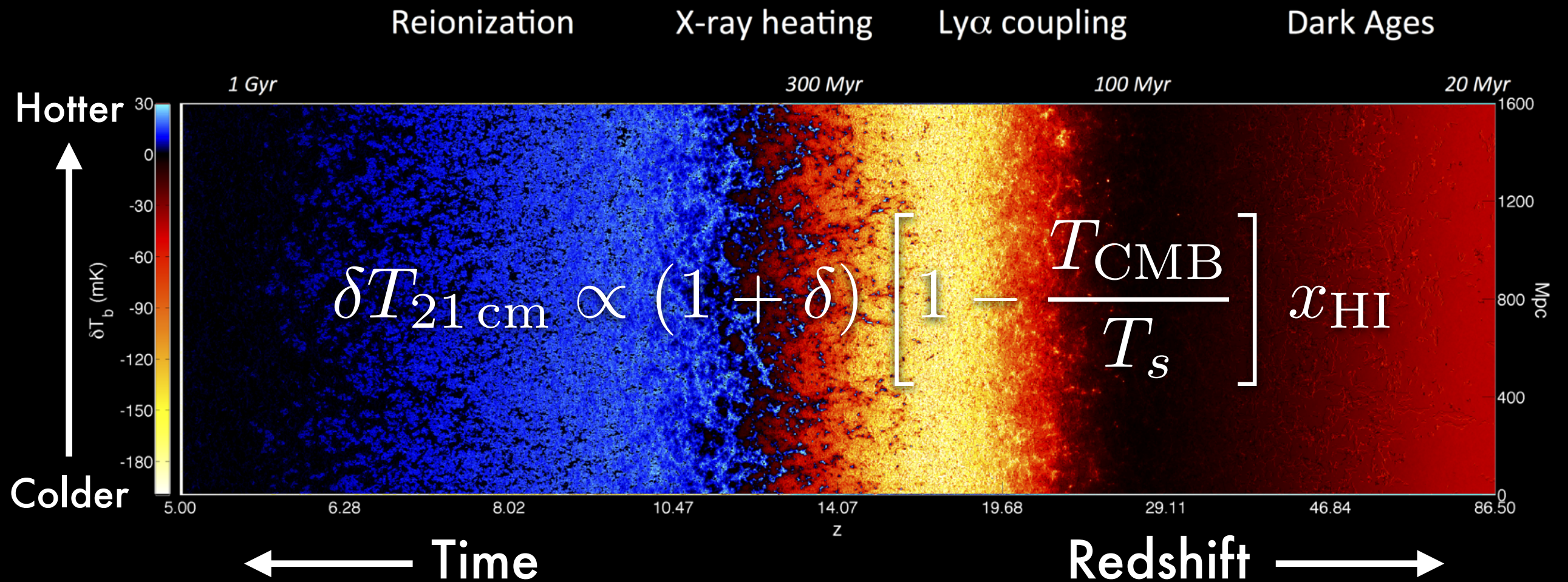
Which means we can precisely measure the ionization history of the universe.



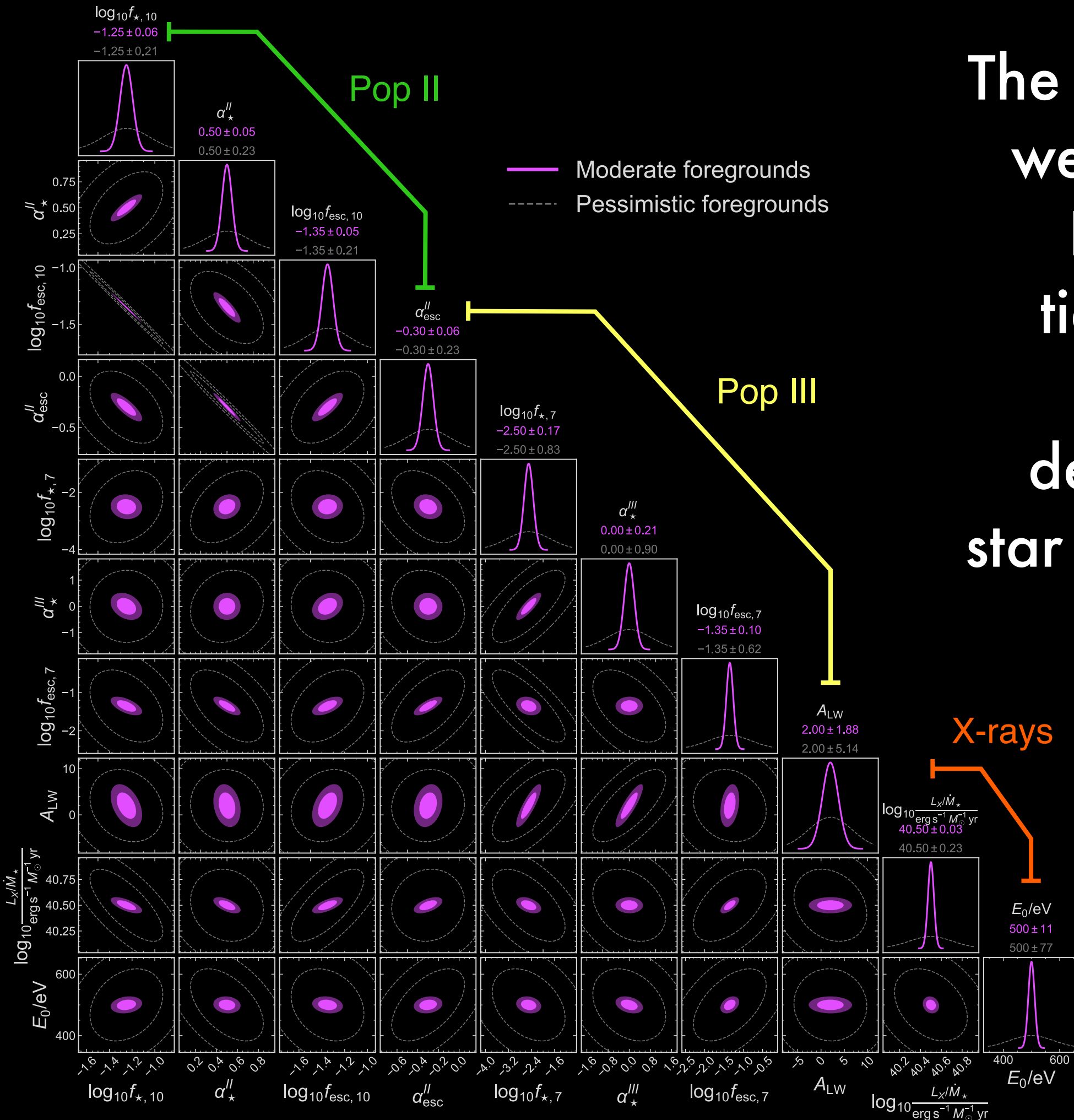
And, perhaps increase the significance of a detection of non-zero Σm_ν with CMB-S4.



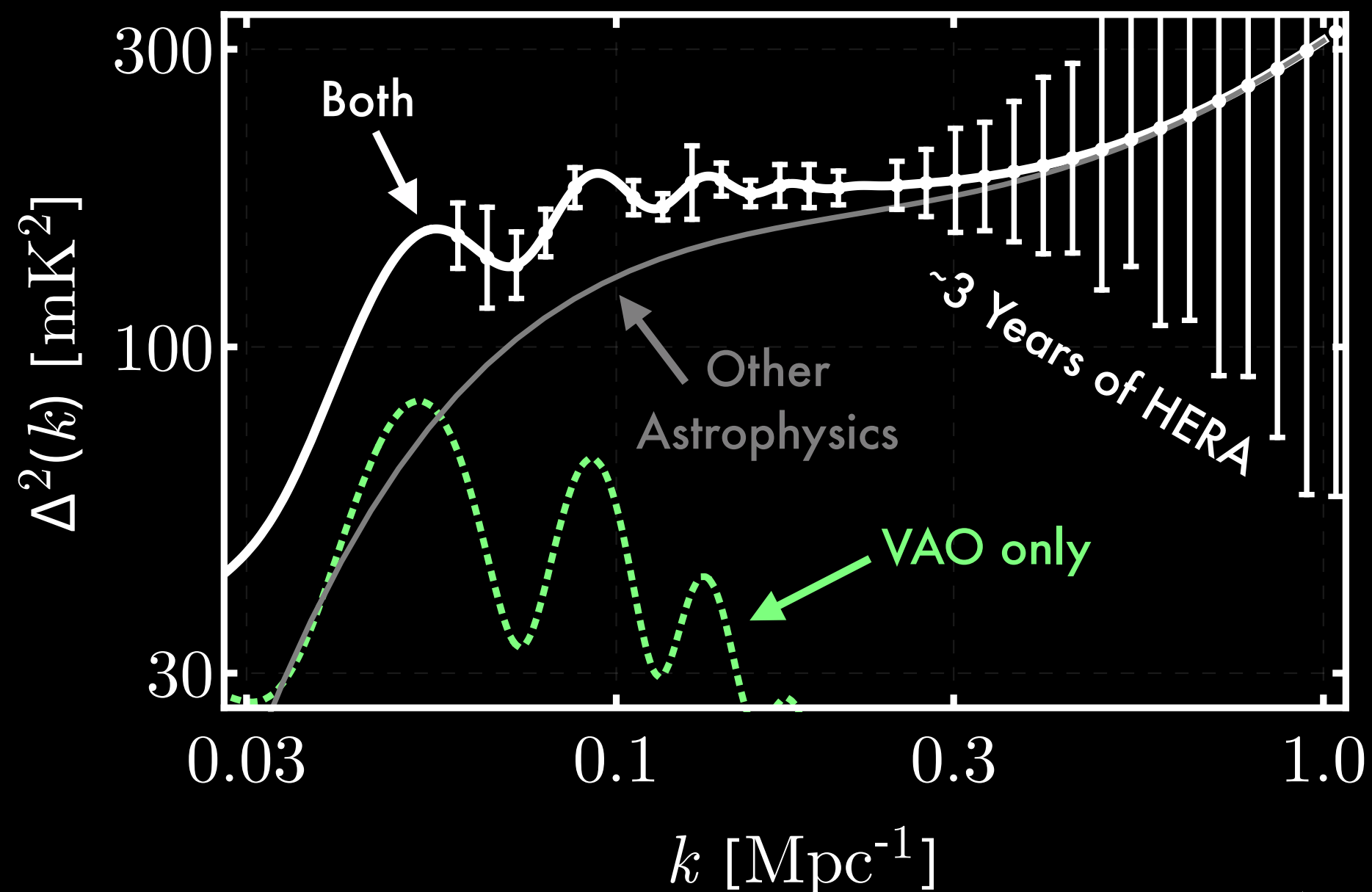
There's also complex, interconnected astrophysics to explore before the EoR, even if EDGES is wrong.



The power spectra we measure with HERA will also tightly constrain parameters describing early star formation and X-ray heating.

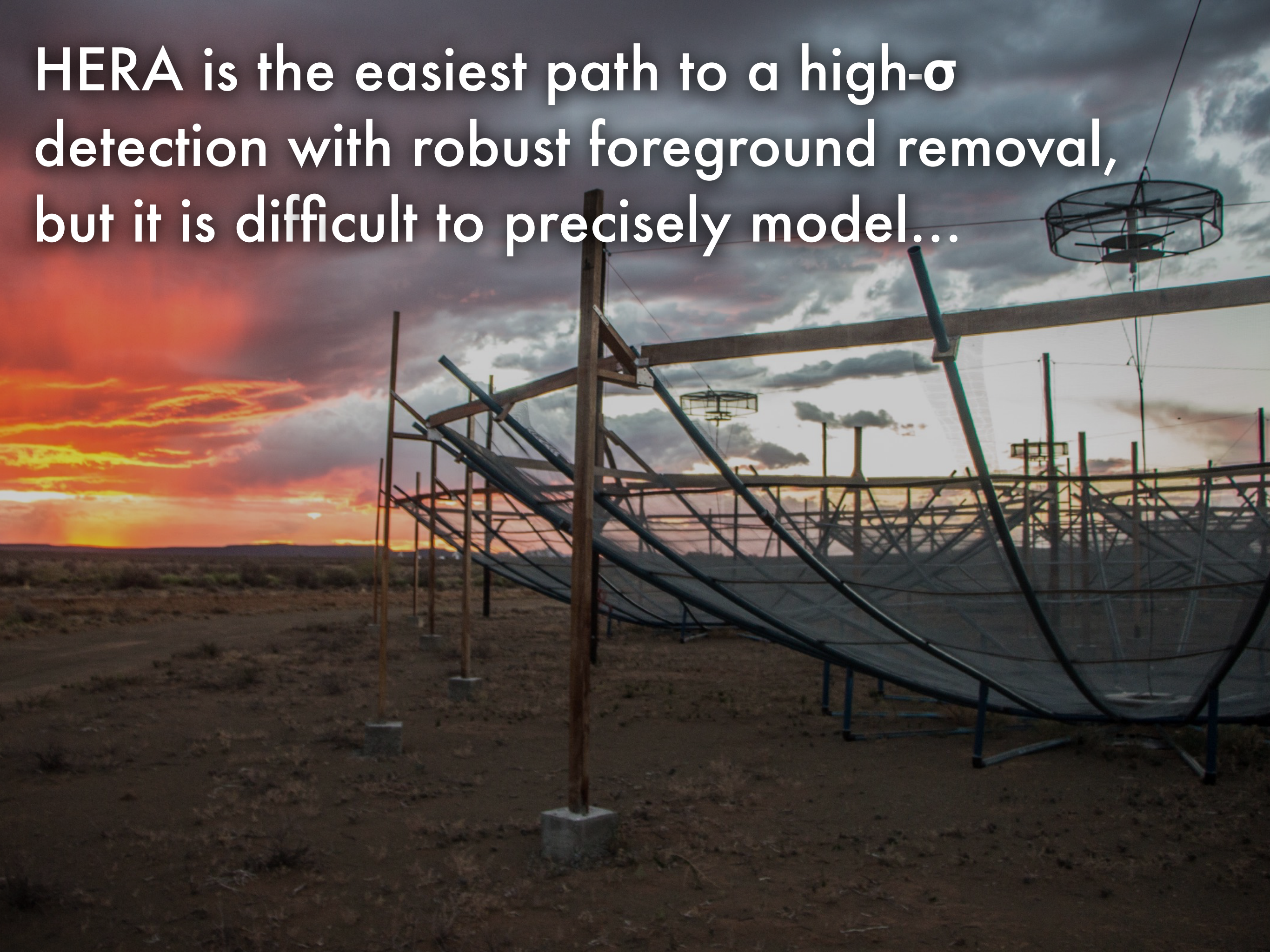


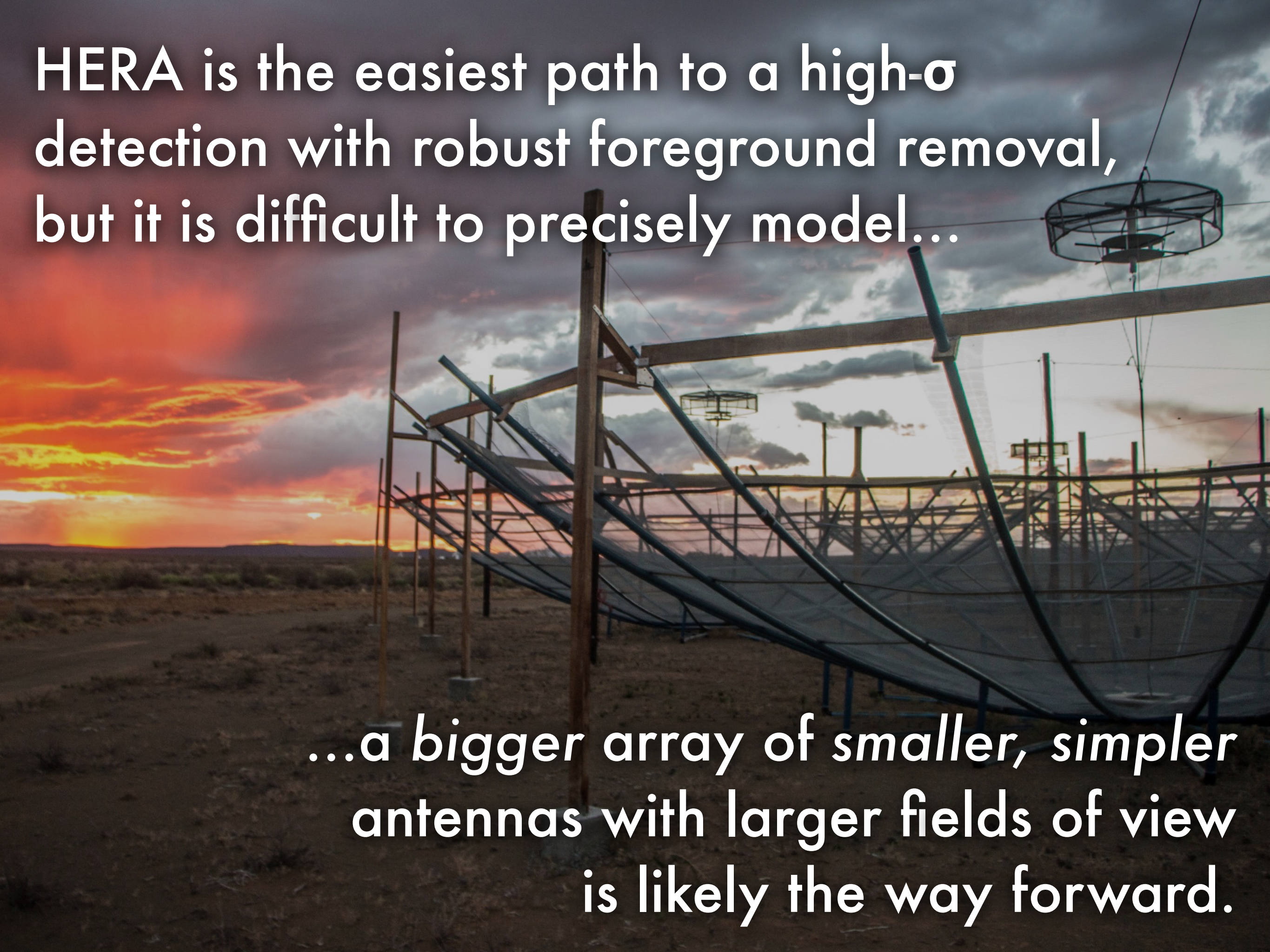
With a few years of observing, we may detect velocity acoustic oscillations, providing a new standard ruler at $z \approx 16$.



What comes after HERA?

HERA is the easiest path to a high- σ detection with robust foreground removal, but it is difficult to precisely model...



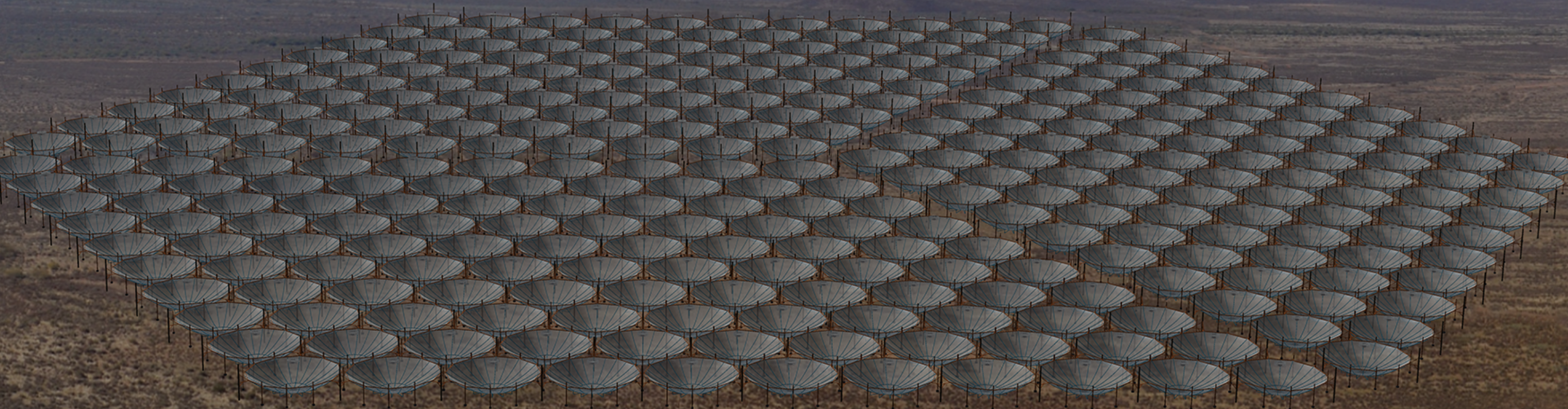
A photograph of a large radio telescope array, likely the Murchison Widefield Array (MWA), at sunset. The sky is a mix of orange, red, and grey clouds. The foreground shows the complex metal structure of the telescope, including a large parabolic dish and various support beams. The ground is a flat, open field.

HERA is the easiest path to a high- σ detection with robust foreground removal, but it is difficult to precisely model...

...a *bigger* array of *smaller, simpler* antennas with larger fields of view is likely the way forward.

There's a problem with how we measure visibilities.

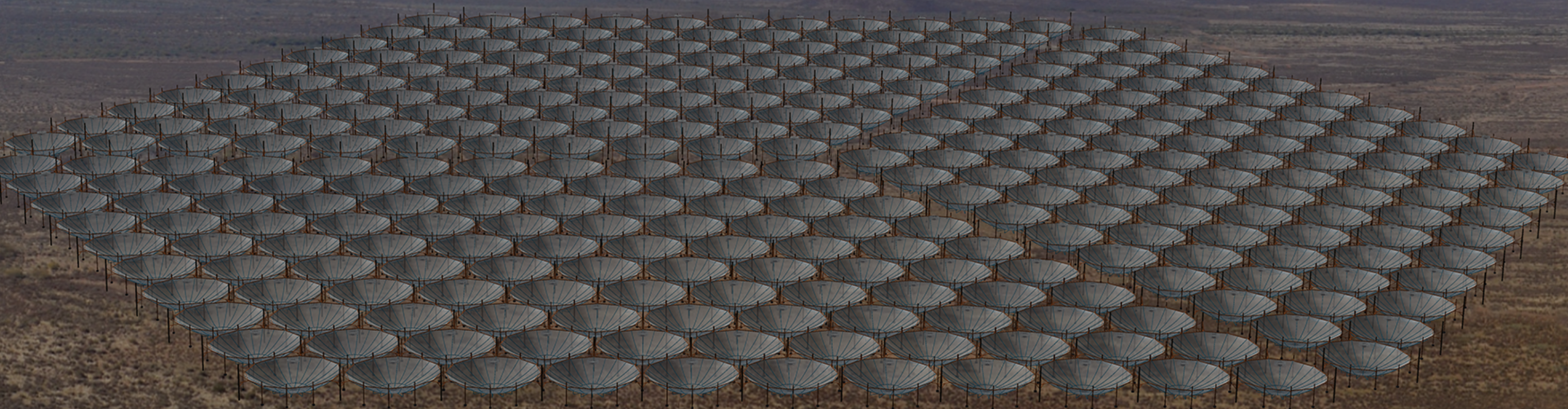
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

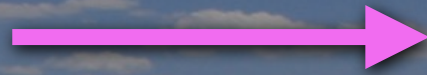
Measure antenna
voltages $v_i(t)$.



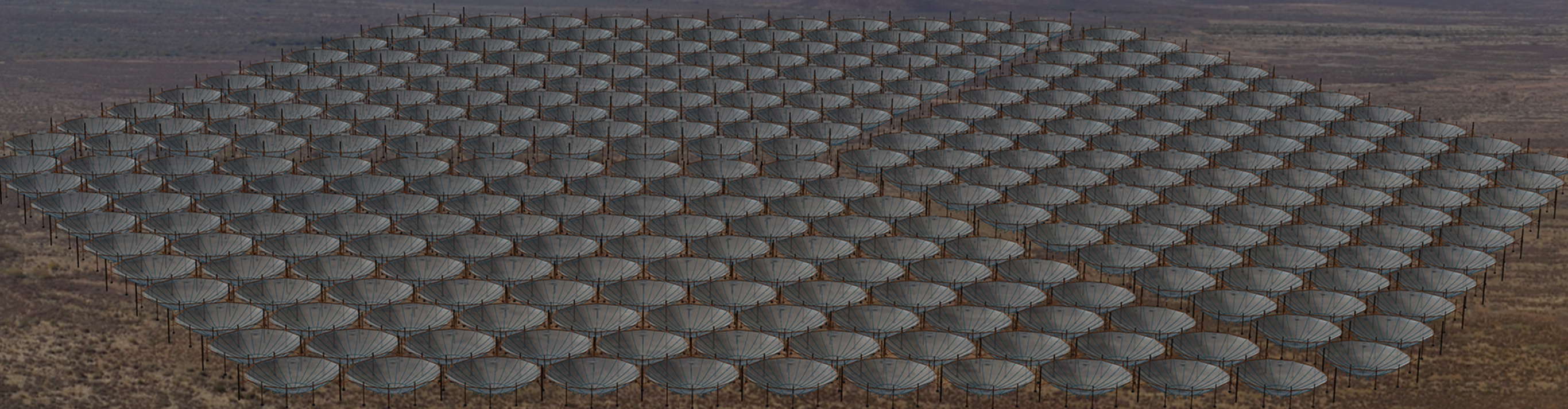
There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Measure antenna
voltages $v_i(t)$.



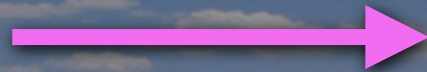
Fourier transform
to frequency: $\tilde{v}_i(\nu)$



There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Measure antenna
voltages $v_i(t)$.

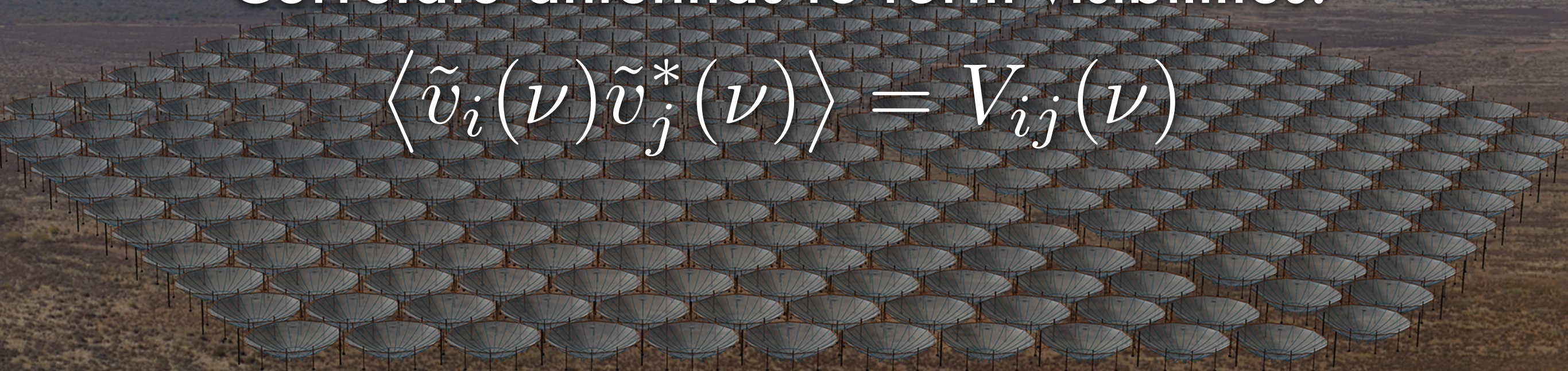


Fourier transform
to frequency: $\tilde{v}_i(\nu)$



Correlate antennas to form visibilities:

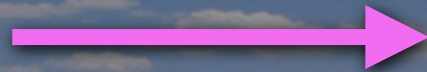
$$\langle \tilde{v}_i(\nu) \tilde{v}_j^*(\nu) \rangle = V_{ij}(\nu)$$



There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Measure antenna
voltages $v_i(t)$.



Fourier transform
to frequency: $\tilde{v}_i(\nu)$

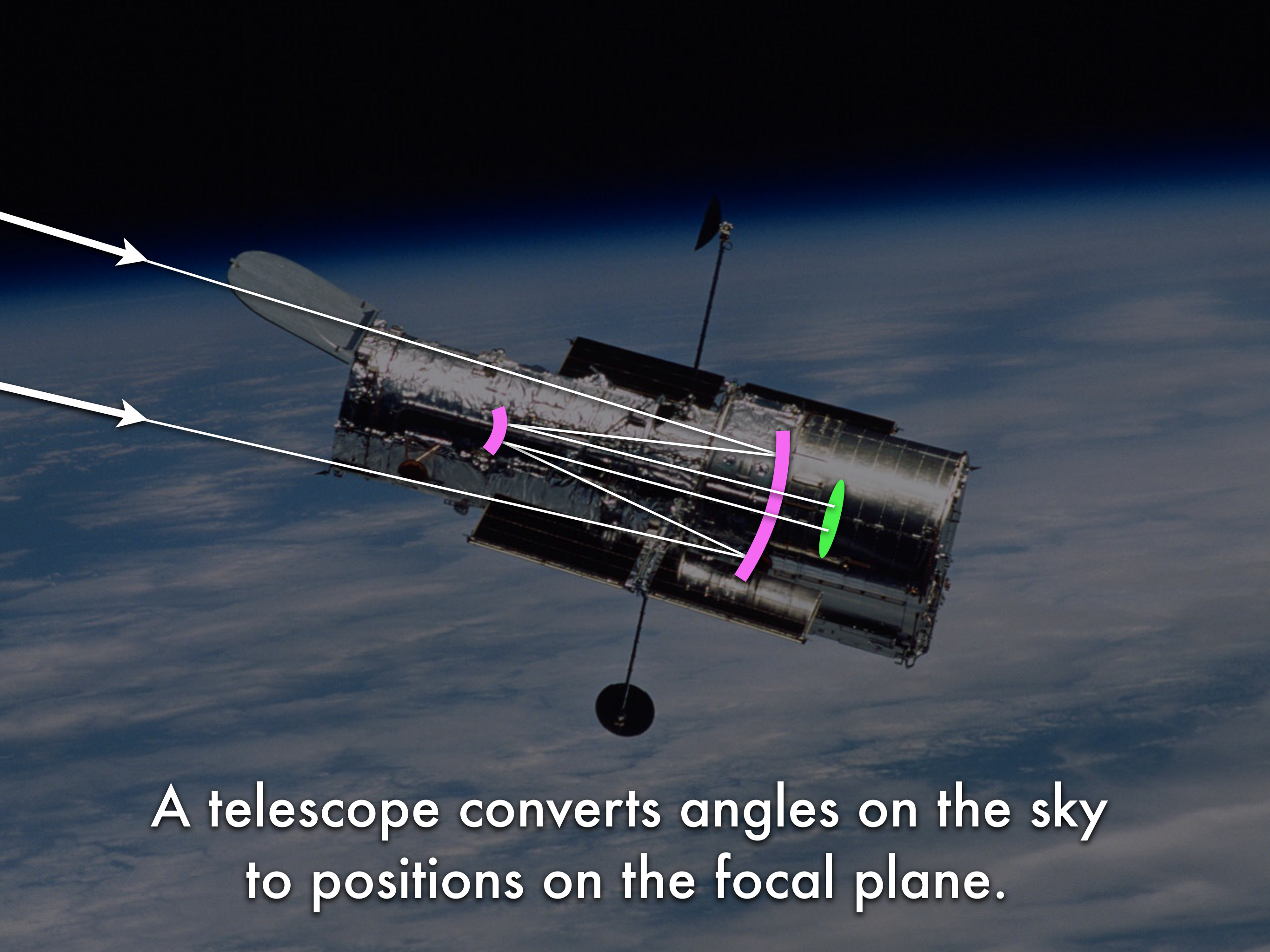


Correlate antennas to form visibilities:

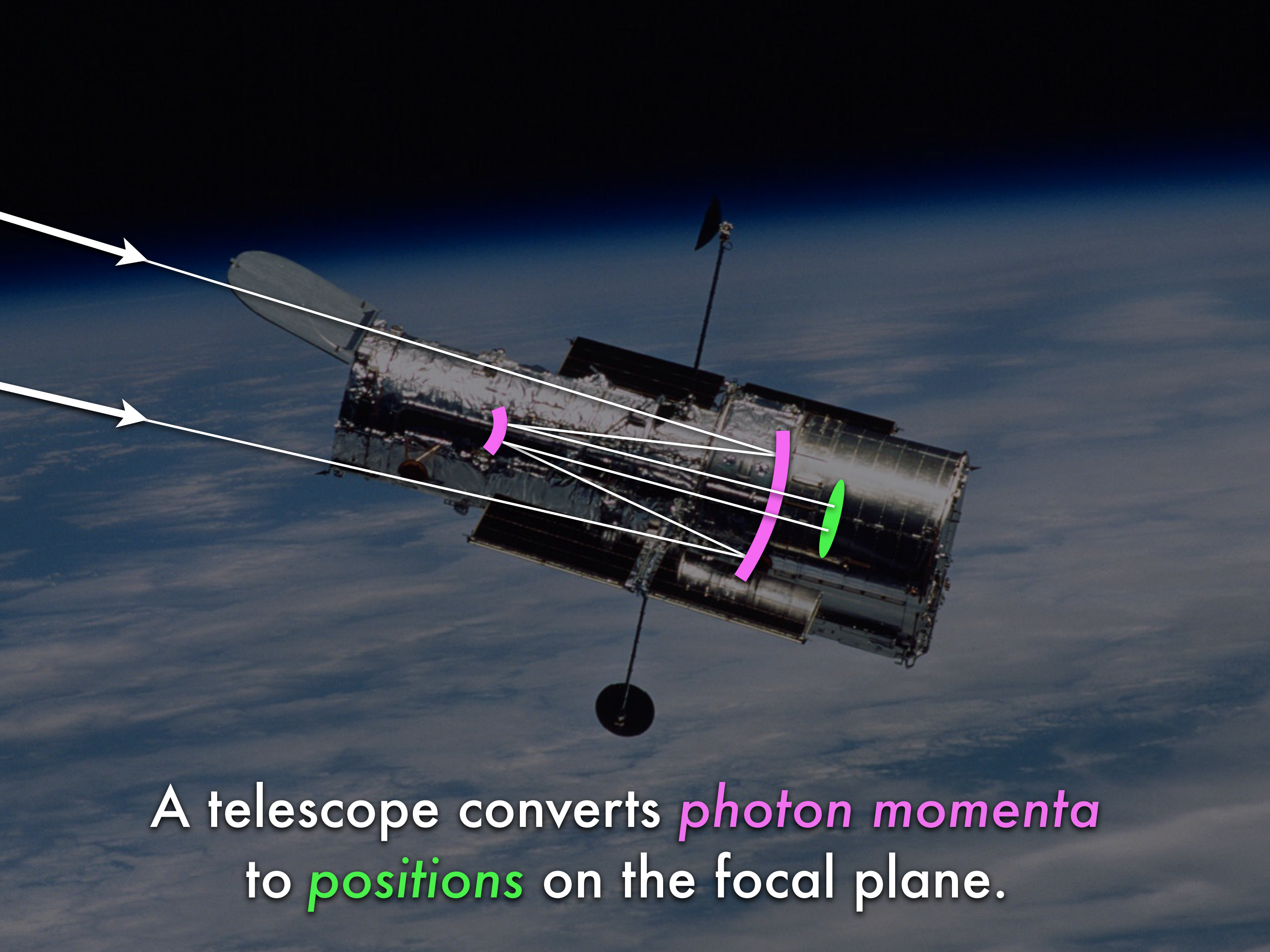
$$\langle \tilde{v}_i(\nu) \tilde{v}_j^*(\nu) \rangle = V_{ij}(\nu)$$

This scales like $O(N^2)$!

All telescopes are
Fourier transformers.



A telescope converts angles on the sky to positions on the focal plane.

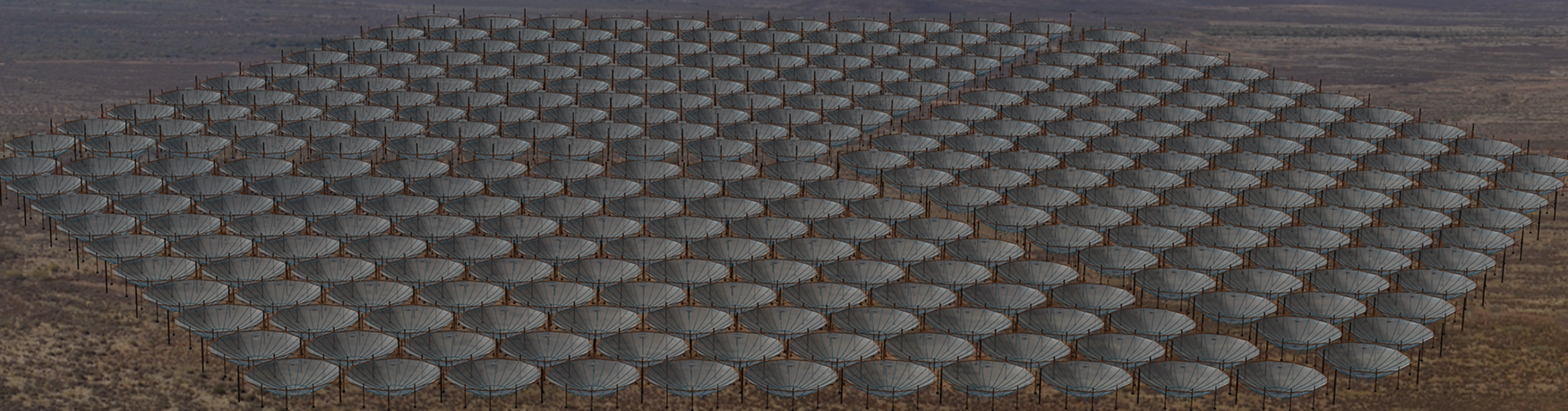


A telescope converts *photon momenta* to *positions* on the focal plane.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

can be rewritten suggestively as...

$$\langle \tilde{v}_i(k) \tilde{v}_j^* \rangle = \int B(\mathbf{k}) I(\mathbf{k}) \exp [i\mathbf{k} \cdot (\mathbf{x}_i - \mathbf{x}_j)] d\Omega$$



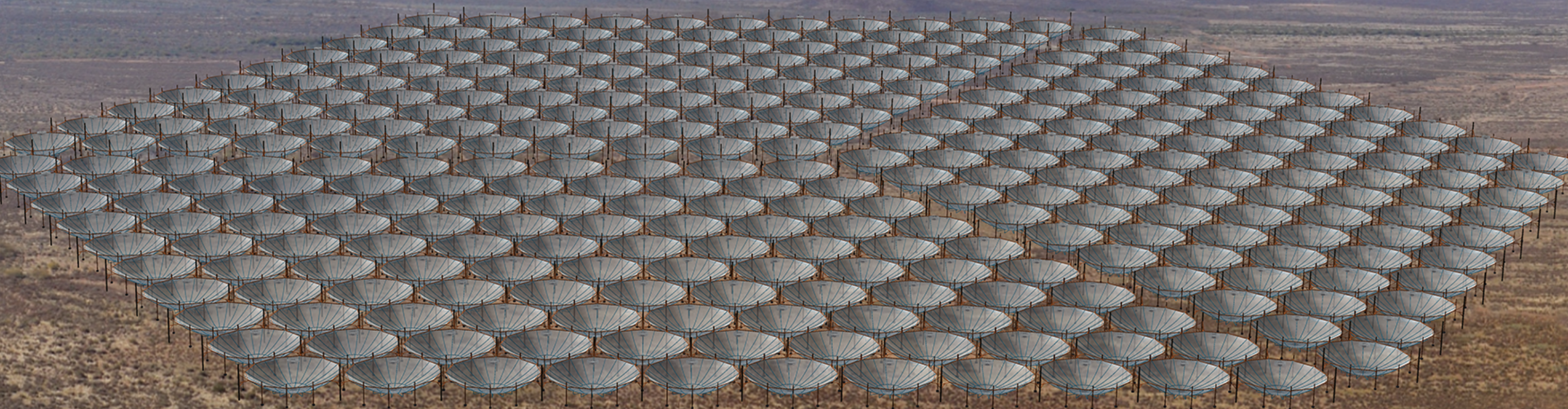
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

can be rewritten suggestively as...

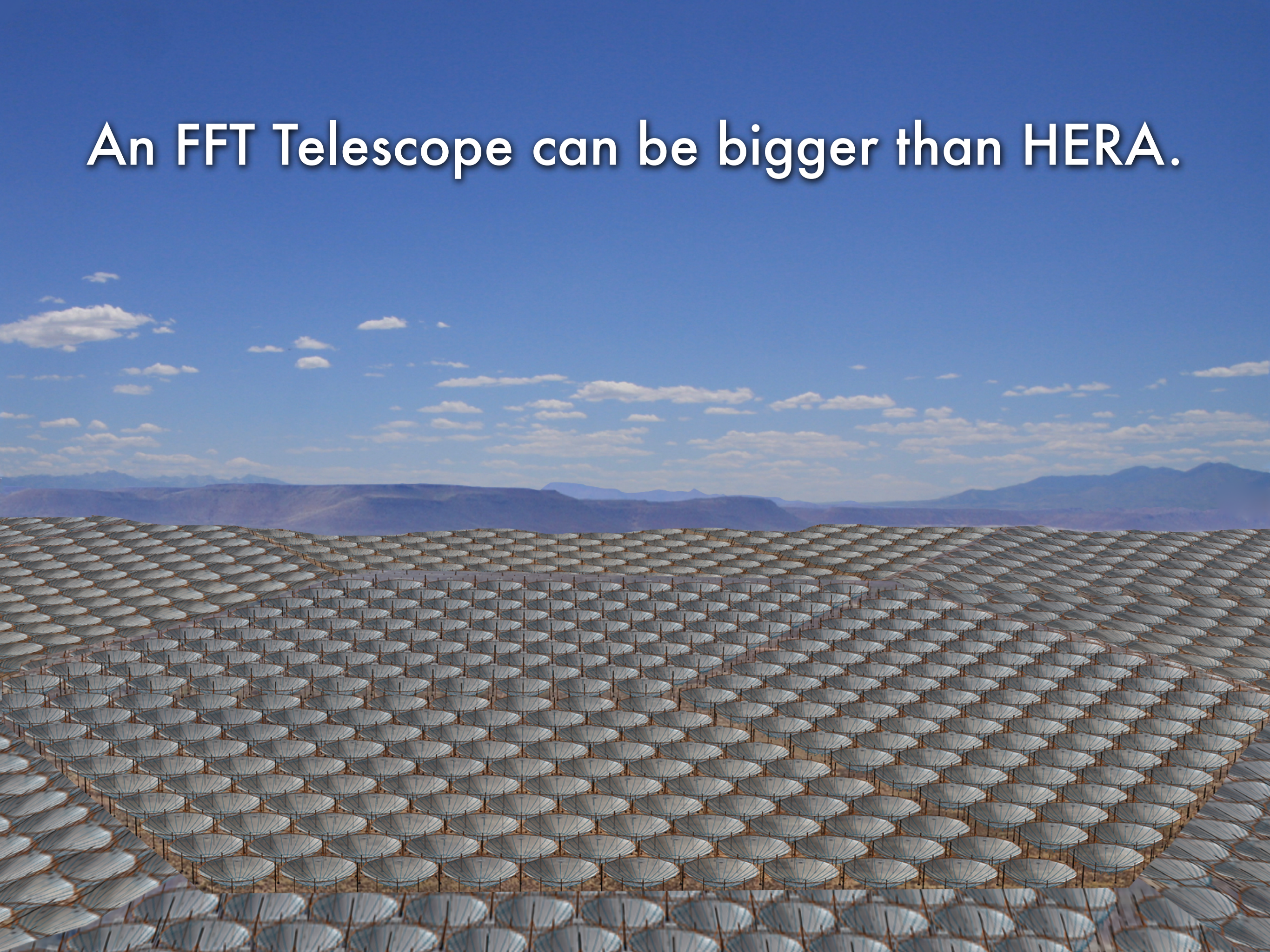
$$\langle \tilde{v}_i(k) \tilde{v}_j^* \rangle = \int B(\mathbf{k}) I(\mathbf{k}) \exp [i\mathbf{k} \cdot (\mathbf{x}_i - \mathbf{x}_j)] d\Omega$$

If antenna positions \mathbf{x}_i are on a regular grid,
we can directly sample the electric field, FFT,
and square to get beam-weighted maps...
effectively correlating in $O(N \log N)$!

An FFT Telescope can be bigger than HERA.

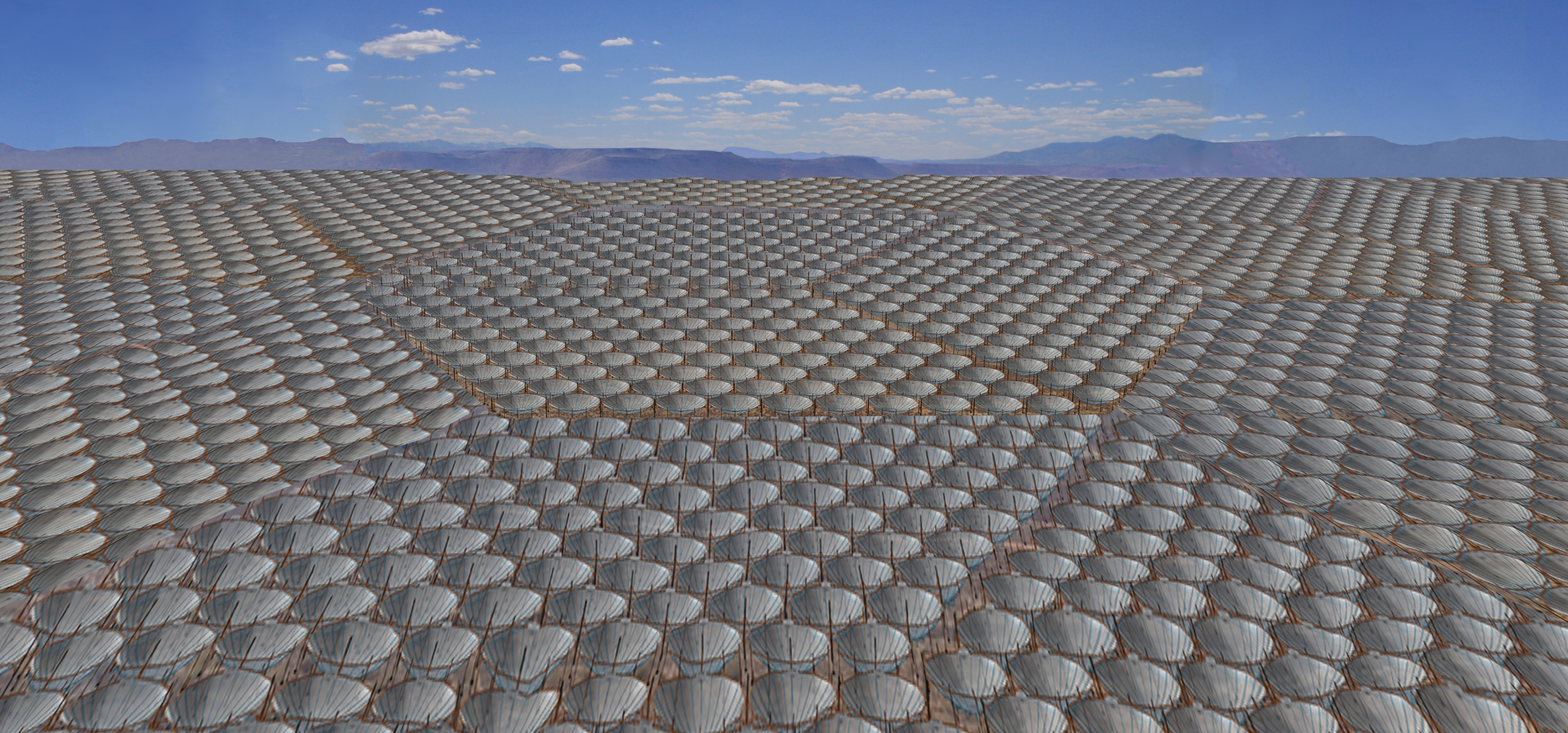


An FFT Telescope can be bigger than HERA.

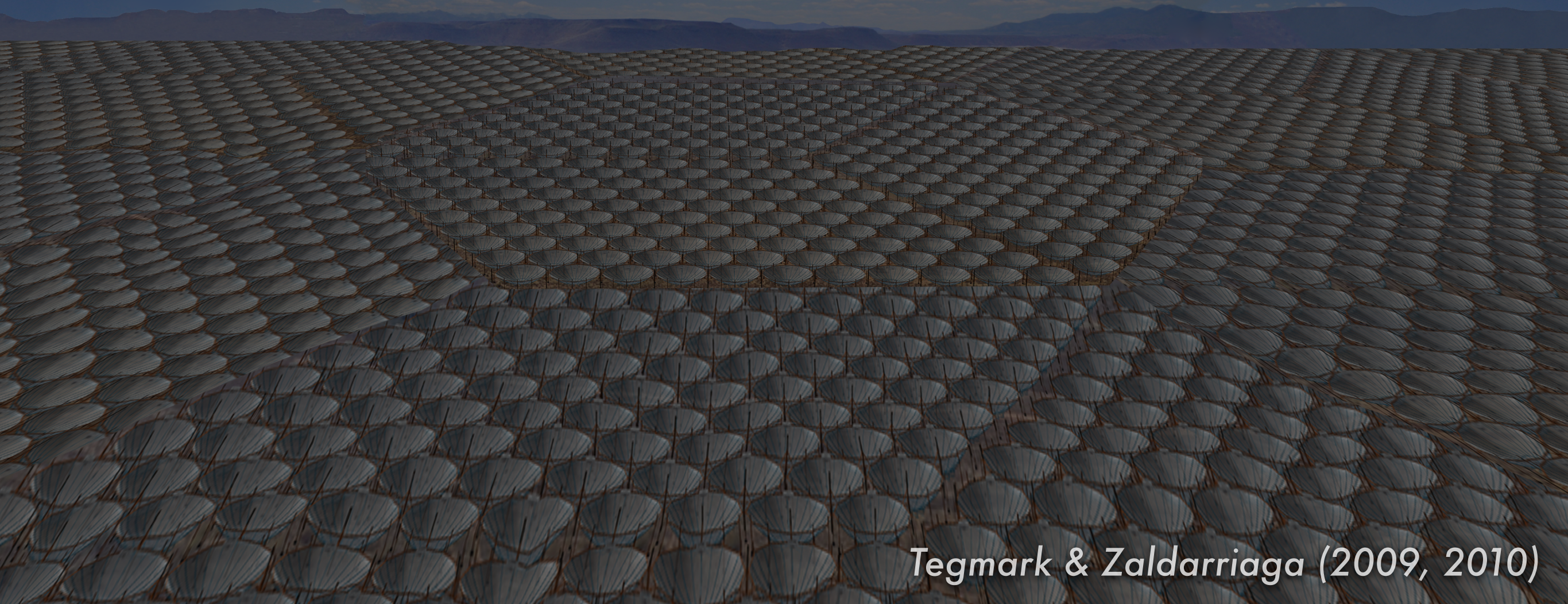


An FFT Telescope can be bigger than HERA.

Much, *much* bigger.



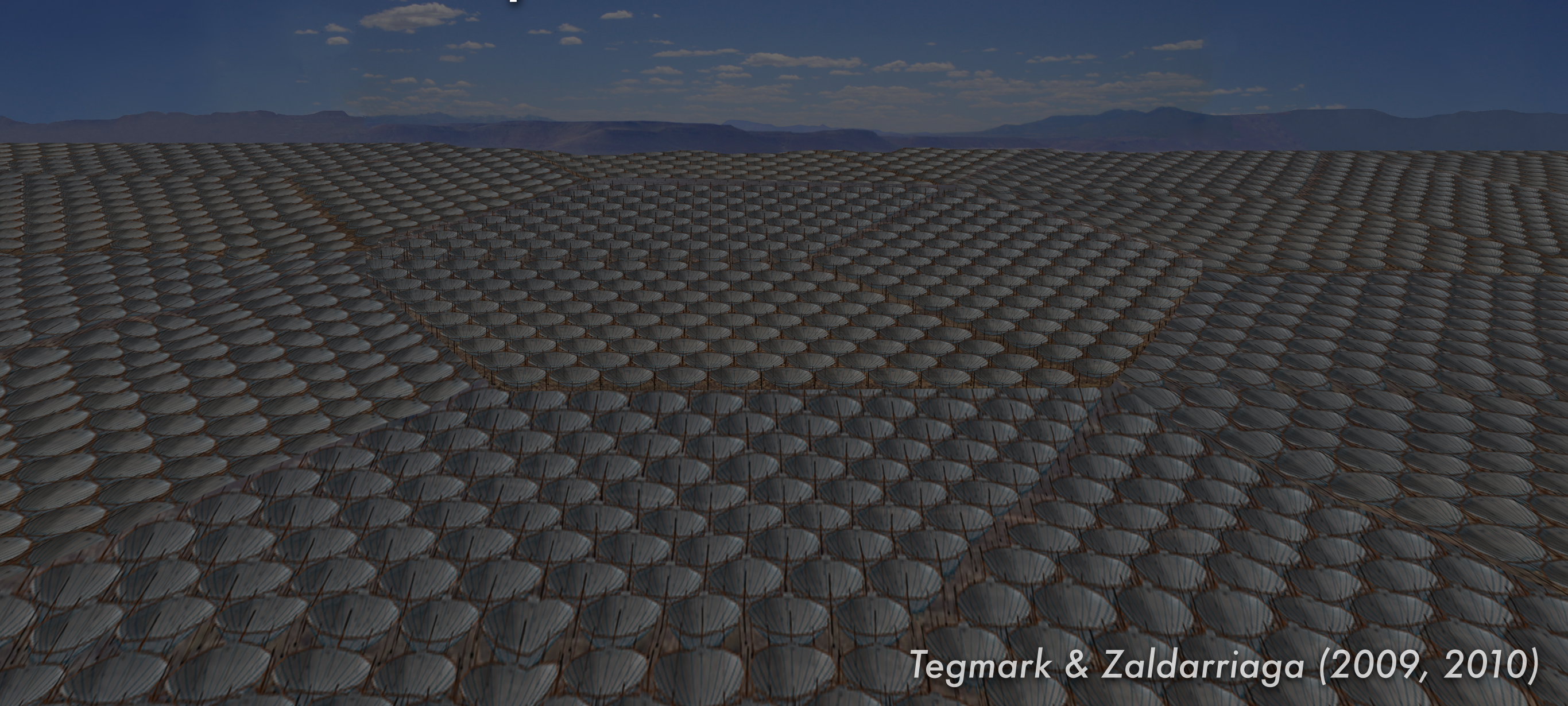
An FFT Telescope needs to be...



Tegmark & Zaldarriaga (2009, 2010)

An FFT Telescope needs to be...

- Co-planar.



An FFT Telescope needs to be...

- Co-planar.
- Made up of identical antenna elements with identical beams.

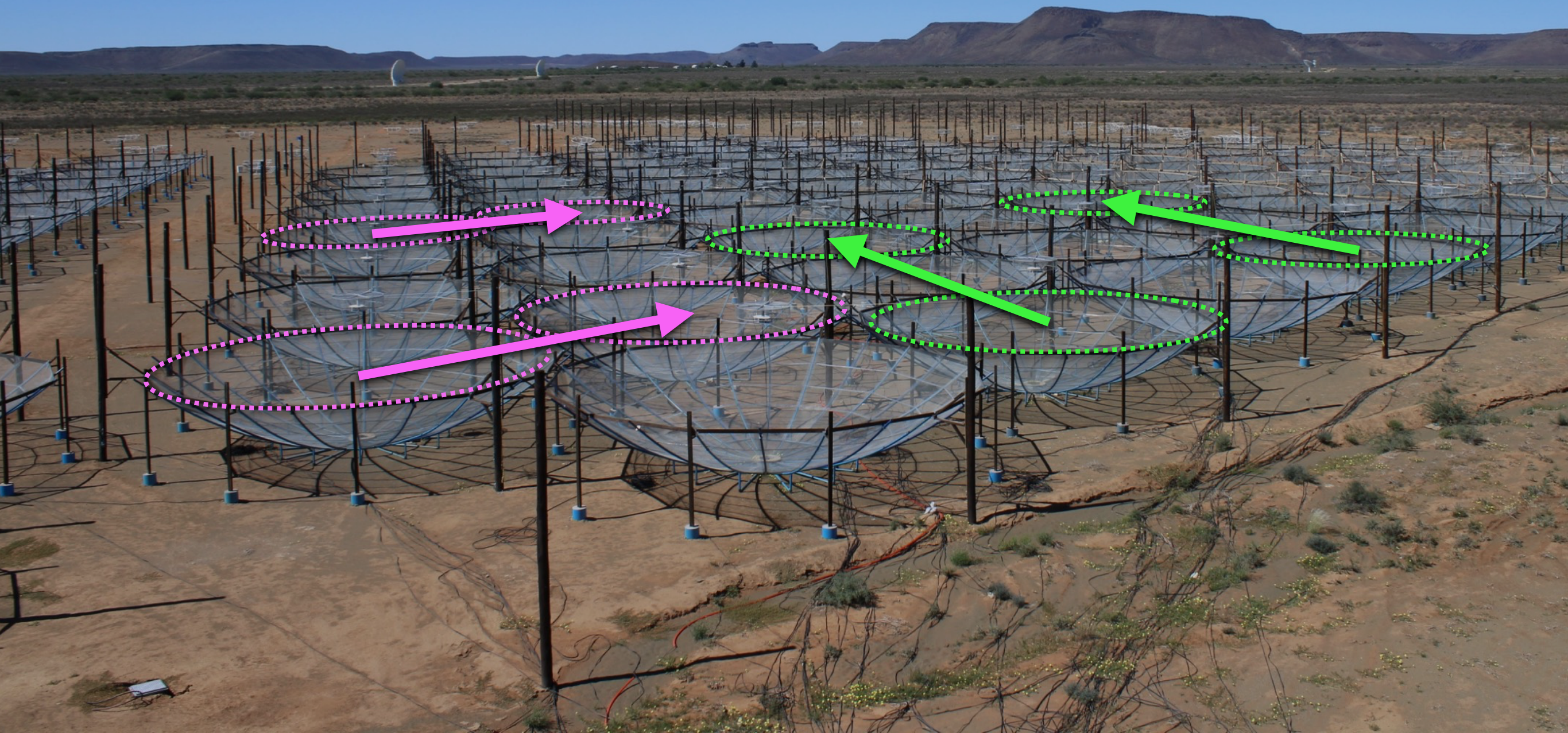
An FFT Telescope needs to be...

- Co-planar.
- Made up of identical antenna elements with identical beams.
- On a regular or hierarchically regular grid.

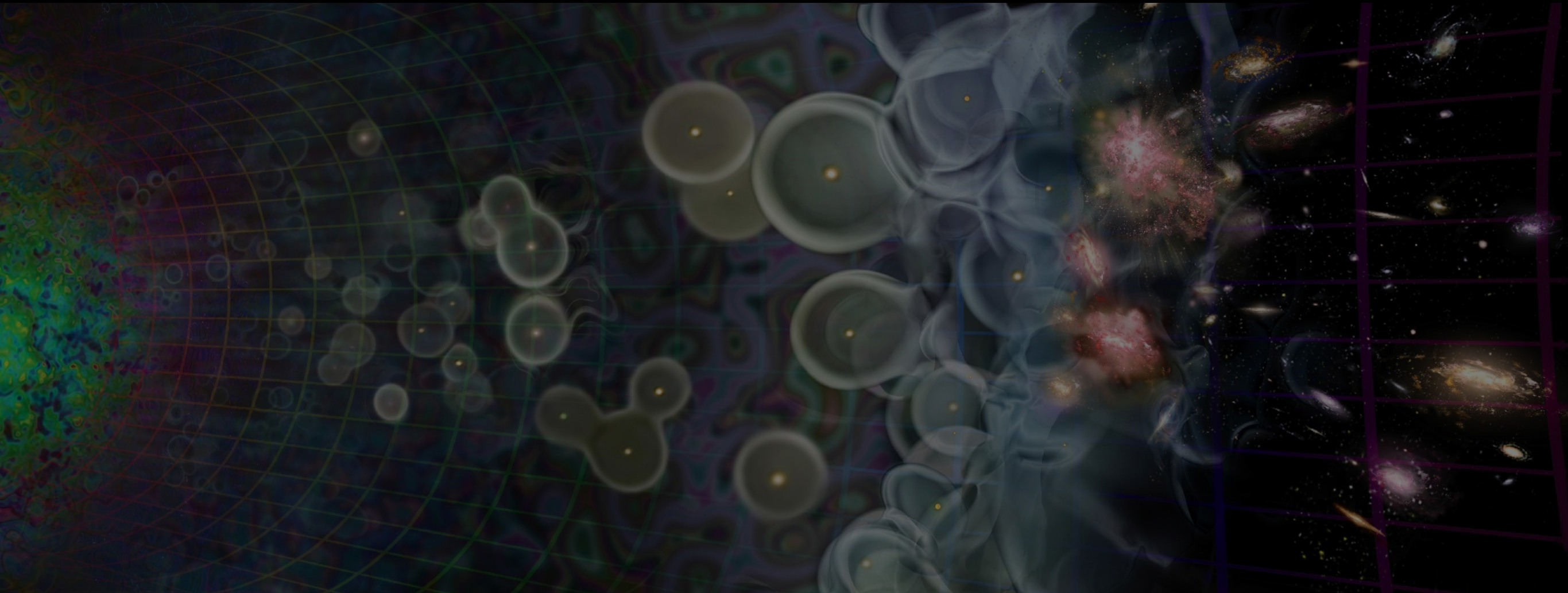
An FFT Telescope needs to be...

- Co-planar.
- Made up of identical antenna elements with identical beams.
- On a regular or hierarchically regular grid.
- Calibrated in real time.

Real-time redundant-baseline calibration of regular arrays is precisely what we're learning to do with HERA!



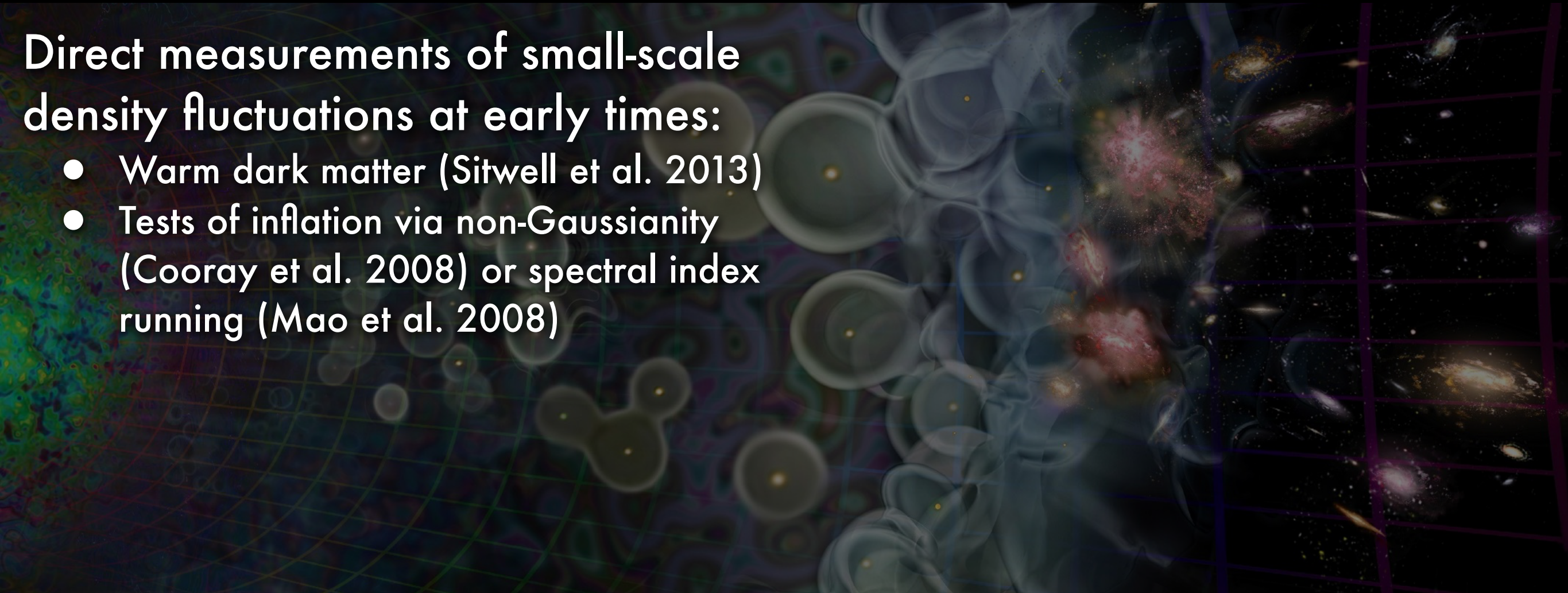
FFTs could map the majority of the volume of the observable universe, giving us...



FFTs could map the majority of the volume of the observable universe, giving us...

Direct measurements of small-scale density fluctuations at early times:

- Warm dark matter (Sitwell et al. 2013)
- Tests of inflation via non-Gaussianity (Cooray et al. 2008) or spectral index running (Mao et al. 2008)



FFTs could map the majority of the volume of the observable universe, giving us...

Direct measurements of small-scale density fluctuations at early times:

- Warm dark matter (Sitwell et al. 2013)
- Tests of inflation via non-Gaussianity (Cooray et al. 2008) or spectral index running (Mao et al. 2008)

A precise thermal history of the universe, constraining:

- Dark matter annihilation and decay (Evoli et al. 2014)
- Primordial black hole evaporation (Mack & Wesley 2008)

FFTs could map the majority of the volume of the observable universe, giving us...

Direct measurements of small-scale density fluctuations at early times:

- Warm dark matter (Sitwell et al. 2013)
- Tests of inflation via non-Gaussianity (Cooray et al. 2008) or spectral index running (Mao et al. 2008)

A precise thermal history of the universe, constraining:

- Dark matter annihilation and decay (Evoli et al. 2014)
- Primordial black hole evaporation (Mack & Wesley 2008)

Unprecedented constraints on the standard model of cosmology:

- Orders of magnitude better than Planck, e.g. $\Delta\Omega_k \approx .0002$ and $\Delta\Sigma v \approx 7$ meV (Mao et al. 2008)

In Summary:



In Summary:

- With a full season of HERA phase I and conservative analysis techniques, we've set a world-leading upper limit on the 21 cm power spectrum and ruled out cold reionization.

In Summary:

- With a full season of HERA phase I and conservative analysis techniques, we've set a world-leading upper limit on the 21 cm power spectrum and ruled out cold reionization.
- With our full array and wider bandwidth, we should have the sensitivity necessary to detect and characterize the 21 cm signal from the EoR and the Cosmic Dawn.

In Summary:

- With a full season of HERA phase I and conservative analysis techniques, we've set a world-leading upper limit on the 21 cm power spectrum and ruled out cold reionization.
- With our full array and wider bandwidth, we should have the sensitivity necessary to detect and characterize the 21 cm signal from the EoR and the Cosmic Dawn.
- One day, an FFTT will draw on the instrumental and analysis legacy of HERA to fulfill the promise of 21 cm cosmology.